# eQE: The embedded Quantum ESPRESSO software package

#### Michele Pavanello & friends

@MikPavanello

Department of Chemistry Rutgers, the State University of New Jersey Newark, NJ

MoISSI workshop at UB 2018

# Pavanello Research Group @ Rutgers



Est. 2012

### Embedding is the future present

### Embedding is the future present

#### The playing field as I see it

- Moore's law forces us to go parallel
- (Super)computer architectures typically provide
  - Compute nodes (12-24 procs/node)
  - Infiniband
  - Coprocessors
- The cloud (probably the future)
  - Low on CPU, high on GPU
  - Google colaboratory / Google cloud: python + tensorflow
  - Not aware of any QC/CMT code tailored to AWS

# Embedding is the future present

#### The playing field as I see it

- Moore's law forces us to go parallel
- (Super)computer architectures typically provide
  - Compute nodes (12-24 procs/node)
  - Infiniband
  - Coprocessors
- The cloud (probably the future)
  - Low on CPU, high on GPU
  - Google colaboratory / Google cloud: python + tensorflow
  - Not aware of any QC/CMT code tailored to AWS

#### The role of Embedding

- Naturally work-parallel
- In some formulations also data-parallel
- Embedding is a paradigm:
  - Can embed mathematical quantities ( $\rho$ ,  $\Psi$ , G)
  - Can embed QC/CMT codes



**Compute nodes** 



Compute nodes





#### The method should:

- Broadcast as little data as possible across compute nodes
- Avoid computing zeros
  - Simple and easy basis set truncation
- Approach excited states and *e<sup>-</sup>*−*N* dynamics

Split the system into (smaller) interacting subsystems

- Split the system into (smaller) interacting subsystems
- Partition of the total electron density into subsystem contributions

$$\rho(\mathbf{r}) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r}) \qquad \rho_I(\mathbf{r}) = \sum_i^\infty n_i^I |\phi_{(i)_I}(\mathbf{r})|^2$$

- Split the system into (smaller) interacting subsystems
- Partition of the total electron density into subsystem contributions

$$\rho(\mathbf{r}) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r}) \qquad \rho_I(\mathbf{r}) = \sum_i^\infty n_i^I |\phi_{(i)_I}(\mathbf{r})|^2$$

The energy functional is almost additive:  $E[\rho] \simeq E[\rho_I] + E[\rho_{II}]$ 

- Split the system into (smaller) interacting subsystems
- Partition of the total electron density into subsystem contributions

$$\rho(\mathbf{r}) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r}) \qquad \rho_I(\mathbf{r}) = \sum_i^\infty n_i^I |\phi_{(i)_I}(\mathbf{r})|^2$$

The energy functional is almost additive:  $E[\rho] \simeq E[\rho_I] + E[\rho_{II}]$ 

 $E_{\text{FDE}}[\rho] = E[\rho_I] + E[\rho_{II}] + T_s^{\text{nadd}}[\rho_I, \rho_{II}] + E_{xc}^{\text{nadd}}[\rho_I, \rho_{II}] + V_{\text{Coul}}^{\text{nadd}}[\rho_I, \rho_{II}]$ 

- Split the system into (smaller) interacting subsystems
- Partition of the total electron density into subsystem contributions

$$\rho(\mathbf{r}) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r}) \qquad \rho_I(\mathbf{r}) = \sum_i^\infty n_i^I |\phi_{(i)_I}(\mathbf{r})|^2$$

The energy functional is almost additive:  $E[\rho] \simeq E[\rho_I] + E[\rho_{II}]$ 

 $E_{\text{FDE}}[\rho] = E[\rho_I] + E[\rho_{II}] + T_s^{\text{nadd}}[\rho_I, \rho_{II}] + E_{xc}^{\text{nadd}}[\rho_I, \rho_{II}] + V_{\text{Coul}}^{\text{nadd}}[\rho_I, \rho_{II}]$ 

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

- Split the system into (smaller) interacting subsystems
- Partition of the total electron density into subsystem contributions

$$\rho(\mathbf{r}) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r}) \qquad \rho_I(\mathbf{r}) = \sum_i^\infty n_i^I |\phi_{(i)_I}(\mathbf{r})|^2$$

The energy functional is almost additive:  $E[\rho] \simeq E[\rho_I] + E[\rho_{II}]$ 

 $E_{\text{FDE}}[\rho] = E[\rho_I] + E[\rho_{II}] + T_s^{\text{nadd}}[\rho_I, \rho_{II}] + E_{xc}^{\text{nadd}}[\rho_I, \rho_{II}] + V_{\text{Coul}}^{\text{nadd}}[\rho_I, \rho_{II}]$ 

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

Frozen Density Embedding (FDE): Coupled Kohn–Sham equations for each subsystem

$$\frac{\delta E_{\text{FDE}}[\rho_I + \rho_{II}]}{\delta \rho_I} = 0 \rightarrow \left[ -\frac{1}{2} \nabla^2 + v_{KS}^I(\mathbf{r}) + v_{emb}^I(\mathbf{r}) \right] \phi_{(i)_I}(\mathbf{r}) = \varepsilon_i^I \phi_{(i)_I}(\mathbf{r})$$

... a slide from 2018 QE workshop...

...a slide from 2018 QE workshop...

The nonadditive functional for two subsystems

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

considering that  $\rho(r) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r})$ .

...a slide from 2018 QE workshop...

The nonadditive functional for two subsystems

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

considering that  $\rho(r) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r})$ .

#### Homework :)

- Compute  $T_s^{nadd}[\rho_I, \rho_{II}]$  in the Thomas-Fermi approximation,  $T_s[\rho] = C_{TF} \int \rho^{5/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_x^{nadd}[\rho_I, \rho_{II}]$  in the Dirac approximation,  $E_x[\rho] = C_x \int \rho^{4/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_H^{nadd}[\rho_I, \rho_{II}], E_H[\rho] = \frac{1}{2} \int \int \rho(\mathbf{r}) \frac{\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r} d\mathbf{r}'$

...a slide from 2018 QE workshop...

The nonadditive functional for two subsystems

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

considering that  $\rho(r) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r})$ .

#### Homework :)

- Compute  $T_s^{nadd}[\rho_I, \rho_{II}]$  in the Thomas-Fermi approximation,  $T_s[\rho] = C_{TF} \int \rho^{5/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_x^{nadd}[\rho_I, \rho_{II}]$  in the Dirac approximation,  $E_x[\rho] = C_x \int \rho^{4/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_H^{nadd}[\rho_I, \rho_{II}], E_H[\rho] = \frac{1}{2} \int \int \rho(\mathbf{r}) \frac{\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r} d\mathbf{r}'$

The embedding potential is given by  $\frac{\delta E^{\text{nadd}}}{\delta \rho_l(\mathbf{r})}$ .

...a slide from 2018 QE workshop...

The nonadditive functional for two subsystems

$$F^{\text{nadd}}[\rho_I, \rho_{II}] = F[\rho] - F[\rho_I] - F[\rho_{II}]$$

considering that  $\rho(r) = \rho_I(\mathbf{r}) + \rho_{II}(\mathbf{r})$ .

#### Homework :)

- Compute  $T_s^{nadd}[\rho_I, \rho_{II}]$  in the Thomas-Fermi approximation,  $T_s[\rho] = C_{TF} \int \rho^{5/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_x^{nadd}[\rho_I, \rho_{II}]$  in the Dirac approximation,  $E_x[\rho] = C_x \int \rho^{4/3}(\mathbf{r}) d\mathbf{r}$
- Compute  $E_H^{nadd}[\rho_I, \rho_{II}], E_H[\rho] = \frac{1}{2} \int \int \rho(\mathbf{r}) \frac{\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r} d\mathbf{r}'$

The embedding potential is given by  $\frac{\delta E^{\text{nadd}}}{\delta \rho_l(\mathbf{r})}$ .

#### ... more homework!

- **T**<sub>s</sub> contribution to  $v_{emb}(\mathbf{r})$
- $E_x$  contribution to  $v_{emb}(\mathbf{r})$
- $E_H$  contribution to  $v_{emb}(\mathbf{r})$



eqe.rutgers.edu

Michele Pavanello & friends @MikPavanello embedded Quantum-ESPRESSO









Rewrite the MPI module of QE	BZ sampling (k-points)
<ul> <li>Subsystem-specific # of CPUs</li> <li>Improved latencies (processes wait for others to complete)</li> <li>Nested DIIS for {ρ<sub>l</sub>(<b>r</b>)}</li> </ul>	<ul> <li>K-point sampling for (semi)conductors.</li> <li>Γ-point for molecules/insulators.</li> </ul>

**Regular QE** 









#### pros & cons

- distributed data communication
- non-polymorphic







- distributed data communication
- non-polymorphic

Michele Pavanello & friends @MikPavanello embedded Quantum-ESPRESSO

gathered data communication

### eQE: a note on coding it



Michele Pavanello & friends @MikPavanello embedded Quantum-ESPRESSO





eQE: An open-source density functional embedding theory code for the condensed phase International Journal of Quantum Chemistry, **117**, e25401 (2017)










#### $30726\,\text{\AA}^3$

#### Speedup compared to regular QE (all PBE)

 $24.5 \times$ 







#### eQE vs QE: Parallel scaling for water 256

Water 256, 256 subsystems



#### QE-eQE gap widens with increasing # of CPUs







#### Diffusion coefficient and dipole moment

$$\langle D \rangle = 2.97(0.4) \cdot 10^{-5} cm^2 s^{-1}$$
  
 $\langle \mu \rangle = 2.8(0.2)D$ 



#### Diffusion coefficient and dipole moment

$$\langle D \rangle = 2.97(0.4) \cdot 10^{-5} cm^2 s^{-1}$$
  
 $\langle \mu \rangle = 2.8(0.2)D$ 

- eQE recovers correct structureeQE recovers correct dynamics
- eQE recovers correct e<sup>-</sup> structure



#### Diffusion coefficient and dipole moment

$$\langle D \rangle = 2.97(0.4) \cdot 10^{-5} cm^2 s^{-1}$$
  
 $\langle \mu \rangle = 2.8(0.2)D$ 

eQE recovers correct structure
eQE recovers correct dynamics
eQE recovers correct e<sup>-</sup> structure
How about e<sup>-</sup> excited states?

#### Subsystem TDDFT and Related Methods

- effective excitations of active subsystem in environment (leads to "FDEu") M.E. Casida, T.A. Wesolowski, Int. J. Quant. Chem. 96 (2004), 577; T.A. Wesolowski, JACS 126 (2004), 11444.
- general subsystem TDDFT formulation for delocalized excitations ("FDEc") J. Neugebauer, J. Chem. Phys. 126 (2007), 134116; J. Neugebauer, J. Chem. Phys. 131 (2009), 084104
- derivation with focus on subsystem response functions M. Pavanello, J. Chem. Phys. 138 (2013), 204118.
- fragment-based TDDFT in the context of partition DFT M.A. Mosquera, D. Jensen, A. Wasserman, *Phys. Rev. Lett.* 111 (2013), 023001.
- time-dependent potential-functional theory for subsystems
   C. Huang, F. Libisch, Q. Peng, E.A. Carter J. Chem. Phys. 140 (2014), 124113.
- real-time subsystem TDDFT

A. Krishtal, D. Ceresoli, M. Pavanello, J. Chem. Phys. 142 (2015), 154116.

subsystem TDDFT with external orthogonality

D.V. Chulhai, L. Jensen, Phys. Chem. Chem. Phys. 18 (2016), 21032.

$$\delta \rho(\mathbf{r},\omega) = \sum_{I} \delta \rho_{I}(\mathbf{r},\omega)$$
 Derives directly from  $\rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r})$ 

$$\delta \rho(\mathbf{r}, \omega) = \sum_{I} \delta \rho_{I}(\mathbf{r}, \omega) \qquad \text{Derives directly from } \rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r})$$
$$\chi(\mathbf{r}, \mathbf{r}', \omega) = \sum_{I} \chi_{I}(\mathbf{r}, \mathbf{r}', \omega) \qquad \text{Spectrum is subsystem-additive ?!?}$$

$$\begin{split} \delta\rho(\mathbf{r},\omega) &= \sum_{I} \delta\rho_{I}(\mathbf{r},\omega) & \text{Derives directly from }\rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r}) \\ \chi(\mathbf{r},\mathbf{r}',\omega) &= \sum_{I} \chi_{I}(\mathbf{r},\mathbf{r}',\omega) & \text{Spectrum is subsystem-additive ?!?} \\ \chi_{I}^{u} &= \chi_{I}^{0} + \chi_{I}^{0} K_{II} \chi_{I}^{u} & \text{Uncoupled/Closed (FDE}_{u}) \end{split}$$

$$\begin{split} \delta\rho(\mathbf{r},\omega) &= \sum_{I} \delta\rho_{I}(\mathbf{r},\omega) & \text{Derives directly from } \rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r}) \\ \chi(\mathbf{r},\mathbf{r}',\omega) &= \sum_{I} \chi_{I}(\mathbf{r},\mathbf{r}',\omega) & \text{Spectrum is subsystem-additive ?!?} \\ \chi_{I}^{u} &= \chi_{I}^{0} + \chi_{I}^{0} K_{II} \chi_{I}^{u} & \text{Uncoupled/Closed (FDE}_{u}) \end{split}$$

 $\chi_I = \chi_I^u$ 

$$\begin{split} \delta\rho(\mathbf{r},\omega) &= \sum_{I} \delta\rho_{I}(\mathbf{r},\omega) & \text{Derives directly from }\rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r}) \\ \chi(\mathbf{r},\mathbf{r}',\omega) &= \sum_{I} \chi_{I}(\mathbf{r},\mathbf{r}',\omega) & \text{Spectrum is subsystem-additive ?!?} \\ \chi_{I}^{u} &= \chi_{I}^{0} + \chi_{I}^{0}K_{II}\chi_{I}^{u} & \text{Uncoupled/Closed (FDE}_{u}) \\ \chi_{I} &= \chi_{I}^{u} + \sum \chi_{I}^{u}K_{IJ}\chi_{J} & \text{Coupled/Open (FDE}_{c}) - \text{Holographic theory} \end{split}$$

The K<sub>IJ</sub> coupling

 $J \neq I$ 

$$K_{IJ}(\mathbf{r}_1, \mathbf{r}_2, t - t') = \frac{\delta(t - t')}{|\mathbf{r}_1 - \mathbf{r}_2|} + \frac{\delta^2 E_{\rm xc}}{\delta\rho(\mathbf{r}_1, t)\delta\rho(\mathbf{r}_2, t')} + \frac{\delta^2 T_{\rm s}}{\delta\rho(\mathbf{r}_1, t)\delta\rho(\mathbf{r}_2, t')}$$

$$\begin{split} \delta\rho(\mathbf{r},\omega) &= \sum_{I} \delta\rho_{I}(\mathbf{r},\omega) & \text{Derives directly from } \rho(\mathbf{r}) = \sum_{I} \rho_{I}(\mathbf{r}) \\ \chi(\mathbf{r},\mathbf{r}',\omega) &= \sum_{I} \chi_{I}(\mathbf{r},\mathbf{r}',\omega) & \text{Spectrum is subsystem-additive ?!?} \\ \chi_{I}^{u} &= \chi_{I}^{0} + \chi_{I}^{0} K_{II} \chi_{I}^{u} & \text{Uncoupled/Closed (FDE}_{u}) \\ \chi_{I} &= \chi_{I}^{u} + \sum_{I} \chi_{I}^{u} K_{IJ} \chi_{J} & \text{Coupled/Open (FDE}_{c}) - \text{Holographic theory} \end{split}$$

The K<sub>IJ</sub> coupling

 $J \neq I$ 

$$K_{U}(\mathbf{r}_{1},\mathbf{r}_{2},t-t') = \frac{\delta(t-t')}{|\mathbf{r}_{1}-\mathbf{r}_{2}|} + \frac{\delta^{2}E_{\mathrm{xc}}}{\delta\rho(\mathbf{r}_{1},t)\delta\rho(\mathbf{r}_{2},t')} + \frac{\delta^{2}T_{\mathrm{s}}}{\delta\rho(\mathbf{r}_{1},t)\delta\rho(\mathbf{r}_{2},t')}$$

#### Subsystem TDDFT is a general theory for open systems Let's apply it to liquid water!

Cooperation and Environment Characterize the Low-Lying Optical Spectrum of Liquid Water

J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

#### Cooperation and Environment Characterize the Low-Lying Optical Spectrum of Liquid Water

J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

Details of the calculation

Real-time subsystem TD-DFT, coded in eQE A. Krishtal, D. Ceresoli, M. Pavanello JCP 142, 154116 (2015)

Cooperation and Environment Characterize the Low-Lying Optical Spectrum of Liquid Water

J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

#### Details of the calculation

- Real-time subsystem TD-DFT, coded in eQE
   A. Krishtal, D. Ceresoli, M. Pavanello JCP 142, 154116 (2015)
- Ultrasoft pseudopotentials
   X. Qian, J. Li, X. Lin, S. Yip PRB 73, 035408 (2006)

#### Cooperation and Environment Characterize the Low-Lying Optical Spectrum of Liquid Water

J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

#### Details of the calculation

- Real-time subsystem TD-DFT, coded in eQE
   A. Krishtal, D. Ceresoli, M. Pavanello JCP 142, 154116 (2015)
- Ultrasoft pseudopotentials
   X. Qian, J. Li, X. Lin, S. Yip PRB 73, 035408 (2006)
- Reduce the # of Plane Waves by 80%
   A. Genova, D. Ceresoli, M. Pavanello JCP 144, 234105 (2016)

#### Cooperation and Environment Characterize the Low-Lying Optical Spectrum of Liquid Water

J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

#### Details of the calculation

- Real-time subsystem TD-DFT, coded in eQE
   A. Krishtal, D. Ceresoli, M. Pavanello JCP 142, 154116 (2015)
- Ultrasoft pseudopotentials
   X. Qian, J. Li, X. Lin, S. Yip PRB 73, 035408 (2006)
- Reduce the # of Plane Waves by 80%
   A. Genova, D. Ceresoli, M. Pavanello JCP 144, 234105 (2016)
- $\Gamma$  point for each water molecule







# Computation of ϵ 10 snapshots of water 64 Average of 640 spectra Match experimental sum-rule in 0-25 eV window






### Computation of $\epsilon$

- 10 snapshots of water 64
- Average of 640 spectra
- Match experimental sum-rule in 0-25 eV window

### We find...

- Overall good agreement across wide window of ω
- KS Exciton binding energy
  - $\rightarrow$  KS gap  $\sim$  7.0 eV
  - $\rightarrow$  Optical gap  $\sim$  6.4 eV



Kumar P., S. & Genova, A. & MP, J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

### Computation of $\epsilon$

- 10 snapshots of water 64
- Average of 640 spectra
- Match experimental sum-rule in 0-25 eV window

### We find...

- Overall good agreement across wide window of ω
- KS Exciton binding energy
  - $\rightarrow$  KS gap  $\sim$  7.0 eV
  - $\rightarrow$  Optical gap  $\sim$  6.4 eV
- Index of refraction n = 1.68



Kumar P., S. & Genova, A. & MP, J. Phys. Chem. Lett., 8 (20), pp 5077-5083 (2017)

### Dielectric Constant - Many Body Effects

Michele Pavanello & friends @MikPavanello embedded Quantum-ESPRESSO

### Dielectric Constant - Many Body Effects

 One Body (closed):

$$\chi = \sum_{I} \chi_{I}^{u}$$

### Dielectric Constant – Many Body Effects

 One Body (closed):

$$\chi = \sum_{I} \chi_{I}^{u}$$

Many Body (open):

$$\chi = \sum_{I} \chi_{I}^{u} + \sum_{I \neq J} \chi_{I}^{u} K_{IJ} \chi_{J}$$









3 Many-body effects do not change first peak position



- 1 Oscillator strength enhanced in low-lying excitations
- 2 Oscillator strength depleted in high-lying excitations
- 3 Many-body effects do not change first peak position
- 4 Consistent with Galli's explanation for increased index of refraction with pressure<sup>a</sup>

### First Absorption Band of Liquid Water: $\omega_1$

### Role of the Environment: Urbach tail and $\omega_1$ peak position

Michele Pavanello & friends @MikPavanello embedded Quantum-ESPRESSO

## First Absorption Band of Liquid Water: $\omega_1$

### Role of the Environment: Urbach tail and $\omega_1$ peak position



- Urbach tail entirely due to environment
- ω<sub>1</sub> and OH stretching do not correlate in the liquid
- Environment Order Parameter (EOP): Regression of 300+ descriptors
- Correlation of EOP to  $\omega_1$  is 0.65

### Environment Order Parameter: Simplified depiction



### Environment Order Parameter: Simplified depiction



#### Environment Order Parameter: Simplified depiction



#### What is "Environment Order Parameter" made of?

Accepted Hydrogen Bonds is the highest contribution to EOP

#### Environment Order Parameter: Simplified depiction



- Accepted Hydrogen Bonds is the highest contribution to EOP
- Distances and angles within the 1<sup>st</sup> solvation shell

#### Environment Order Parameter: Simplified depiction



- Accepted Hydrogen Bonds is the highest contribution to EOP
- Distances and angles within the 1<sup>st</sup> solvation shell
- Linear and Neural Network regressions lead to the same result

#### Environment Order Parameter: Simplified depiction



- Accepted Hydrogen Bonds is the highest contribution to EOP
- Distances and angles within the 1<sup>st</sup> solvation shell
- Linear and Neural Network regressions lead to the same result
- Same correlations between EOP and KS eigenvalue difference

#### Environment Order Parameter: Simplified depiction



### What is "Environment Order Parameter" made of?

- Accepted Hydrogen Bonds is the highest contribution to EOP
- Distances and angles within the 1<sup>st</sup> solvation shell
- Linear and Neural Network regressions lead to the same result
- Same correlations between EOP and KS eigenvalue difference

... we have dissected liquid water...

### Acknowledgments

### Postdocs, Students & Collaborators

#### **Current PRG members**

Postdocs:

- Dr. Wenhui Mi
- Dr. Muhammed Acikgoz
- Dr. Pablo Ramos

Graduate Students:

- Rupali Chawla
- Alina Umerbekova
- Jack Maranhao

Funding:

- NSF CAREER
- DOE CTC

Alumni & Collaborators

#### Alumni:

- Johannes Tölle (Münster University)
- Dr. Alessandro Genova (@ Kitware)
- Prof. Alisa Krishtal (@ NJIT)
- Dr. Debalina Sinha (@ L'Oreal)

Collaborators:

- Dr. Davide Ceresoli (CNR)
- Prof. Rob DiStasio (Cornell)
- Dr. Andre Gomes (CNRS)
- Prof. Oliviero Andreussi (North Texas)
- Prof. Henk Eshuis (Montclair State)
- Dr. Damien Riedel (Paris Sud)