



# Aperiodic Fragments in Periodic Solids

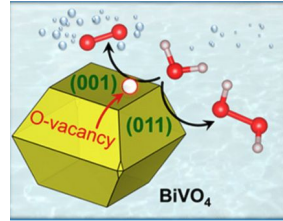
Rob Lavroff, Daniel Kats, Lorenzo Maschio, Nikolay Bogdanov, Ali Alavi, Anastassia Alexandrova, and Denis Usvyat



VISTA Seminar, 9/18/2024

# Defects in solids: a playground for all kinds of chemistry

- Catalysis active sites:
  - Oxygen vacancies in metal-oxide catalysts promote thermal and photocatalysis
  - Single-atom dopants (ie gr-N4 + TM) are at the forefront of electrocatalysis



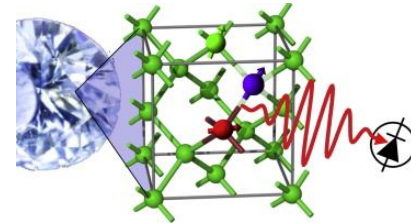
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- Quantum information:
  - NV centers in diamond and silicon as sensors and spin qubits
  - Carbon dopants in hBN as SPEs for quantum communication



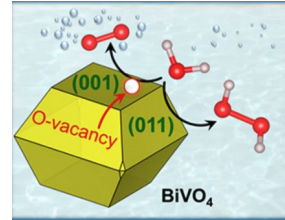
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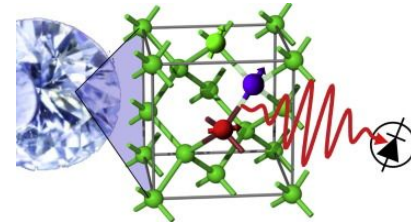
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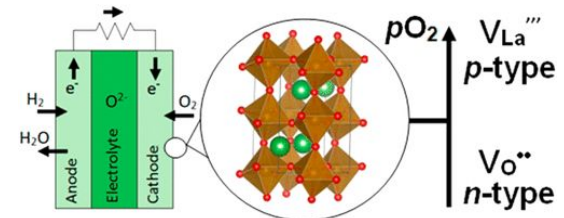
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- **Quantum information:**
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- **Clean energy:**
  - Tuning transport in batteries and semiconductors via defect engineering
  - Defects as anchoring sites in fuel cells



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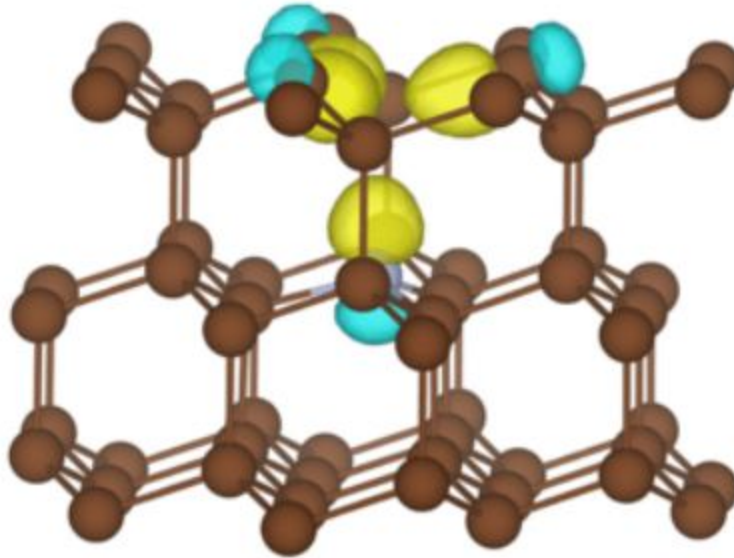
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# The role of theory in defect chemistry

- Ab-initio calculations can circumvent trial-and-error experiments
  - Binding energies, excited states, spin-couplings, dynamics, etc etc
  - Insights often inaccessible or difficult to obtain via experiment: direct wavefunction information



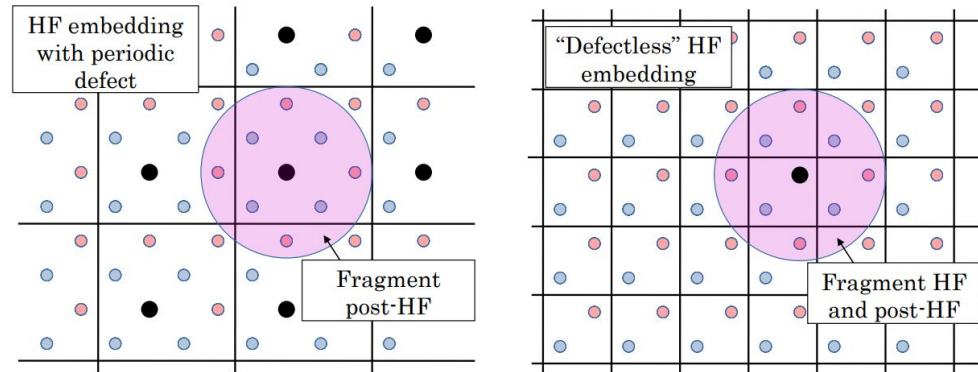
PhysRevB.105.235104

# The role of theory in defect chemistry

- Ab-initio calculations can circumvent trial-and-error experiments
  - Binding energies, excited states, spin-couplings, dynamics, etc etc
  - Insights often inaccessible or difficult to obtain via experiment: direct wavefunction information
- But can current models always give an accurate picture of defects?
  - Relies on periodic boundary conditions, defect is always repeated across unit cells
  - Defects can be spaced at nano- or microscale, routine DFT unit cells < a few nanometers
  - Wavefunction methods (MP2, CC) are more trustworthy, but much smaller unit cells required
  - If a defect is charged (i.e.  $NV^-$  center in diamond), need to throw out diverging terms
- The ever-evolving field of **quantum embedding**
  - Can treat defects with near-exact wavefunction methods, pair to cheaper surroundings
  - Suffers from same issue of defect repetition in PBC...

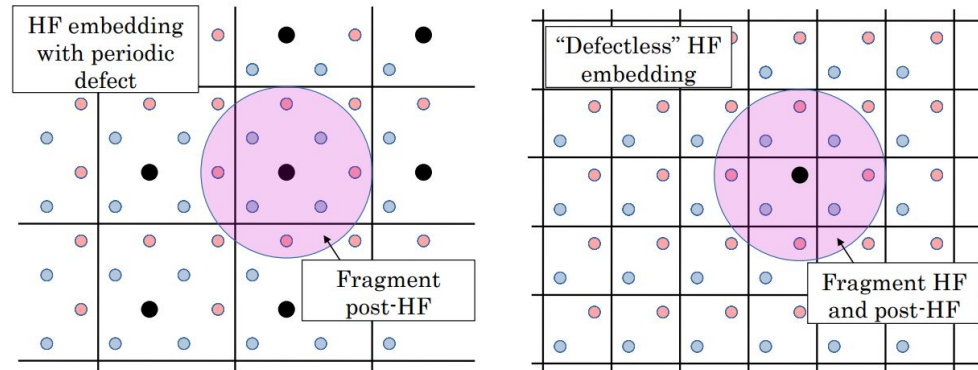
# A quantum embedding approach without defect repetition

- Introducing an “aperiodic fragment” approach for the defect could
  - Allow the solid’s unit cell used to be only large enough to house the defect -> **savings**
  - Eliminate the term removal of “compensating background charges” -> **trustworthiness**
  - Assure that defects in different unit cells aren’t artificially interacting -> **accuracy**



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  - Assure that defects in different unit cells aren’t artificially interacting -> **accuracy**
- This approach would need to
  - Converge to the same properties as the periodic approach in the thermodynamic limit (TDL)
  - Take into account the relaxation of the environment due to the defect
  - Have flexibility in ab-initio approaches to tackle a variety of defects





# Aperiodic fragment in frozen mean-field surroundings

- Embedding often uses HF for surroundings: no double counting of  $E_{\text{corr}}$ 
  - Self-consistency or expansion of correlated fragment fixes lack of correlation later
- Fragment HF energy same form as a molecule's, but with modified  $h$  and  $E_{\text{nuc}}$

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$$h_{\mu'\nu'}^{\text{frag}} = \left\langle \mu' \left| -\frac{1}{2} \nabla^2 \right| \nu' \right\rangle - \left\langle \mu' \left| \sum_{K' \in \text{frag}} \frac{Z_{K'}}{|\mathbf{r} - \mathbf{R}_{K'}|} \right| \nu' \right\rangle \\ - \left\langle \mu' \left| \sum_{K \notin \text{frag}} \frac{Z_K}{|\mathbf{r} - \mathbf{R}_K|} \right| \nu' \right\rangle + \sum_{i \notin \text{frag}} [2(\mu'\nu'|ii) - (\mu'i|iv')] \\ = F_{\mu'\nu'}^{\text{per}} - \sum_{i \in \text{frag}} [2(\mu'\nu'|ii) - (\mu'i|iv')] \\ + \sum_{K \in \text{frag}} \left\langle \mu' \left| \frac{Z_K}{|\mathbf{r} - \mathbf{R}_K|} \right| \nu' \right\rangle - \sum_{K' \in \text{frag}} \left\langle \mu' \left| \frac{Z_{K'}}{|\mathbf{r} - \mathbf{R}_{K'}|} \right| \nu' \right\rangle$$

$\mu', \nu', \dots$ : fragment's basis orbitals;  $K, L, \dots$ : atoms of the fragment before introduction of the defect;  $K', L', \dots$ : atoms of the fragment after introduction of the defect;  $i, j, \dots$ : the occupied orbitals before the fragment's SCF;  $i', j', \dots$ : the occupied orbitals after the fragment's SCF;  $V(\mathbf{R}_{K'})$ : the periodic electrostatic potential at the point  $\mathbf{R}_{K'}$ ,

# Aperiodic fragment in frozen surroundings

- Embedding often uses HF
  - Self-consistent
- Fragment SCF

counting of  $E_{\text{corr}}$   
 correlation later  
 with modified  $h$  and  $E_{\text{nuc}}$

$E_{\text{HF}}^{\text{frag}}$

In the Fock matrix, we remove the contributions from the pristine crystal inside the fragment and add those of the "defected" crystal, then do SCF

$$\begin{aligned}
 & \left\langle \mu' \left| \sum_{K' \in \text{frag}} \frac{Z_{K'}}{|\mathbf{r} - \mathbf{R}_{K'}|} \right| \nu' \right\rangle \\
 & \left\langle \mu' \left| \sum_{K \notin \text{frag}} \frac{Z_K}{|\mathbf{r} - \mathbf{R}_K|} \right| \nu' \right\rangle + \sum_{i \notin \text{frag}} [2(\mu' \nu' | ii) - (\mu' i | i \nu')] \\
 & = F_{\mu' \nu'}^{\text{per}} - \sum_{i \in \text{frag}} [2(\mu' \nu' | ii) - (\mu' i | i \nu')] \\
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# Aperiodic fragment in frozen mean-field surroundings

- Embedding often in frozen mean-field surroundings: no double counting of  $E_{\text{corr}}$ 
  - Self-consistency: fragment fixes lack of correlation later
- Fragment HF energy:  $E_{\text{HF}}^{\text{frag}}$  but with modified  $h$  and  $E_{\text{nuc}}$

$$E_{\text{HF}}^{\text{frag}} = 2 \sum_{i' \in \text{frag}} h_{i'i'} + \frac{1}{2} \sum_{i' \in \text{frag}} \sum_{j' \in \text{frag}} \langle i' | \sum_{\mu' \in \text{env}} \frac{Z_{K'}}{|\mathbf{r} - \mathbf{R}_{K'}|} | j' \rangle + \frac{1}{2} \sum_{K' \in \text{frag}} \sum_{L' \in \text{frag}} \frac{Z_{K'}}{|\mathbf{R}_{K'} - \mathbf{R}_{L'}|} + \sum_{K' \in \text{frag}} Z_{K'} \cdot V(\mathbf{R}_{K'}) + 2 \sum_{i \in \text{frag}} \left\langle i \left| \sum_{K' \in \text{frag}} \frac{Z_{K'}}{|\mathbf{r} - \mathbf{R}_{K'}|} \right| i \right\rangle - \sum_{L \in \text{frag}} \sum_{K' \in \text{frag}} \frac{Z_L Z_{K'}}{|\mathbf{R}_L - \mathbf{R}_{K'}|} + \sum_{K \in \text{frag}} \dots$$

In the total energy we remove the environment: pristine-fragment interactions and add the environment: "defected"-fragment interactions

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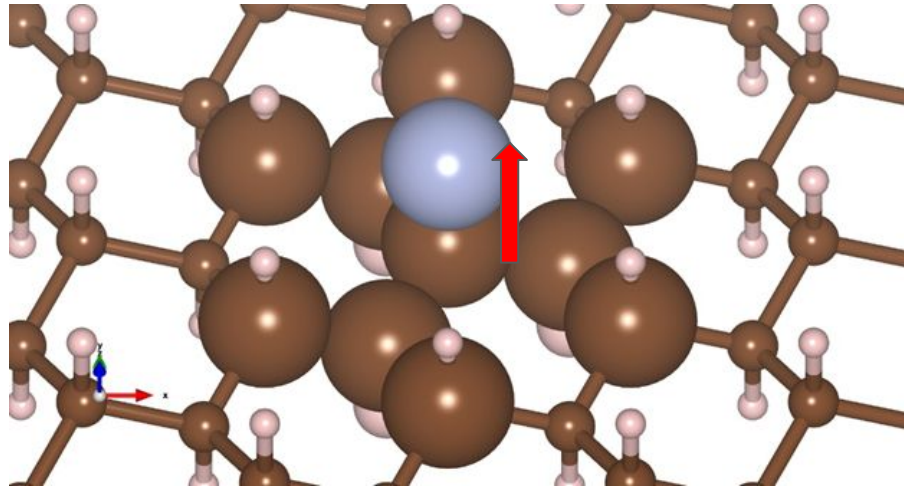
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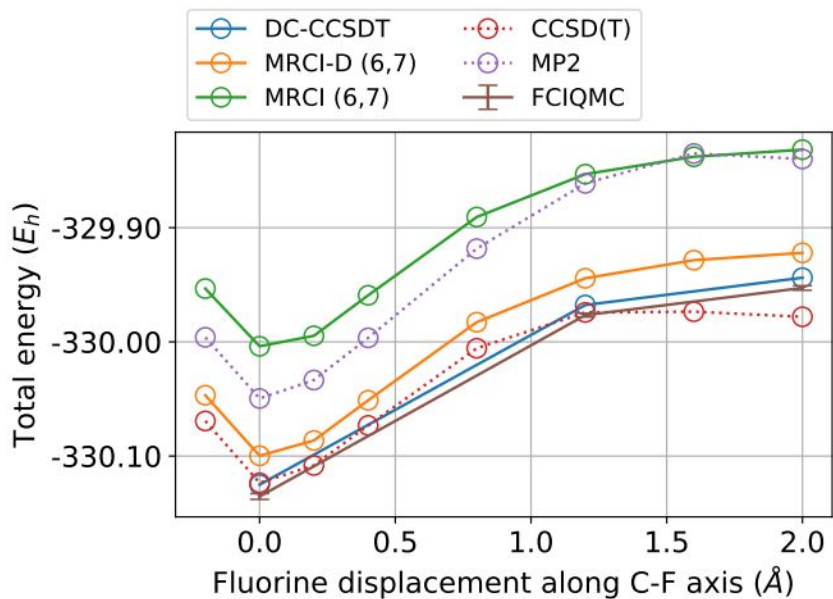
# Putting things to the test: Fluorographane

- Real material (10.1038/s41565-019-0582-z) and an interesting testing ground
- Substitute a H in graphane with a F, then start breaking the C-F bond
  - As bond is stretched, static correlation increases, dynamic correlation of F valence electrons



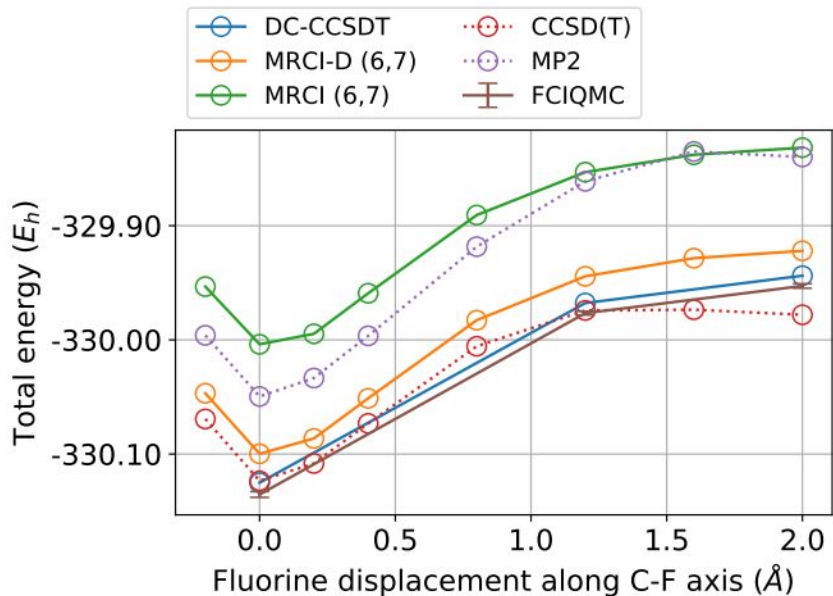


# Periodic embedding vs aperiodic

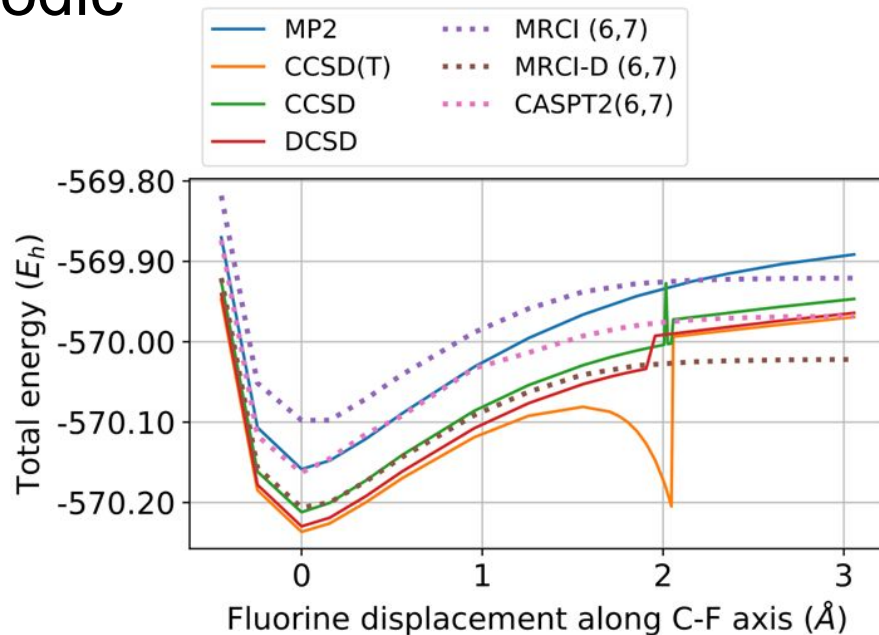


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- Expected behavior of MP2 and CCSD(T) in strong correlation

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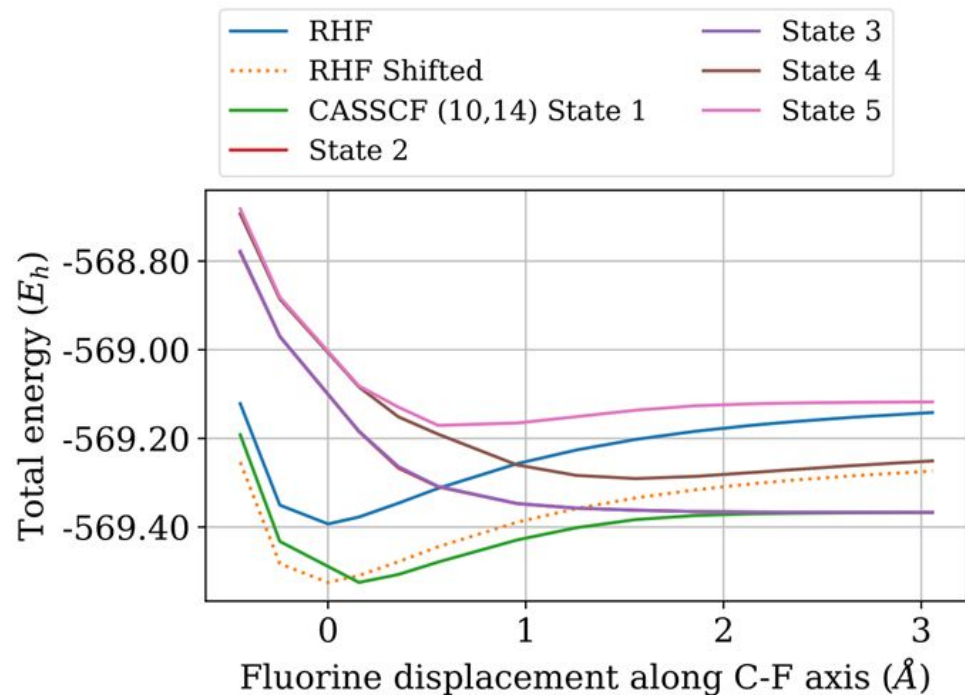


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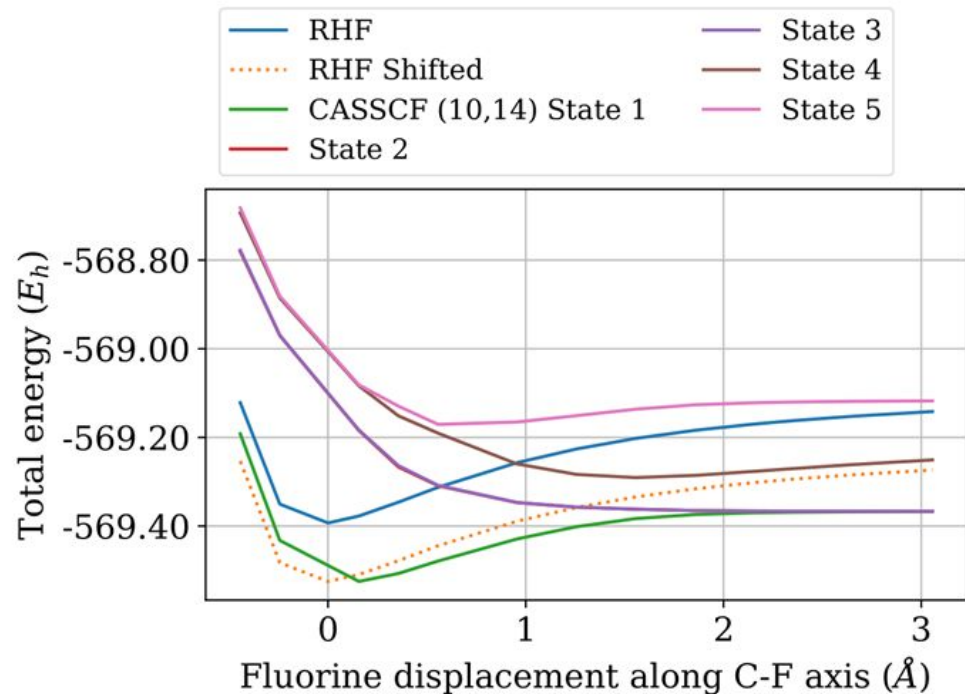
- Aperiodic approach, same 14-atoms
- MP2 behaves like HF, CC/DC has major issues at 2Å but like MP2 after
  - An avoided crossing?

# A crossing! Does aperiodic reproduce periodic in the TDL?

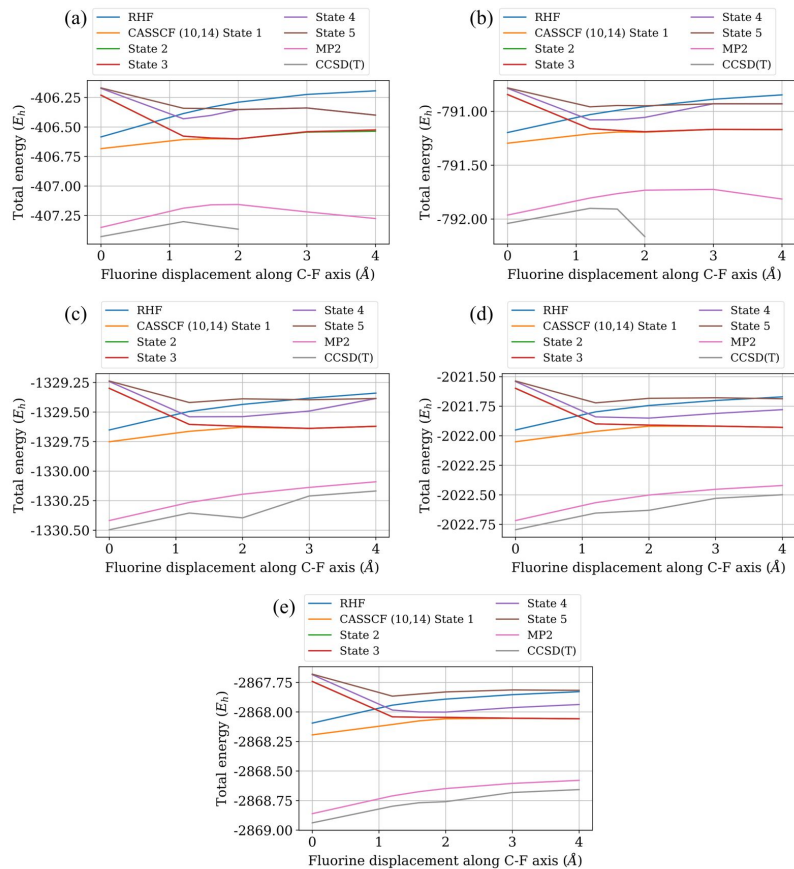


Using aperiodic embedding, RHF zips along a covalent-ionic crossing (follows ionic diabats)

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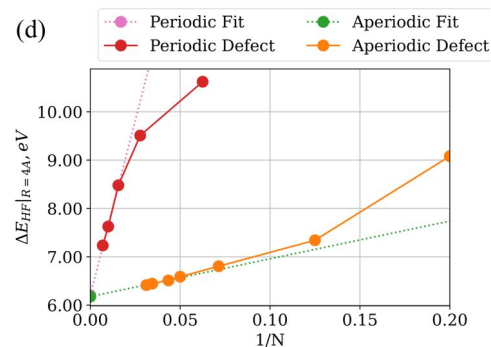
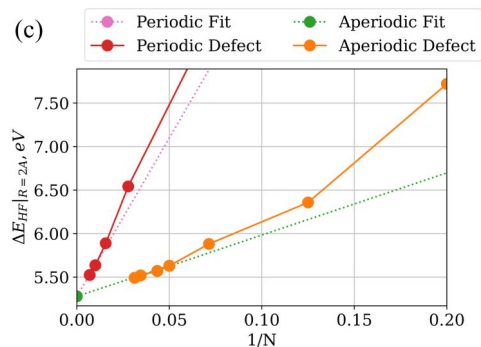
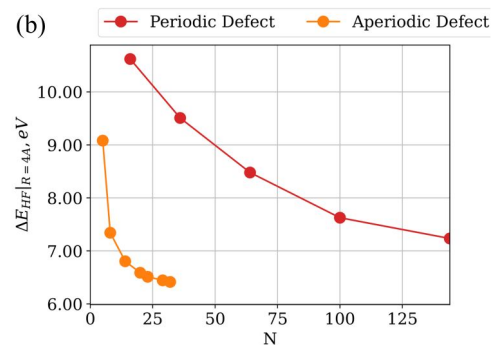
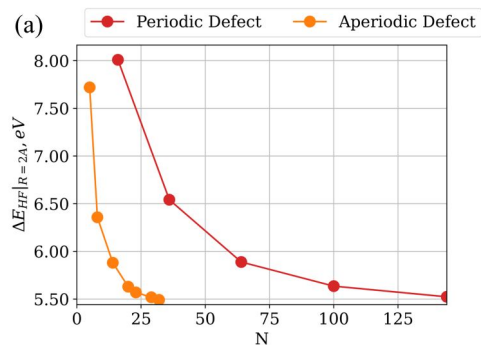


As we take periodic scheme from 2x2 to 6x6, excited state manifold (and MP2/CC) match!

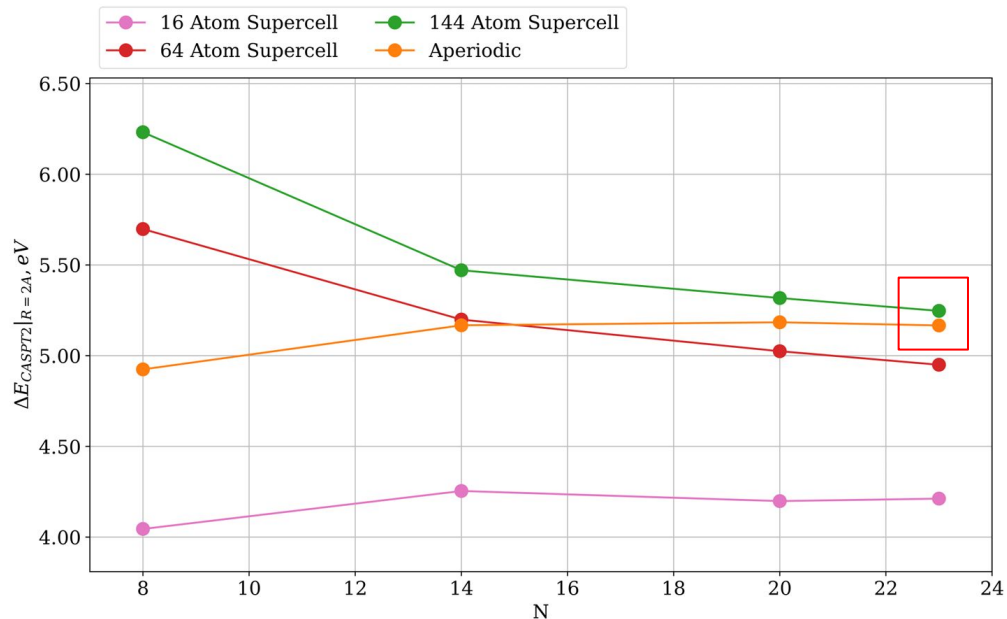
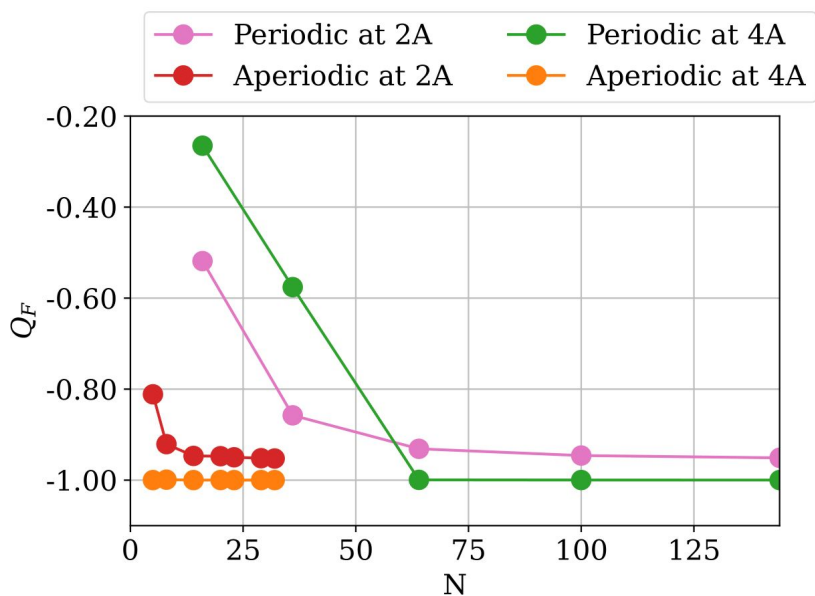


# What about quantitative agreement in the TDL?

- HF dissociation energy at bond displacements of 2 and 4 Å at N atoms
- At large enough fragments/supercells, a linear regime in 1/N begins: exact match



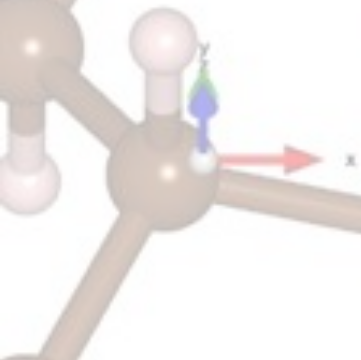
# HF Mulliken F Charge and CASPT2 Dissociation Energies



For CASPT2, an embedded fragment must be used for periodic too, so  $N$ =number of atoms in fragment

# Conclusions and next steps

- Aperiodic embedding can:
  - Give correct ground and excited state PES using a minimal unit cell, while periodic can require very large supercells to get here
  - Converge to the same TDL values of dissociation energies and Mulliken charges as periodic
    - Periodic is over 400x more costly for this 2-D system
  - Capture response of the environment in this system by simply expanding the fragment
    - HF's ionic dissociation induces a large dipole, yet TDL properties match fully periodic ones





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- **Future steps/works in progress:**
  - Testing on charged systems and defects in 3-D crystals (is the speedup even greater?)
  - Analytic gradients of the fragment HF energy: fast geometry optimizations of defects
  - Computing SOCs and NACs: heavy atoms and nonadiabatic dynamics of defects
  - Self-consistency between fragment and environment to allow smaller fragments
    - Applying the approach to metals - fragment is an open system

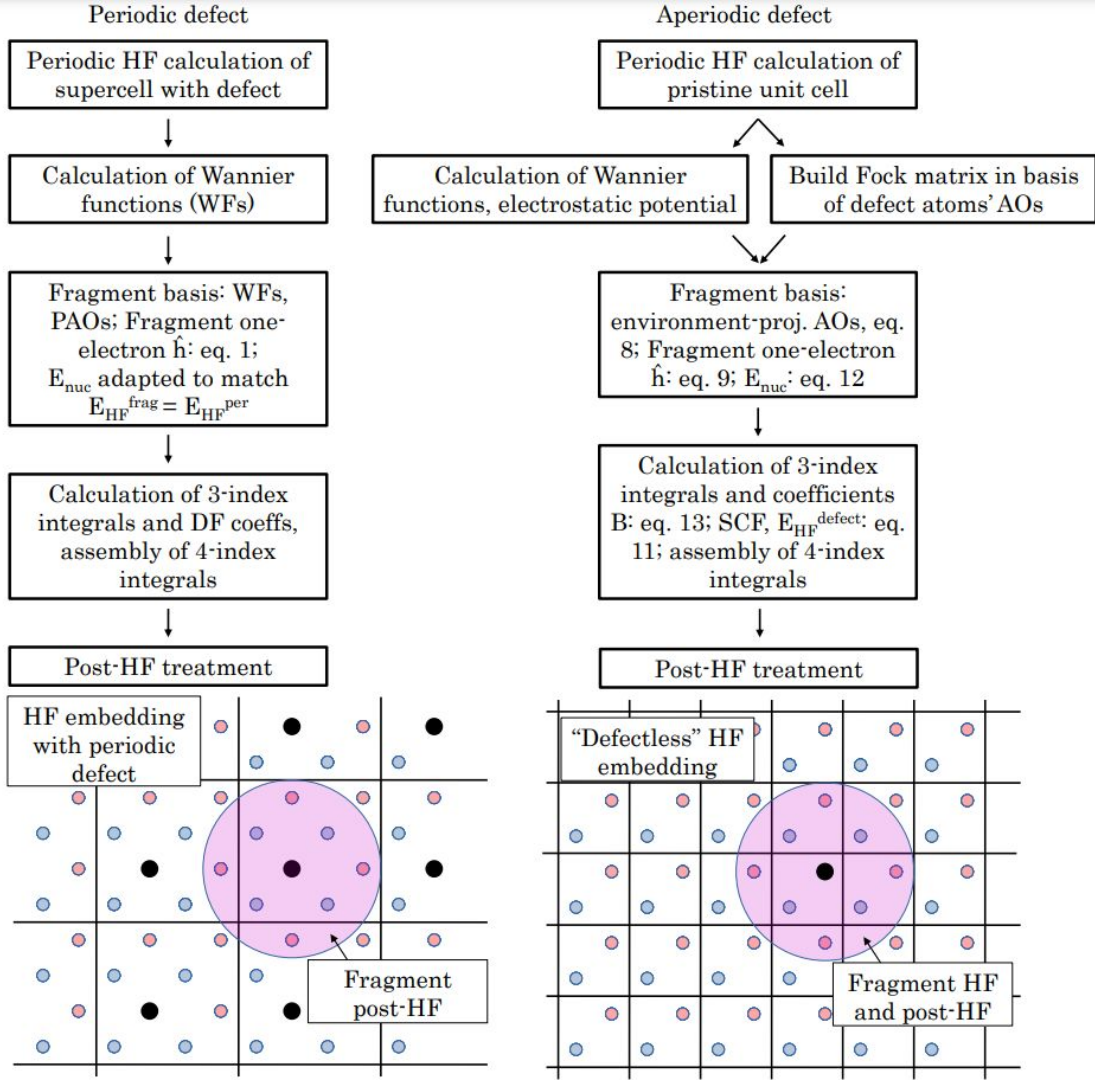


The background of the slide is a complex, repeating molecular structure. It consists of a network of interconnected spheres and rods. The spheres are primarily light brown or tan, with some smaller white spheres interspersed. The rods are thin and connect the spheres in a lattice-like pattern. The overall appearance is that of a crystalline or molecular lattice. In the bottom-left corner, there is a small, detailed inset of a molecular structure. This inset shows a central atom with a blue lobe above it and a green lobe below it, with a red arrow pointing to the right, labeled with the letter 'x'.

# Thank you! Questions?

arXiv:2406.03373

# Workflow



Definition of the electrostatic potential and  $E_{\text{per}}^{\text{HF}}$

$$Z_K \left[ -2 \sum_i \left\langle i \left| \frac{1}{|\mathbf{r} - \mathbf{R}_K|} \right| i \right\rangle + \sum_L' \frac{Z_L}{|\mathbf{R}_K - \mathbf{R}_L|} \right]$$

$$= Z_K \cdot V(\mathbf{R}_K)$$

$$E_{\text{HF}}^{\text{per}} = 2 \sum_{i \in \text{cell}} h_{ii} + \sum_{i \in \text{cell}} (2J_{ii}^{\text{per}} - K_{ii}^{\text{per}}) + E_{\text{nuc}}^{\text{per}}$$

$$= 2 \sum_{i \in \text{cell}} \left\langle i \left| -\frac{1}{2} \nabla^2 \right| i \right\rangle + 2 \sum_{i \in \text{cell}} \left\langle i \left| -\sum_K \frac{Z_K}{|\mathbf{r} - \mathbf{R}_K|} \right| i \right\rangle$$

$$+ \sum_{i \in \text{cell}} \sum_j [2(ii|jj) - (ij|ji)] + \frac{1}{2} \sum_{K \in \text{cell}} \sum_L \frac{Z_K Z_L}{|\mathbf{R}_K - \mathbf{R}_L|} \cdot$$

