

Non-equilibrium electron dynamics from real-time quantum embedding

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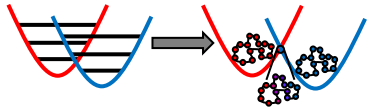
Kretchmer Group: New simulation methods for electron dynamics in complex environments

Method Development

Multi-faceted research group working at the interface of electronic structure and quantum dynamics

Classically Isomorphic Methods

Map a complicated quantum system to a computationally tractable classical system

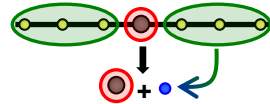


Quantum mechanics Classical mechanics

Real-Time Electronic Structure

Solve the time *dependent* Schrödinger equation – Quantum embedding

Subsystem coupled to *large* environment



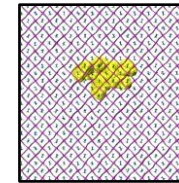
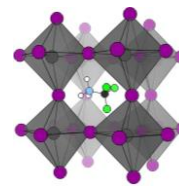
Subsystem coupled to *small* environment



Transport Dynamics in Correlated Materials

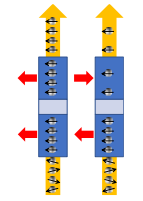
Examine electron-nuclear, electron-electron, and spin-spin interactions in governing the non-equilibrium transport of electrons

Charge-transport in perovskites



Explicit electron in fully atomistic simulation

Spin-transport across interfaces



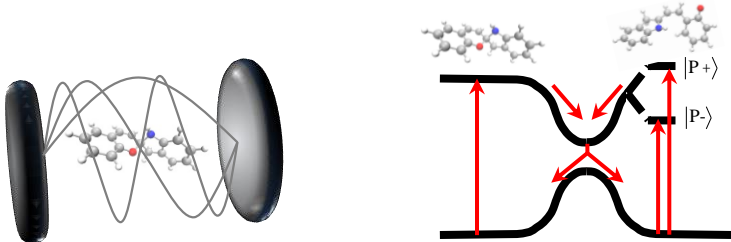
Magnetic tunnel junctions



Chirality-induced spin selectivity

Chemical Dynamics in Optical Cavities

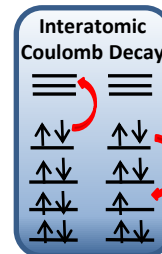
Investigate the influence of *quantum* light on chemical dynamics



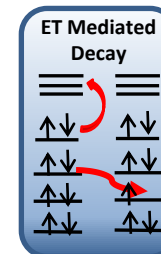
- Strong coupling between a molecular system and quantized modes of light can modify the *energy landscape*
- High-temperature Bose-Einstein condensates as a “quantum solvent” to modify chemistry based on charge-transfer

Electron Dynamics in Weakly Bound Systems

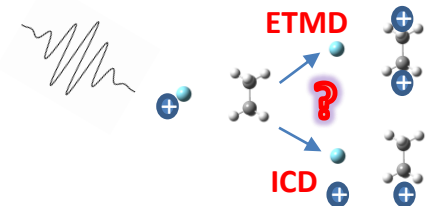
Ionization of a core electron initiates competing ultrafast electronic processes that lead to fragmentation in VdW and H-bonded systems



A B

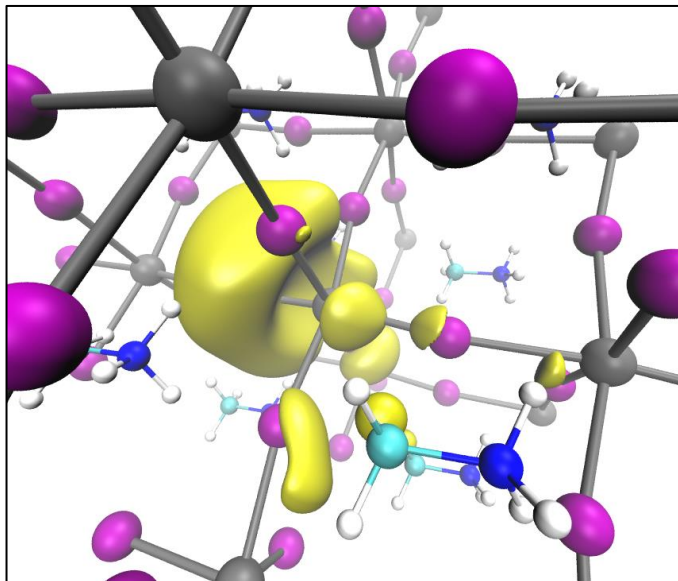


A B



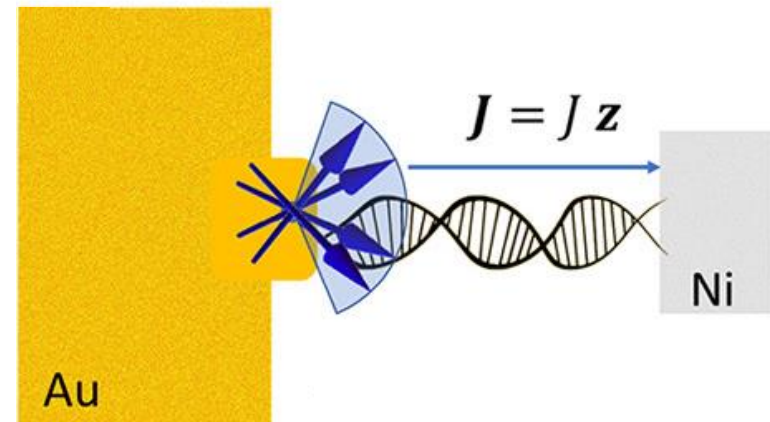
Non-Equilibrium Electron Dynamics

Photoexcited Materials

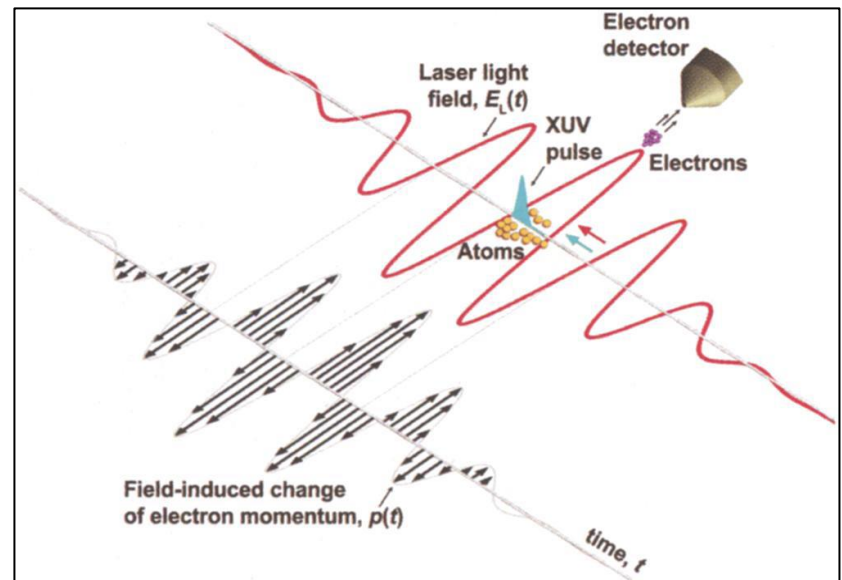


Electronic motion even in the absence of nuclear motion

Spin Dynamics and Selectivity

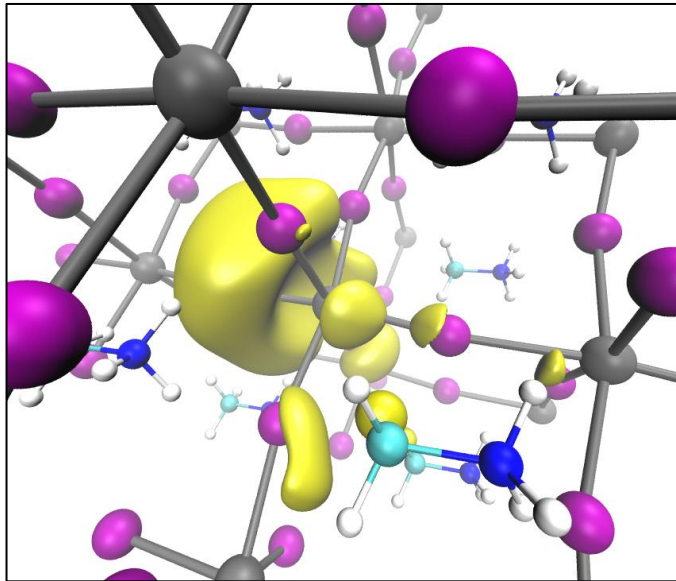


Laser Driven Electron Dynamics



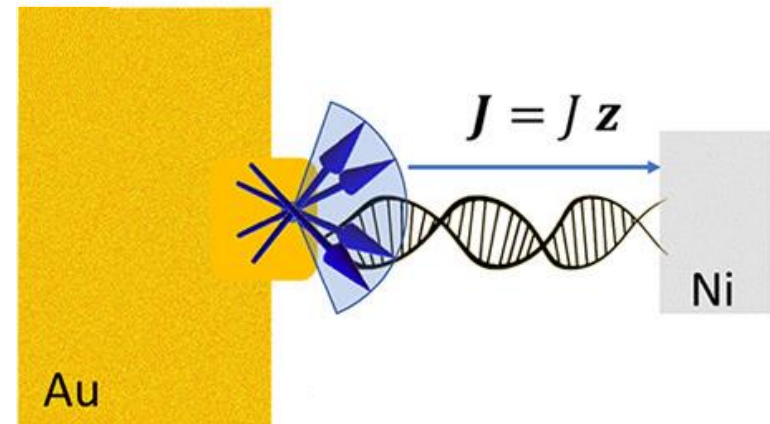
Non-Equilibrium Electron Dynamics

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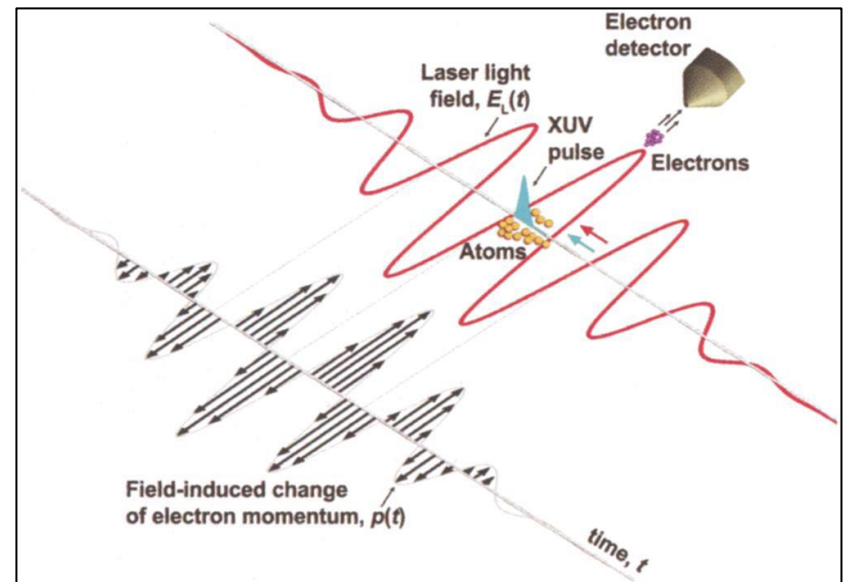


Develop efficient and accurate method for the direct simulation of the *electronic wavefunction* in strongly correlated systems

Spin Dynamics and Selectivity

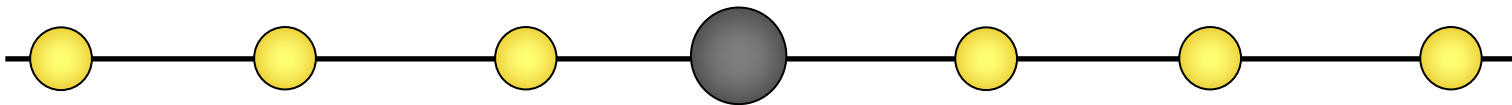
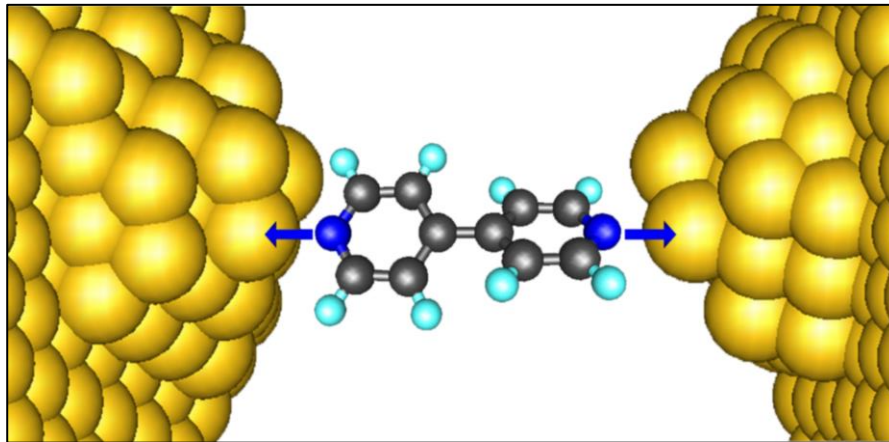


Laser Driven Electron Dynamics



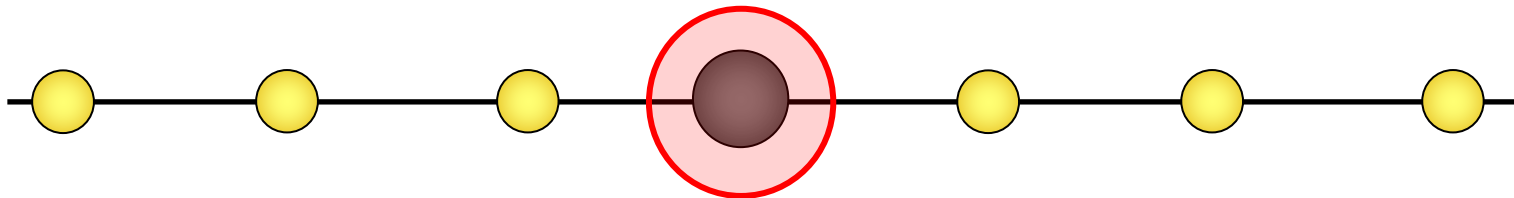
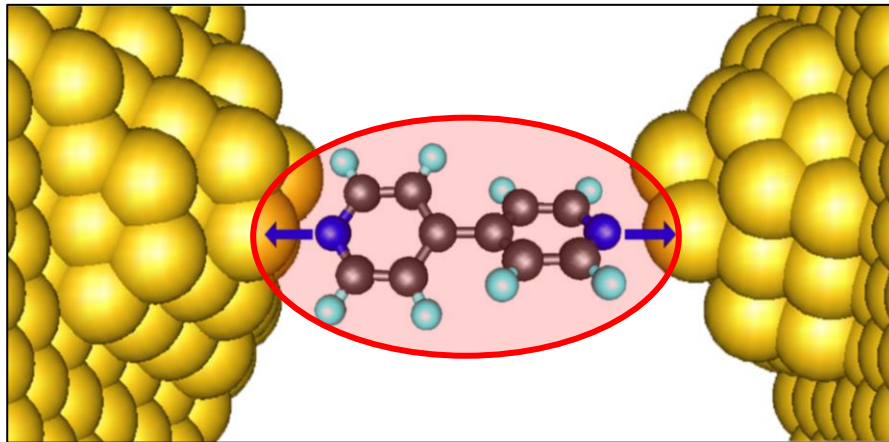
Quantum Embedding

Obtain accurate properties of a small subsystem without performing expensive calculation on the full system



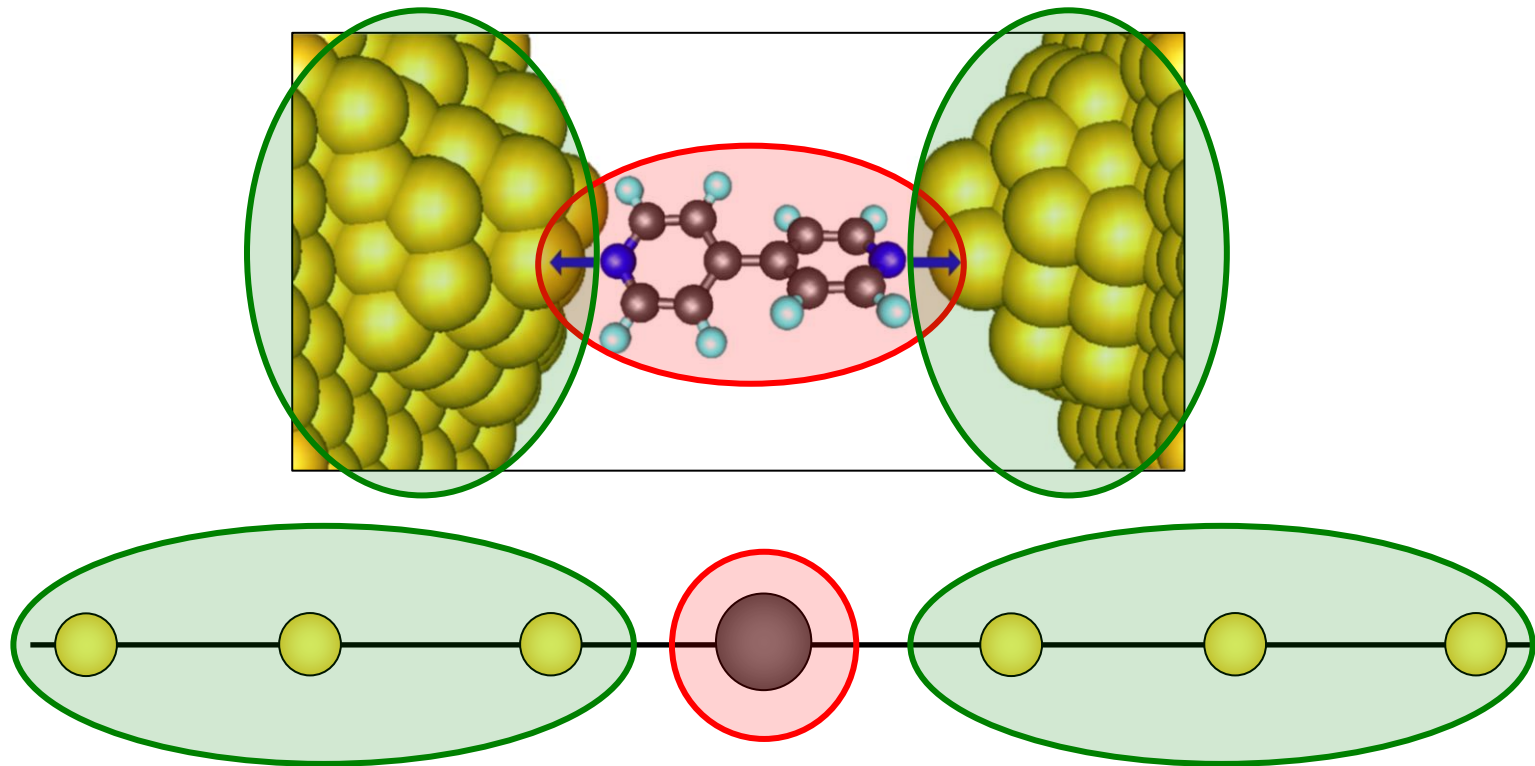
Quantum Embedding

Obtain accurate properties of a **small subsystem** without performing expensive calculation on the full system



Quantum Embedding

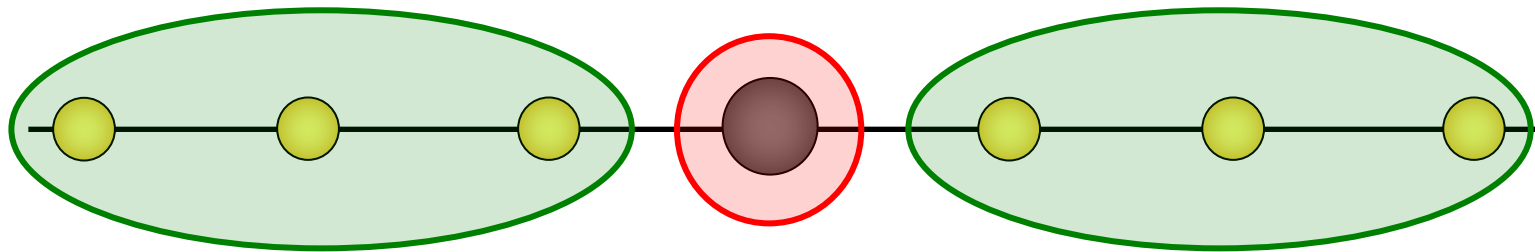
Obtain accurate properties of a **small subsystem** without performing expensive calculation on the full system



How to account for effects of the **surrounding environment**?

Density Matrix Embedding Theory

Surrounding **environment** treated as a quantum bath of the same size as the subsystem – Schmidt Decomposition



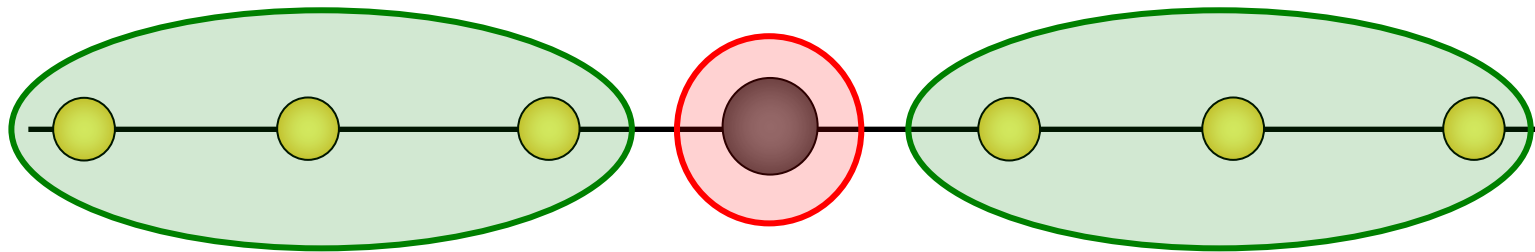
Schmidt decomposition:

$$|\Psi\rangle = \sum_i^{N_A} \sum_j^{N_B} \Psi_{ij} |A_i\rangle |B_j\rangle$$

Start with product of **subsystem** (A_i) and **environment** (B_j) states in **LARGE** space of dimension $N_a \times N_b$

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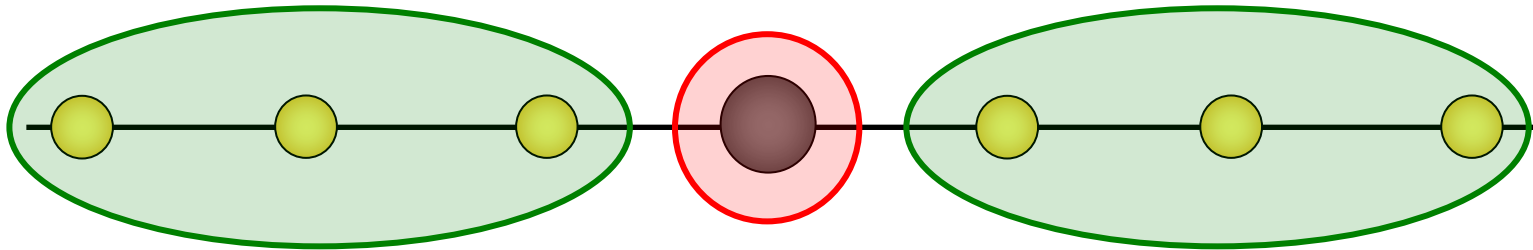
Schmidt decomposition:

$$|\Psi\rangle = \sum_i^{N_A} \sum_j^{N_B} \Psi_{ij} |A_i\rangle |B_j\rangle = \sum_i^{N_A} \sum_j^{N_B} \sum_{\alpha}^{\min(N_A, N_B)} U_{i\alpha} \lambda_{\alpha} V_{\alpha j}^{\dagger} |A_i\rangle |B_j\rangle$$

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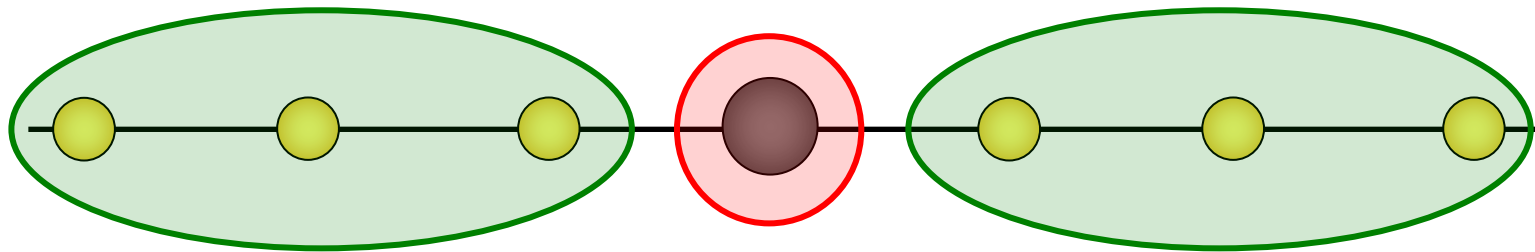
$$|\Psi\rangle = \sum_i^{N_A} \sum_j^{N_B} \Psi_{ij} |A_i\rangle |B_j\rangle = \sum_i^{N_A} \sum_j^{N_B} \sum_{\alpha}^{\min(N_A, N_B)} U_{i\alpha} \lambda_{\alpha} V_{\alpha j}^{\dagger} |A_i\rangle |B_j\rangle = \sum_{\alpha}^{\min(N_A, N_B)} \lambda_{\alpha} |\tilde{A}_{\alpha}\rangle |\tilde{B}_{\alpha}\rangle$$

Start with product of **subsystem** (A_i) and **environment** (B_j) states in **LARGE** space of dimension $N_a \times N_b$

End with product of **subsystem** (\tilde{A}_{α}) and **new environment** (\tilde{B}_{α}) states in **SMALL** space of dimension $N_a \times N_a$

Density Matrix Embedding Theory

Surrounding **environment** treated as a **quantum bath** of the same size as the subsystem – Schmidt Decomposition



Schmidt decomposition:

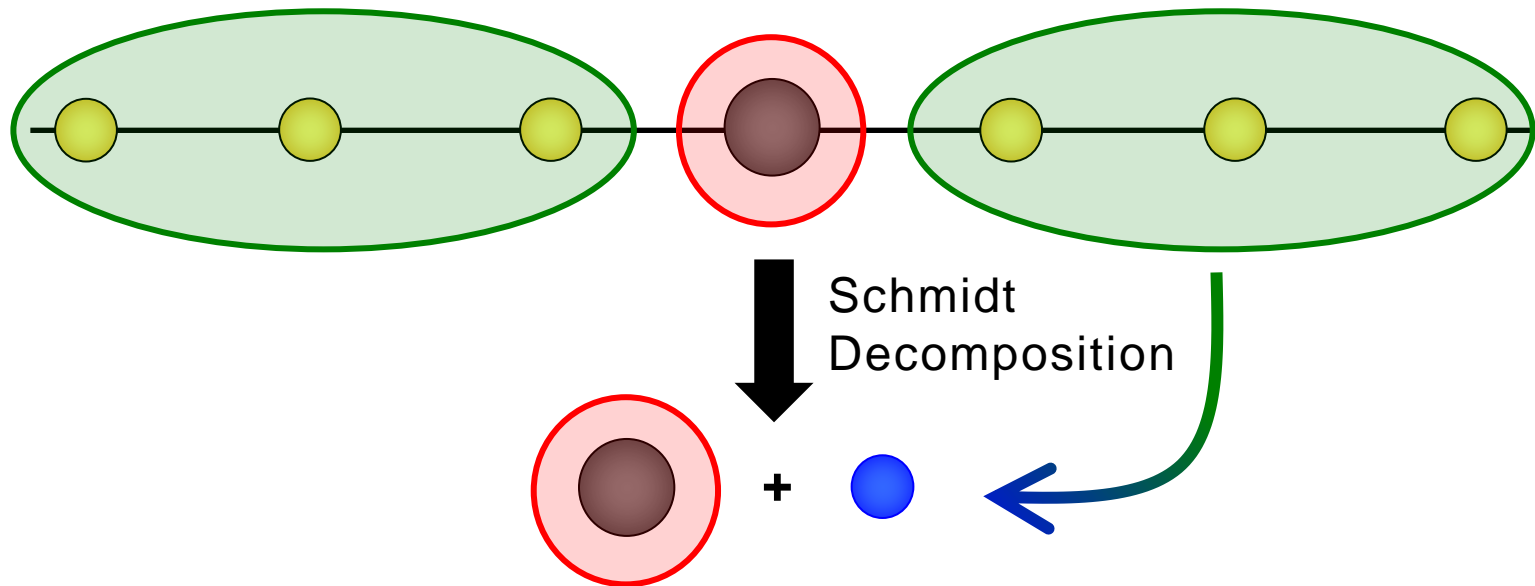
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Start with product of **subsystem** (A_i) and **environment** (B_j) states in **LARGE** space of dimension $N_a \times N_b$

End with product of **subsystem** (\tilde{A}_{α}) and **multi-electron bath** (\tilde{B}_{α}) states in **SMALL** space of dimension $N_a \times N_a$

Density Matrix Embedding Theory

Surrounding **environment** treated as a **quantum bath** of the same size as the subsystem – Schmidt Decomposition

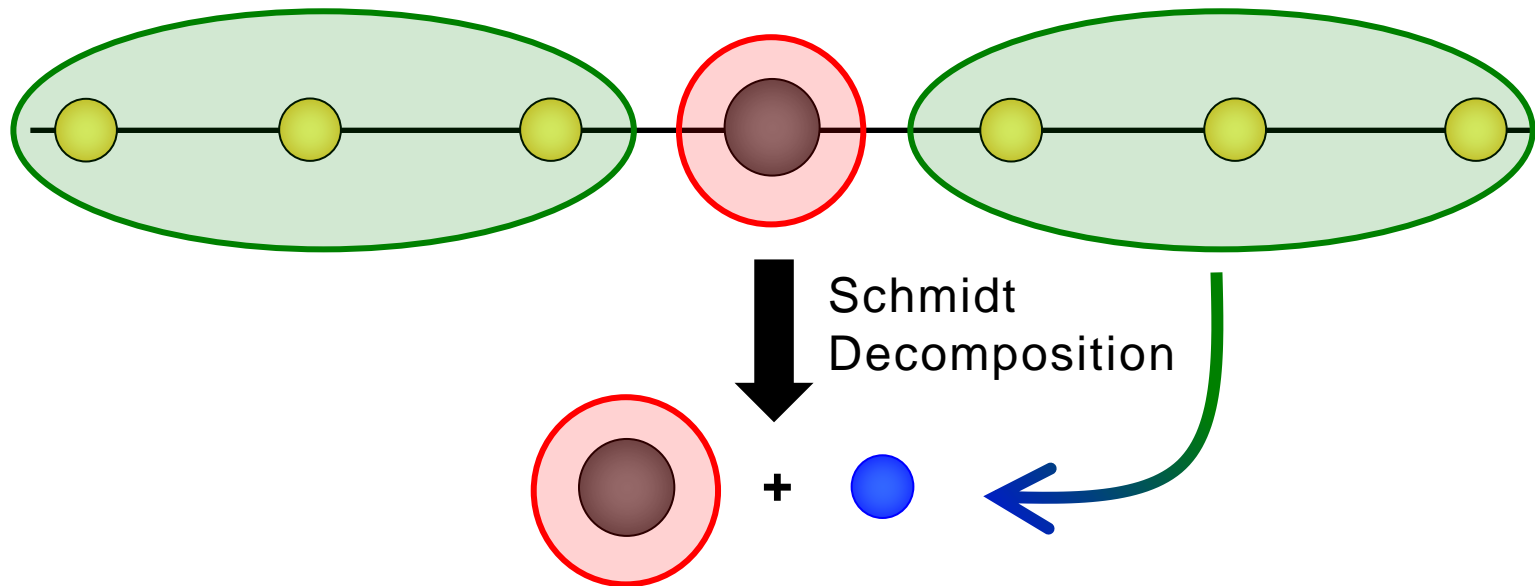


Start with product of **subsystem** (A_i) and **environment** (B_j) states in **LARGE** space of dimension $N_a \times N_b$

End with product of **subsystem** (\tilde{A}_α) and **multi-electron bath** (\tilde{B}_α) states in **SMALL** space of dimension $N_a \times N_a$

Density Matrix Embedding Theory

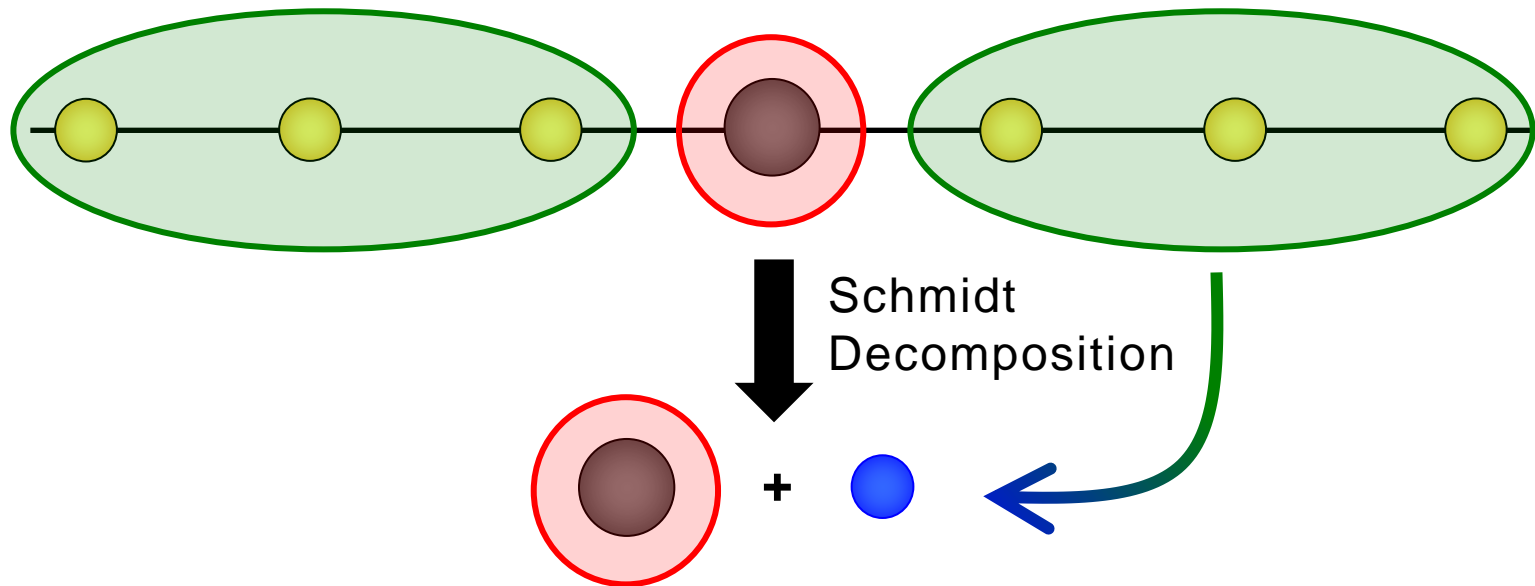
Surrounding **environment** treated as a **quantum bath** of the same size as the subsystem – Schmidt Decomposition



The good: in principle have a method to exactly embed a fragment in a bath of the same size as the fragment itself

Density Matrix Embedding Theory

Surrounding **environment** treated as a **quantum bath** of the same size as the subsystem – Schmidt Decomposition

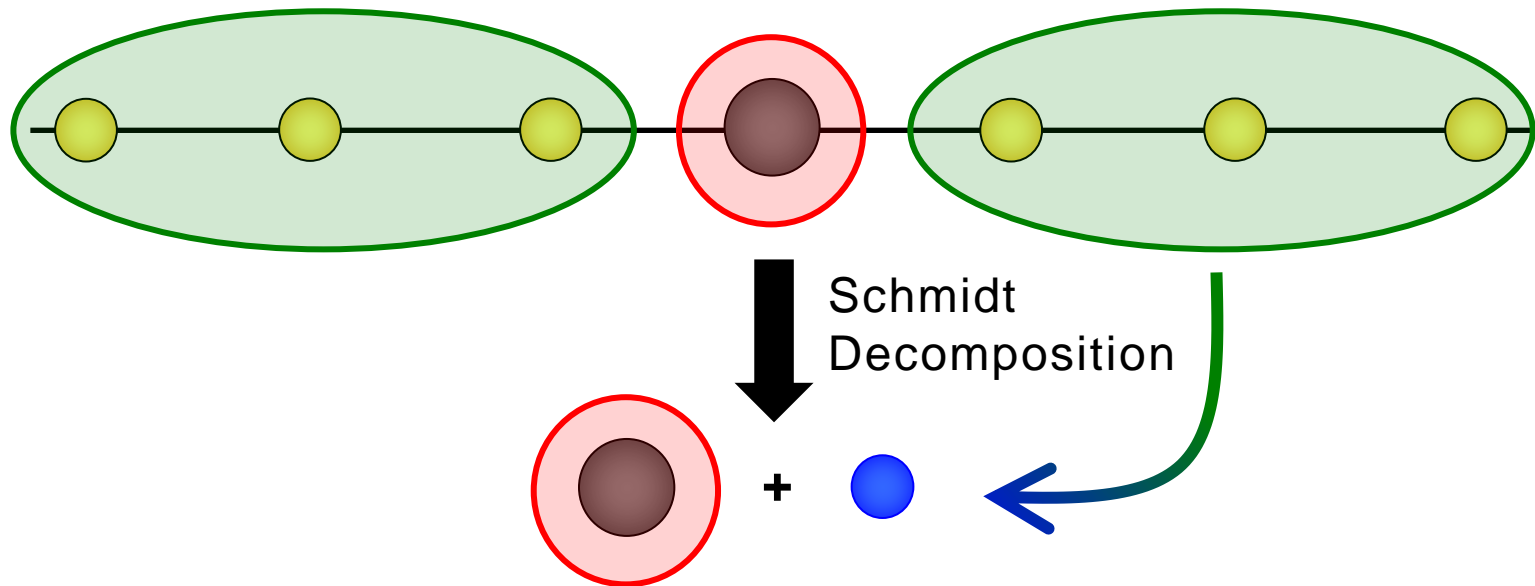


The good: in principle have a method to exactly embed a fragment in a bath of the same size as the fragment itself

The bad: knowledge of the exact wavefunction is needed to construct the bath states

Density Matrix Embedding Theory

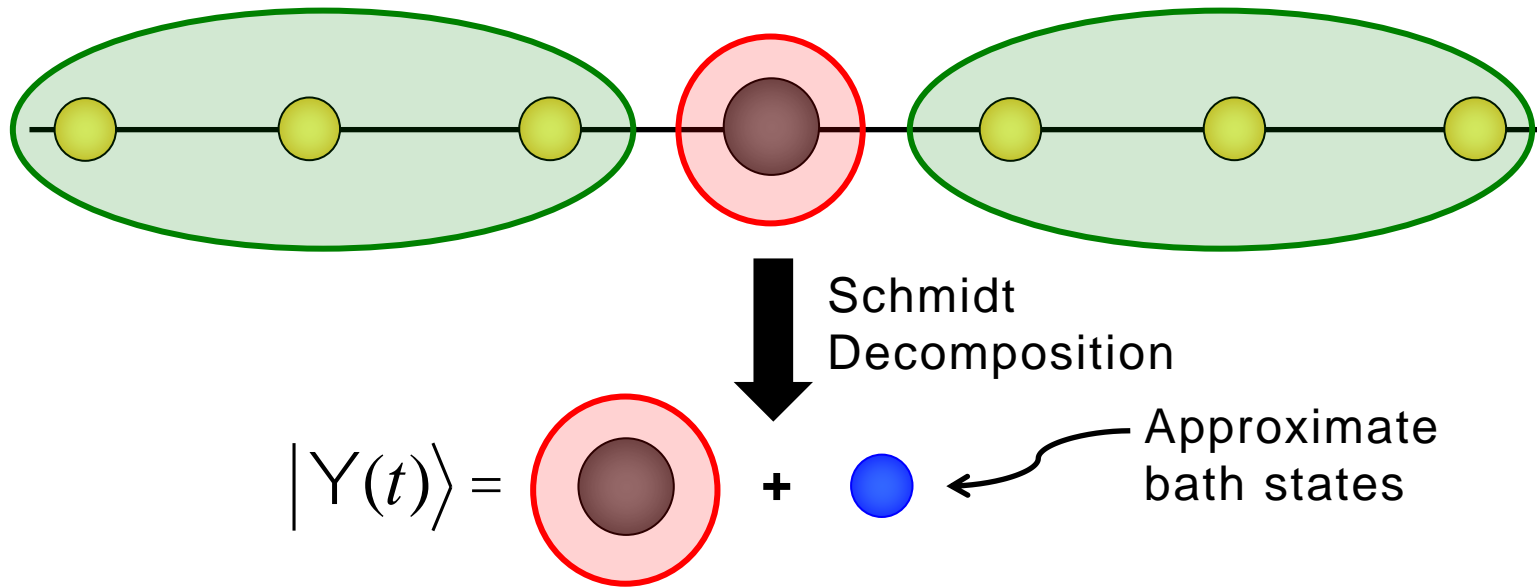
Surrounding **environment** treated as a **quantum bath** of the same size as the subsystem – Schmidt Decomposition



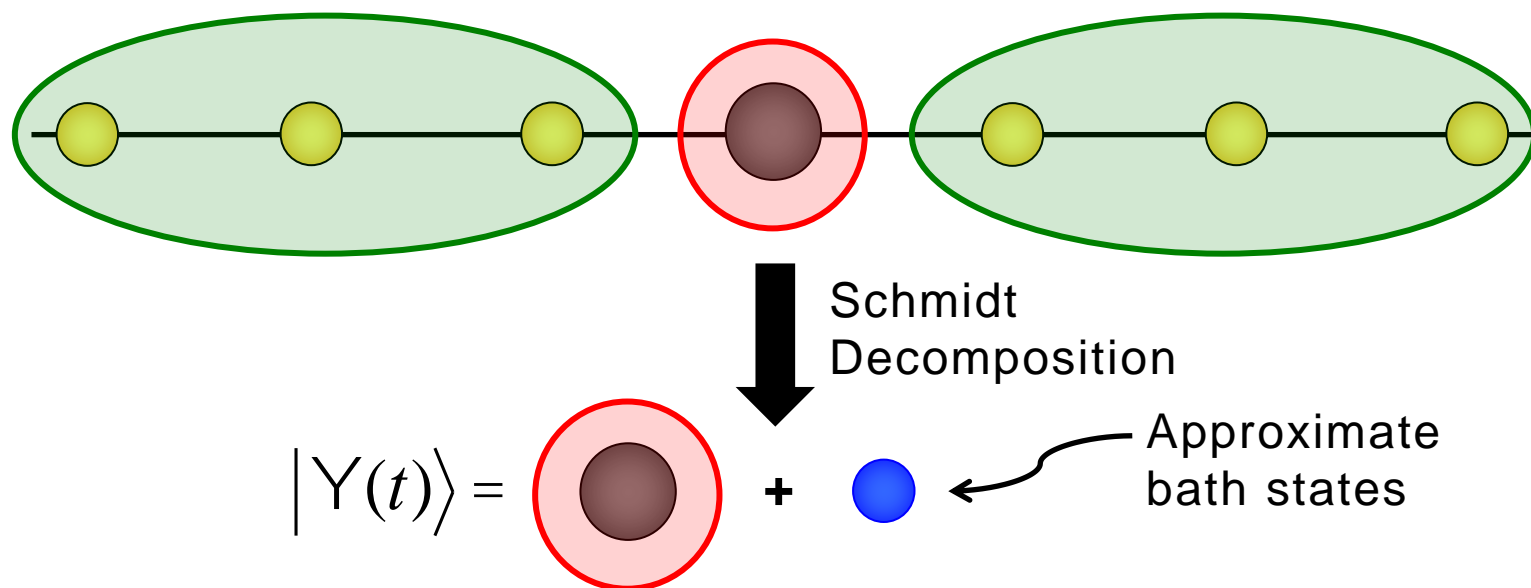
The Fix – Static DMET:

1. Construct the bath states from *mean-field* wavefunction on the total system
2. Perform *correlated* calculation utilizing these bath states
3. Introduce self-consistency between mean-field wavefunction and correlated calculation – one-electron reduced density matrix

Real-Time DMET

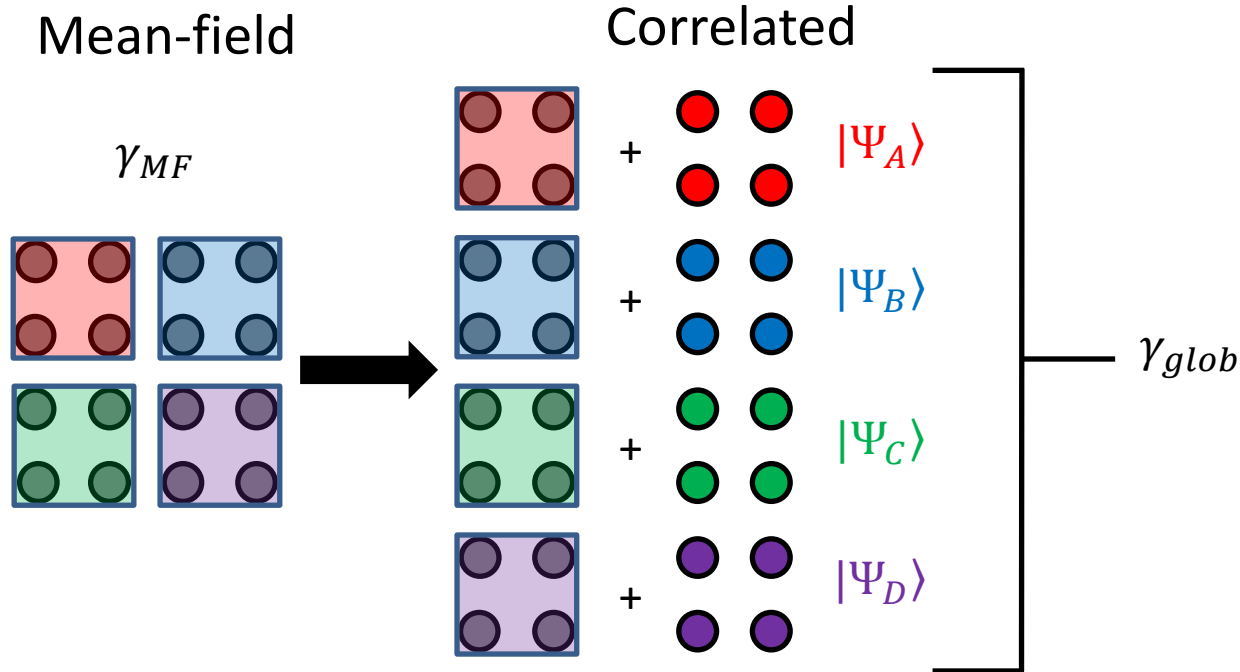


Real-Time DMET



- Initial formulation utilized time-dependent variational principle to derive equations of motion for embedding wavefunction
- Equations of motion bare resemblance to time-dependent CAS

Multi-Fragment Extension

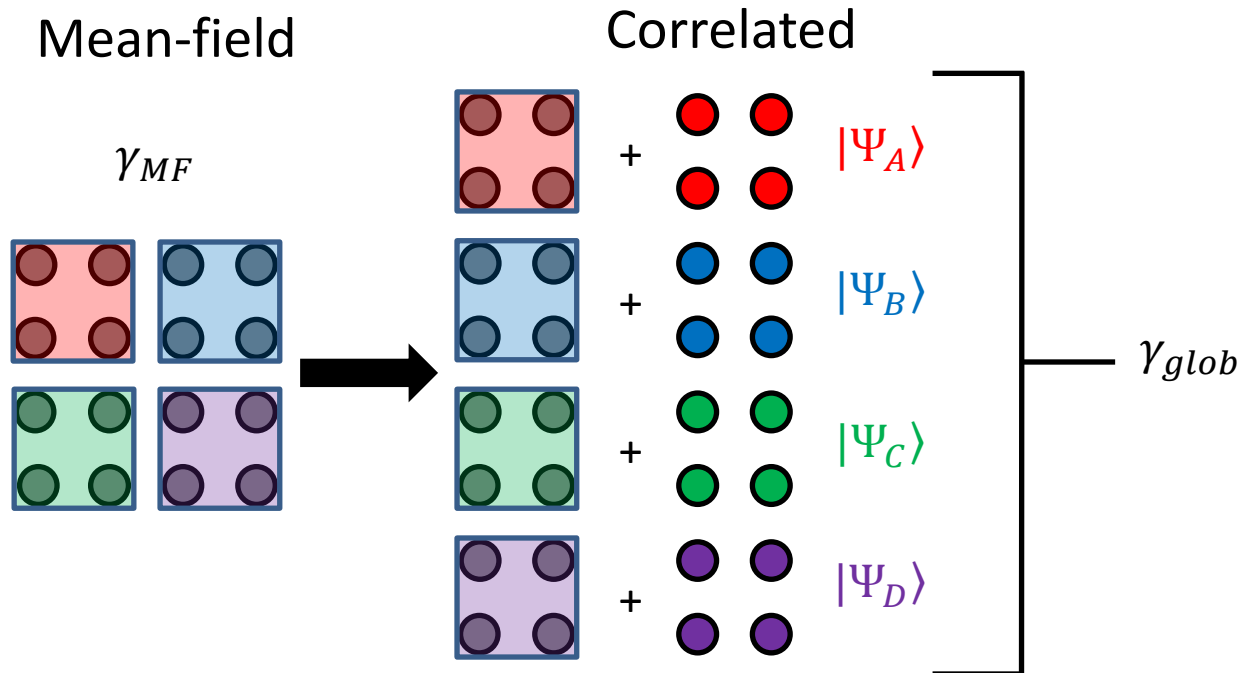


Separate correlated embedding calculations all derived from same total-system mean-field calculation

“Stitch” together embedding calculations through democratic partitioning of operators

$$\langle \hat{a}_i^\dagger \hat{a}_j + \hat{a}_j^\dagger \hat{a}_i \rangle = \langle \Psi_{x(i)} | \hat{a}_i^\dagger \hat{a}_j | \Psi_{x(i)} \rangle + \langle \Psi_{x(j)} | \hat{a}_j^\dagger \hat{a}_i | \Psi_{x(j)} \rangle$$

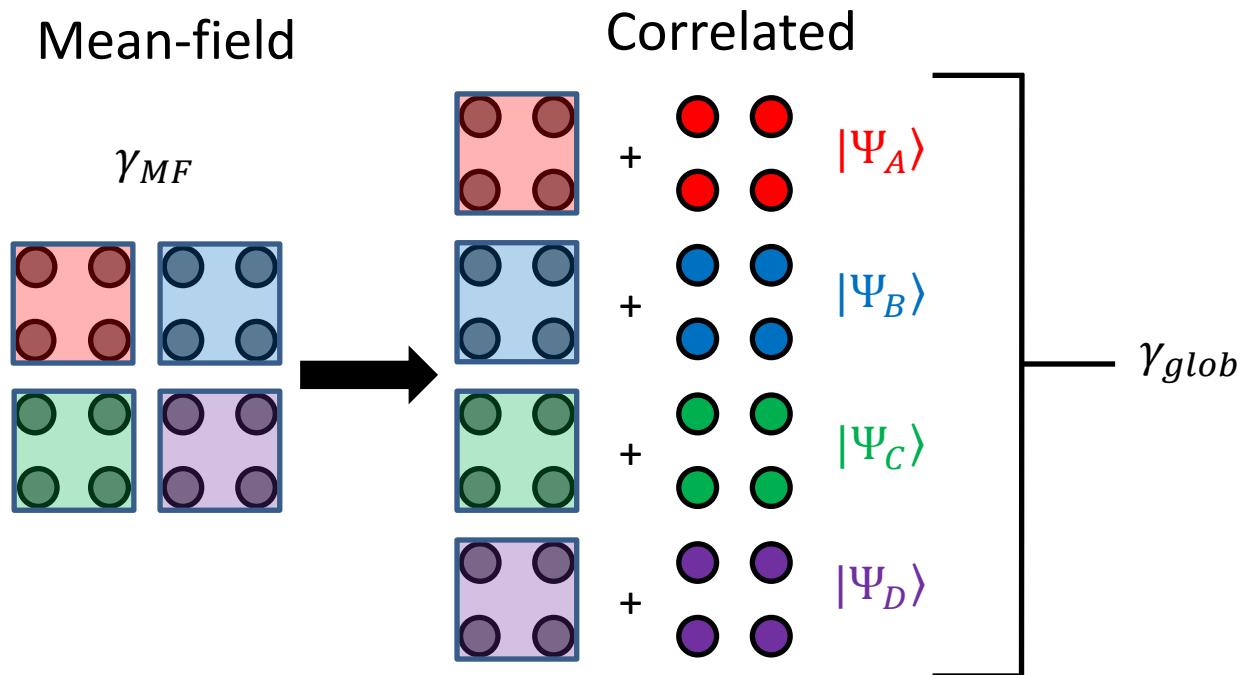
Multi-Fragment Extension



Difficulty for Dynamics:

- Conventional DMET matching condition between mean-field and embedding calculations breaks down
- TDVP does not allow for multiple fragments

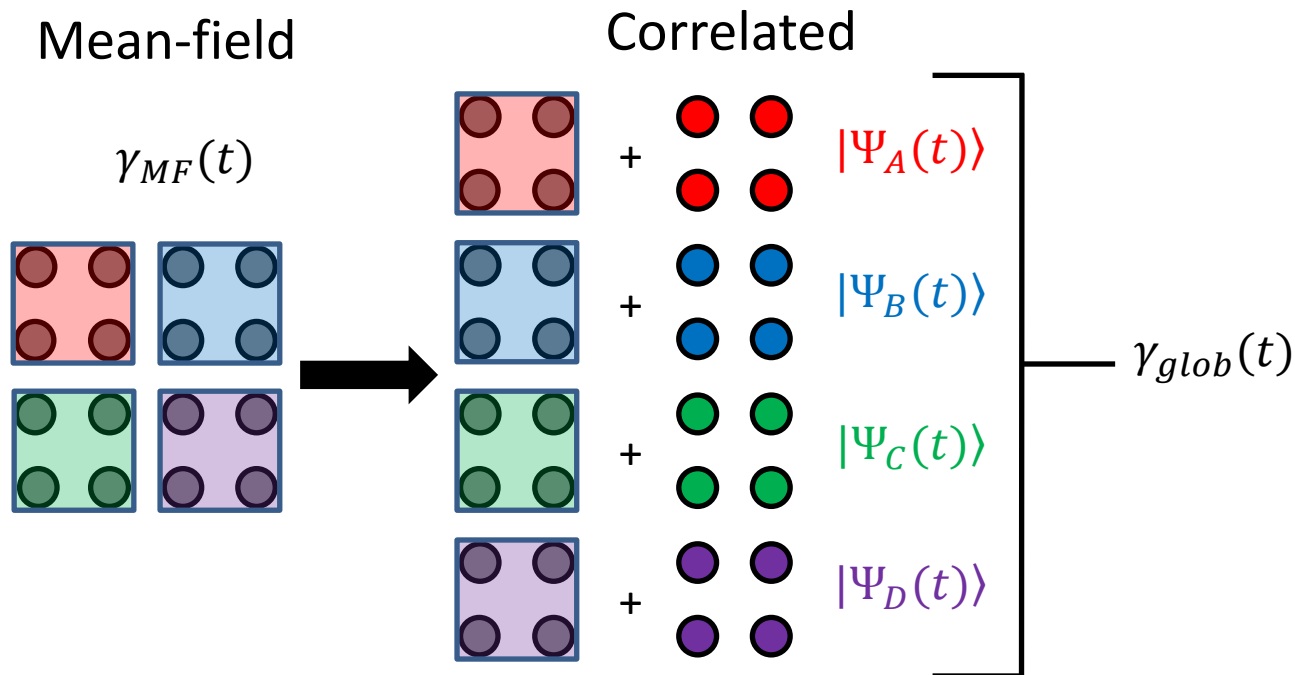
Multi-Fragment Extension



Solution:

- Take advantage of a different form of static DMET – projected DMET
- pDMET has an analytical matching condition between γ_{MF} and γ_{glob}

Real-Time pDMET



General Idea: Simultaneously propagate mean-field wavefunction of full system and correlated wavefunction for each fragment and its bath, such that the pDMET conditions are met at each point in time

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Condition 1: Bath for each fragment obtained from Schmidt decomposition of mean-field 1RDM

$$\gamma_{\text{env}}^{\text{MF}}(t)|a^{\text{A}}(t)\rangle = \lambda_a^{\text{A}}|a^{\text{A}}(t)\rangle$$

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Condition 2: Mean-field 1RDM obtained from natural orbitals of global 1RDM formed from all fragments

$$\begin{aligned}\gamma^{\text{MF}}(t) &= 2C(t)C(t)^\dagger \\ \gamma^{\text{glob}}(t) &= C(t)\Lambda(t)C^\dagger(t)\end{aligned}$$

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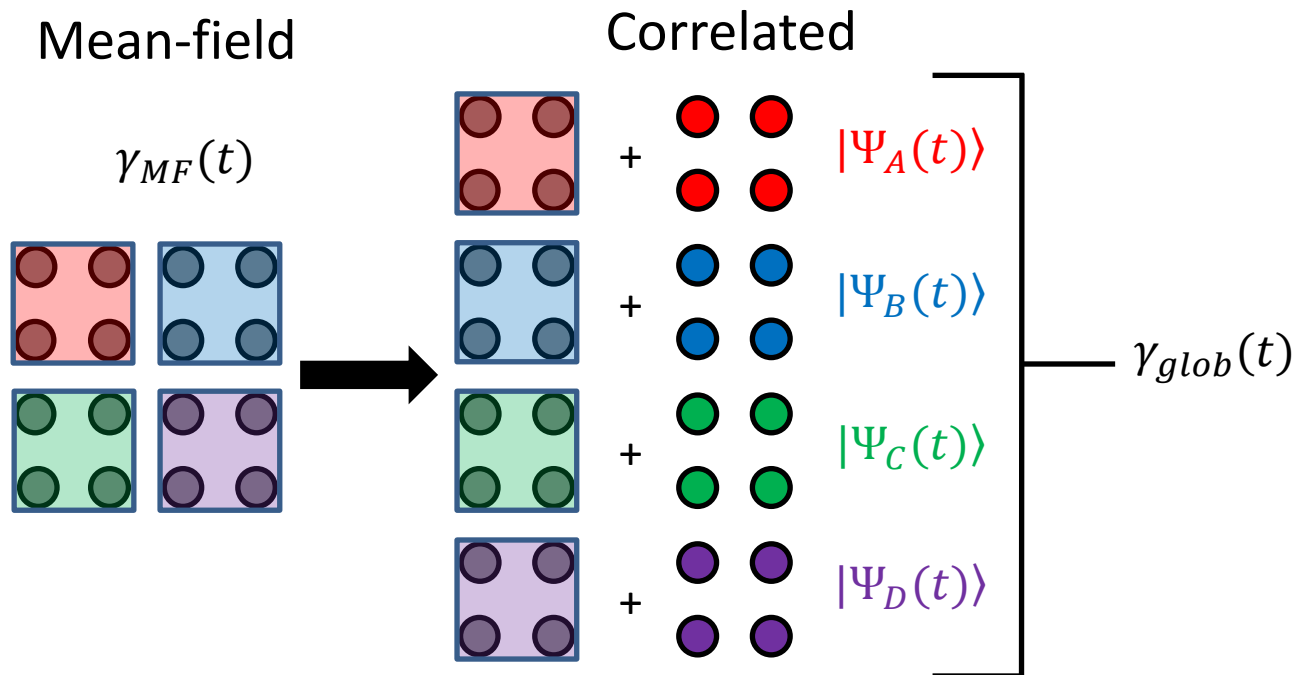
$$\begin{aligned}\gamma^{\text{MF}}(t) &= 2C(t)C(t)^\dagger \\ \gamma^{\text{glob}}(t) &= C(t)\Lambda(t)C^\dagger(t)\end{aligned}$$

Time-dependence of DMET wavefunction for each fragment:

$$i|\dot{\Psi}^{\text{A}}\rangle = \sum_m i\dot{c}_m^{\text{A}}|m^{\text{A}}\rangle + c_m^{\text{A}}X^{\text{A}}|m^{\text{A}}\rangle$$

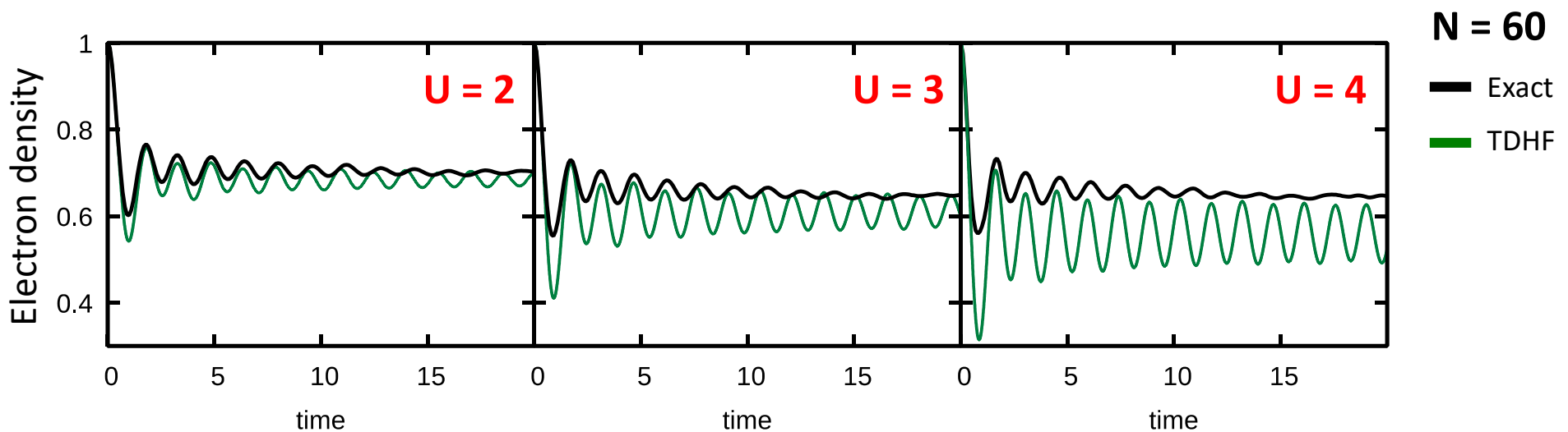
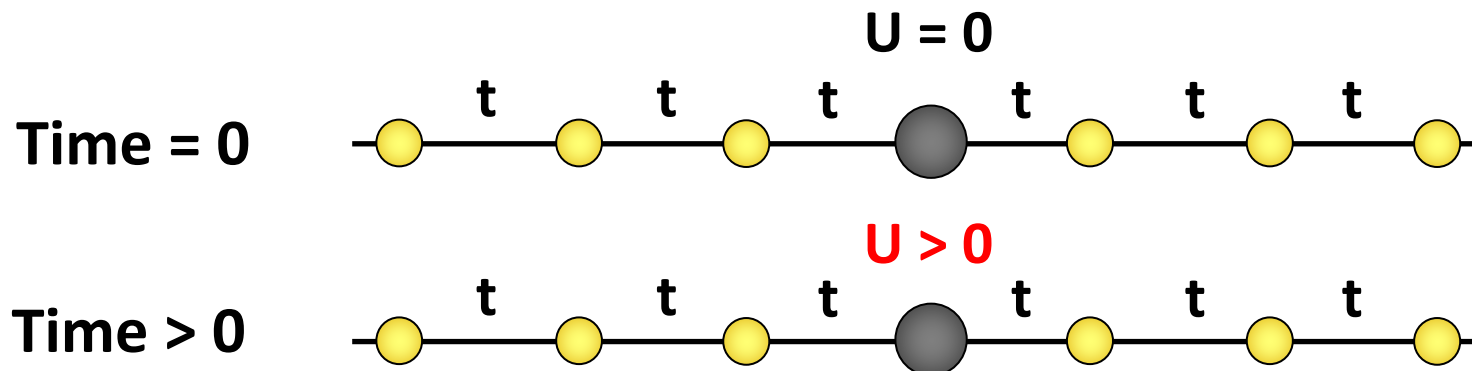
\Rightarrow \dot{c}_m^{A} obtained from TDSE projected into the embedding space
 X^{A} obtained from pDMET conditions at each point in time

Real-Time pDMET

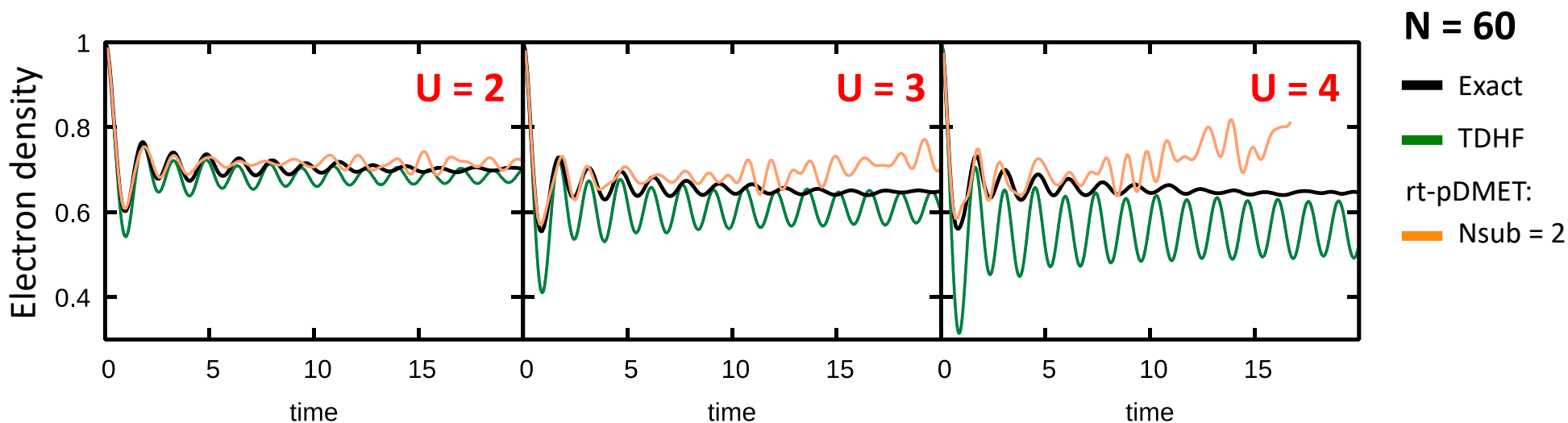
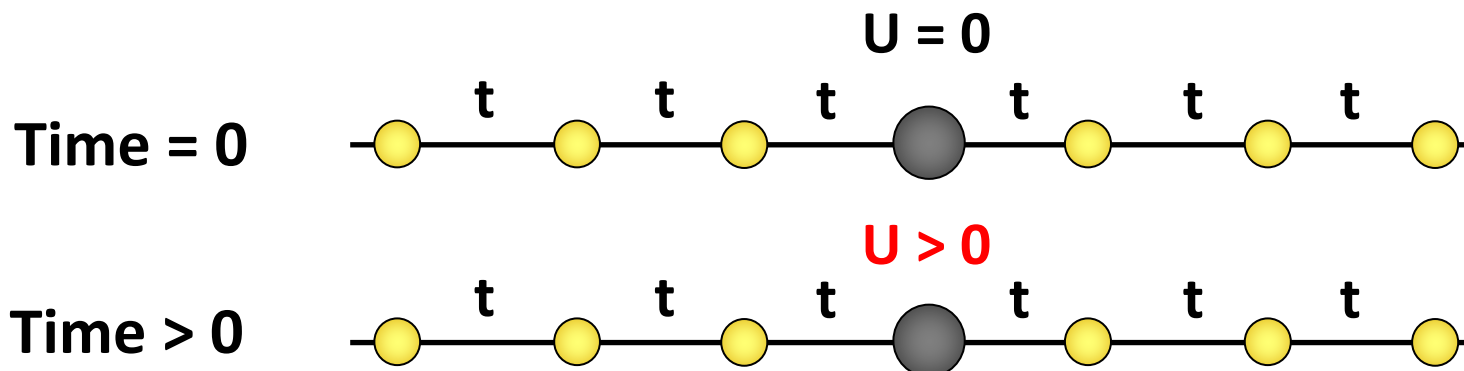


- Exact in non-interacting limit
- Exact in large subsystem limit
- Allows for correlation between subsystem and environment
- Do not need to specify specific area of high accuracy

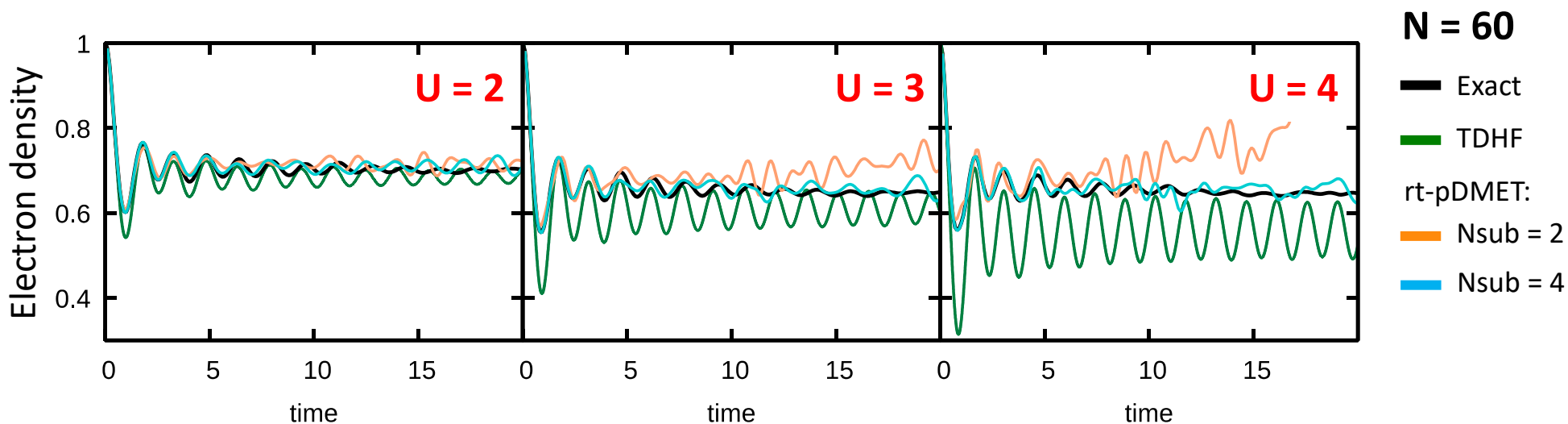
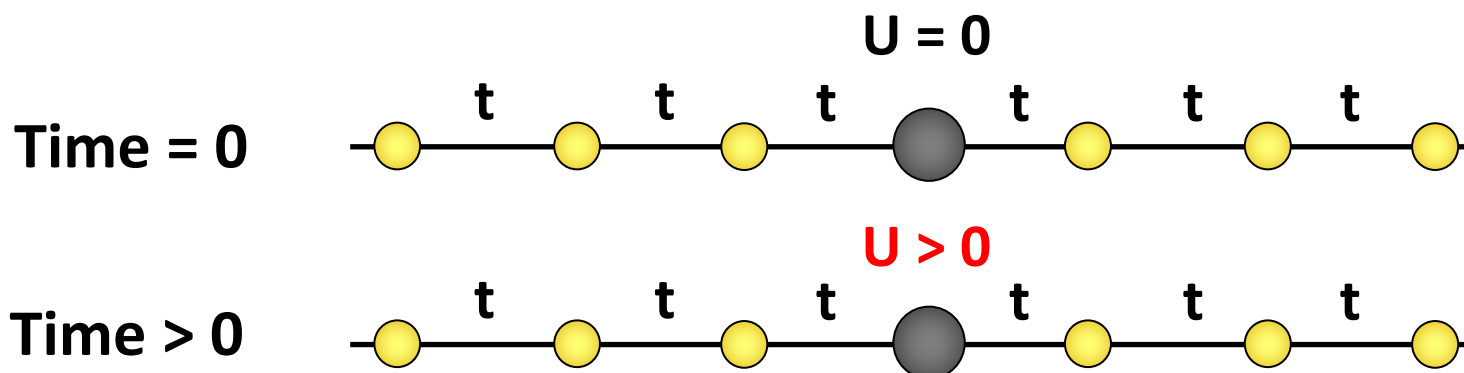
Single Impurity Anderson Model



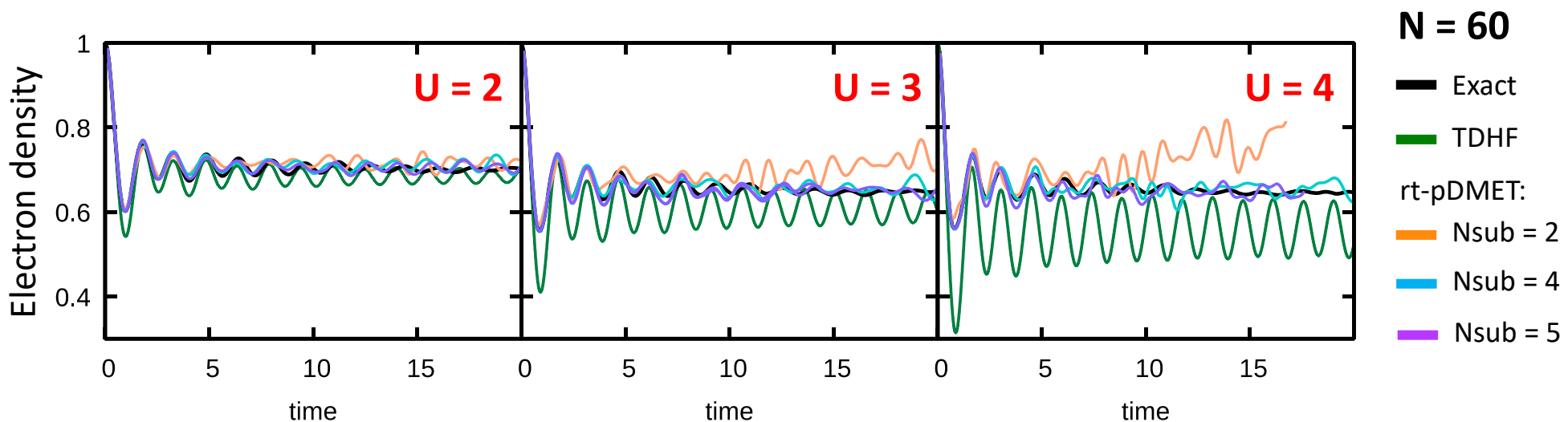
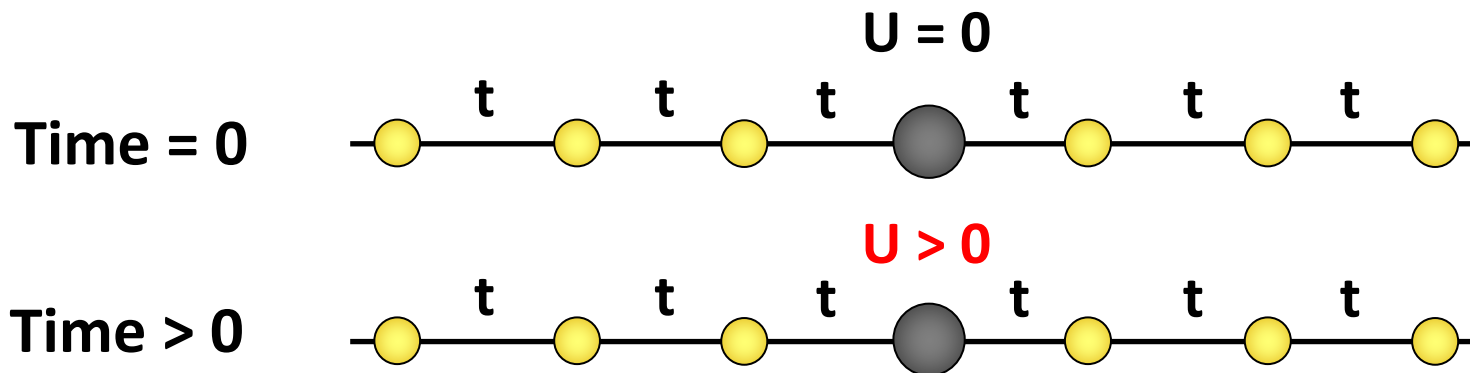
Single Impurity Anderson Model



Single Impurity Anderson Model

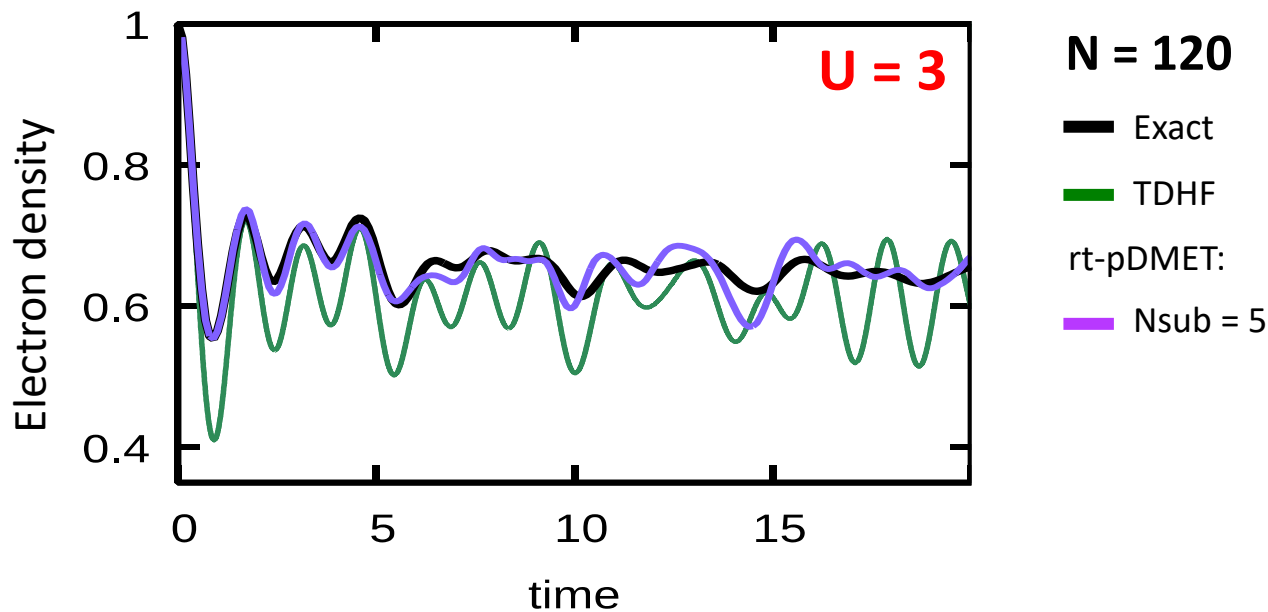
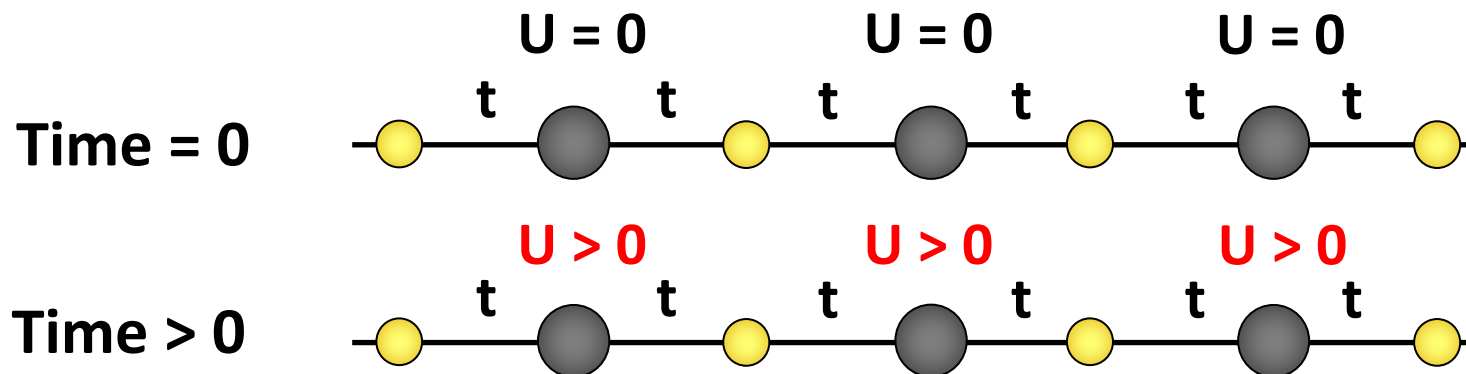


Single Impurity Anderson Model

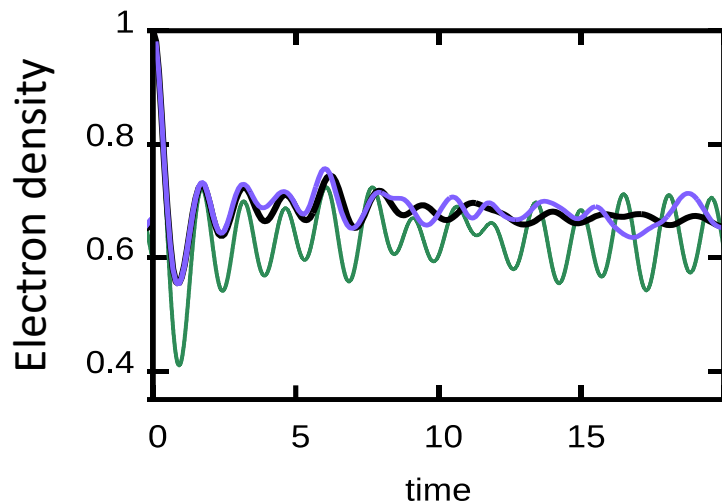
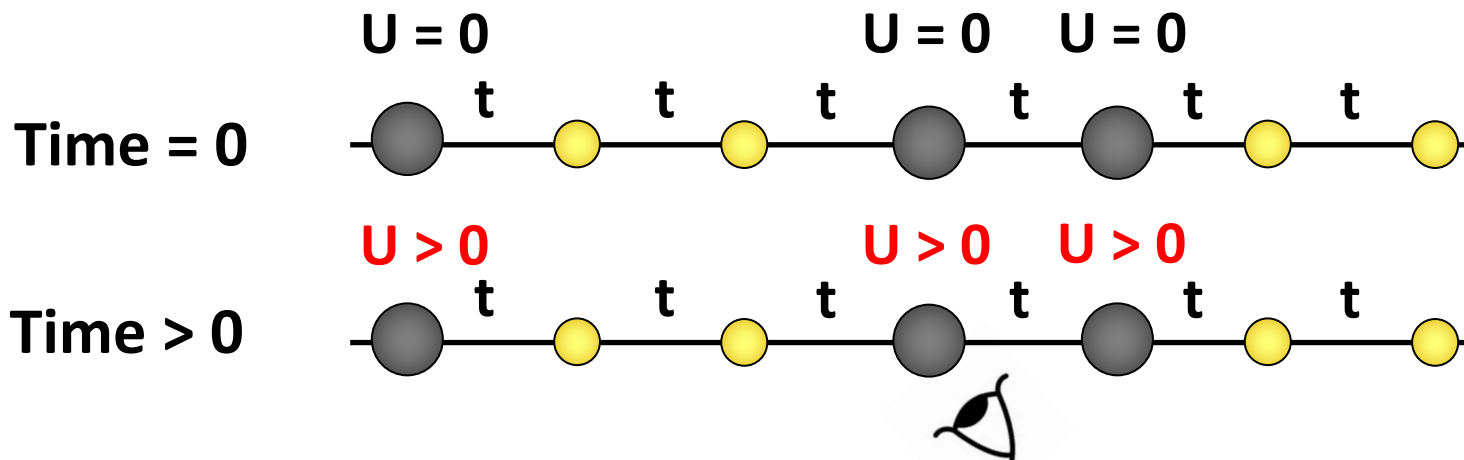


Real-time pDMET converges to the exact answer even in the strongly correlated regime

Multi-Impurity Anderson Model



Multi-Impurity Anderson Model



N = 120

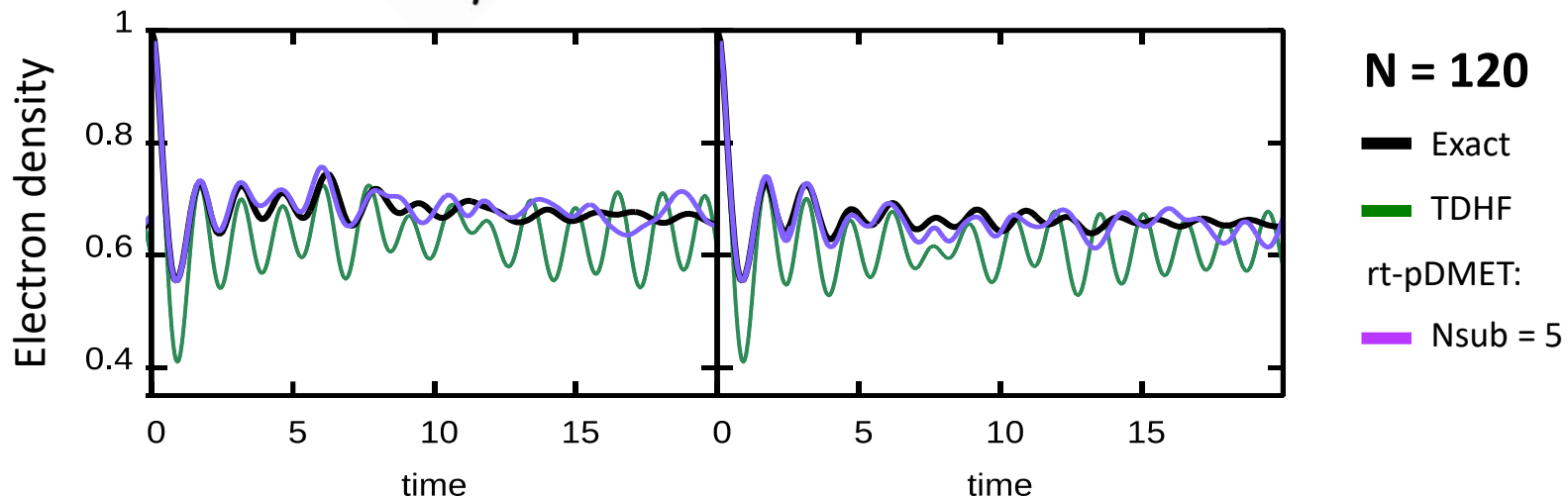
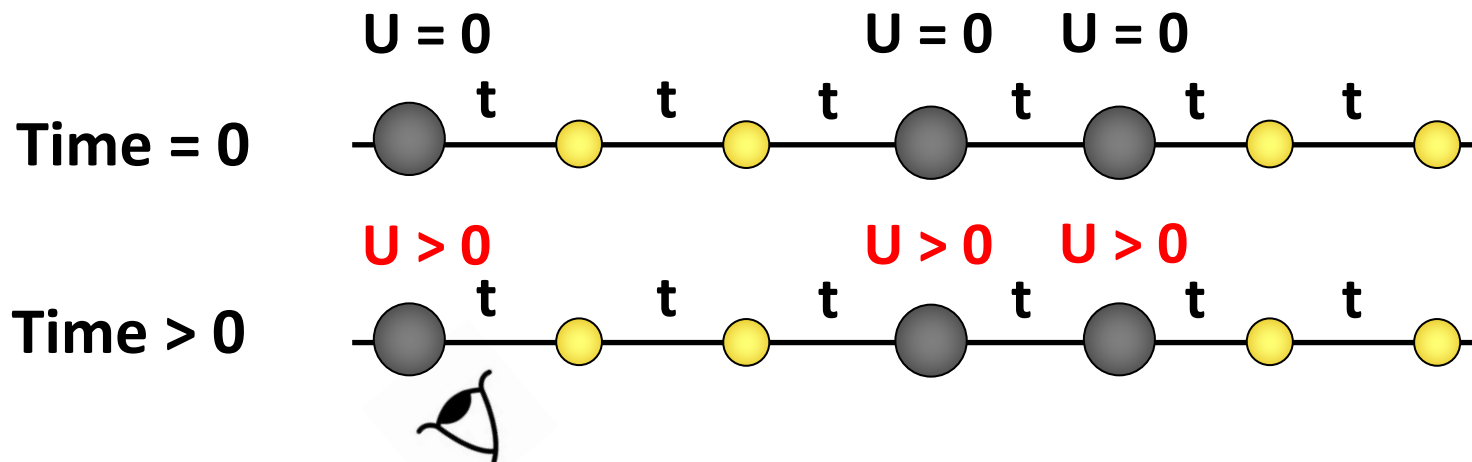
— Exact

— TDHF

rt-pDMET:

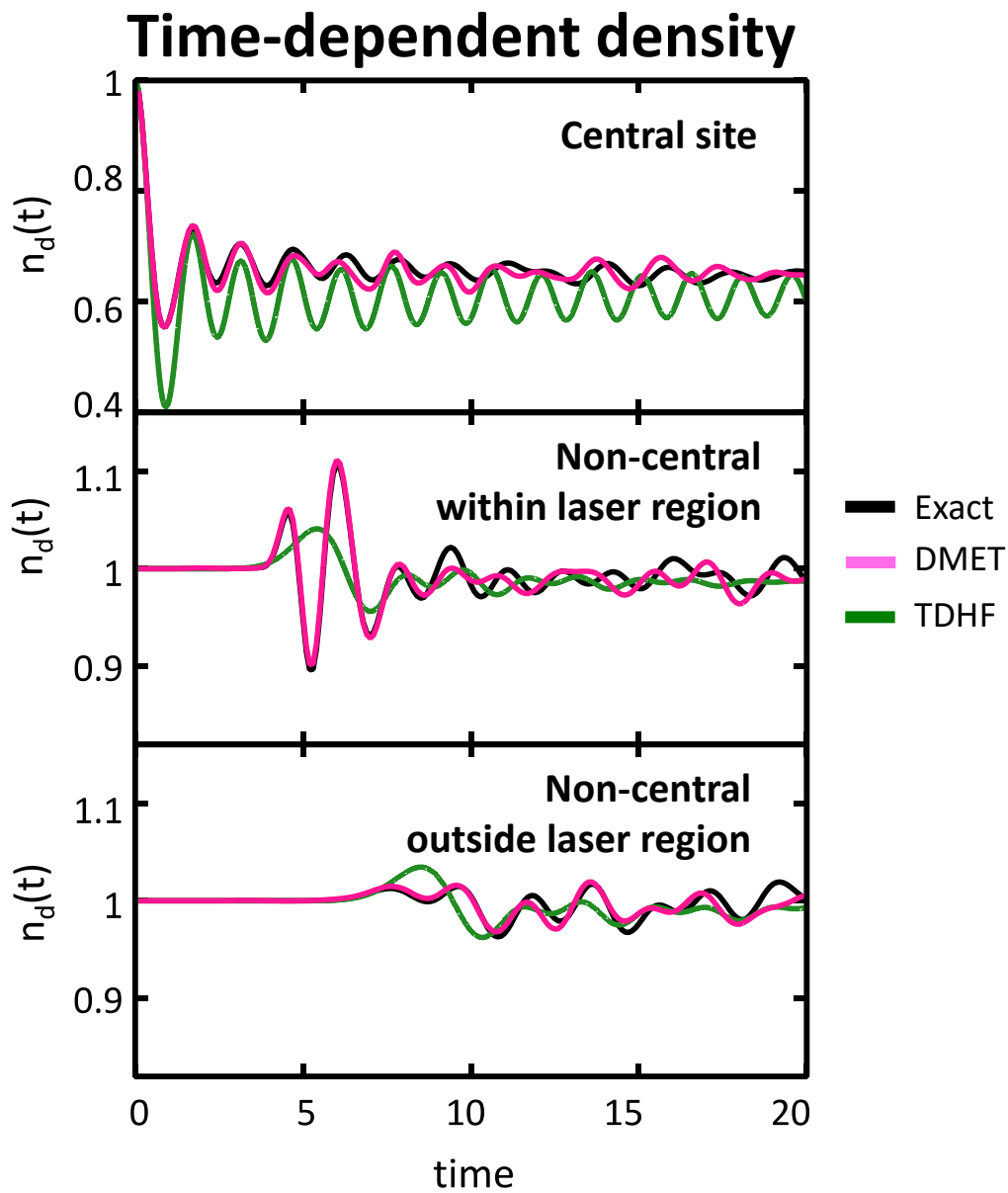
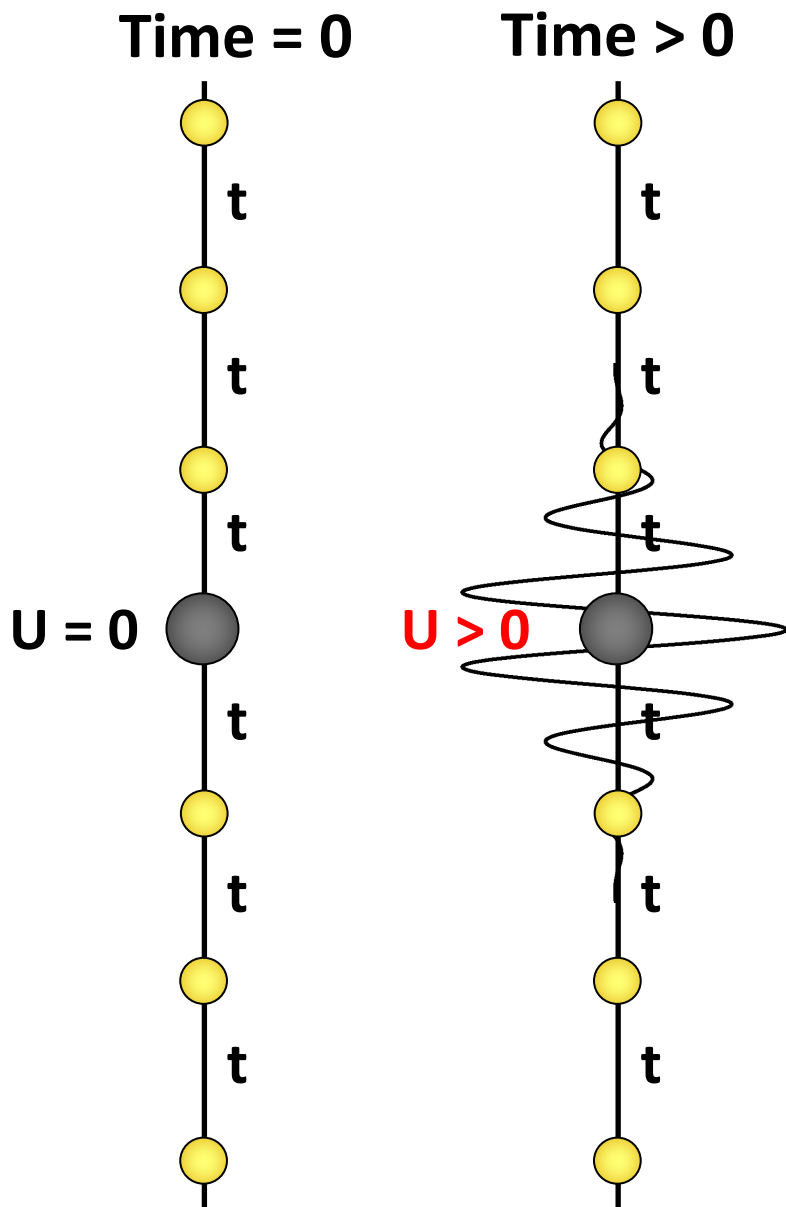
— Nsub = 5

Multi-Impurity Anderson Model



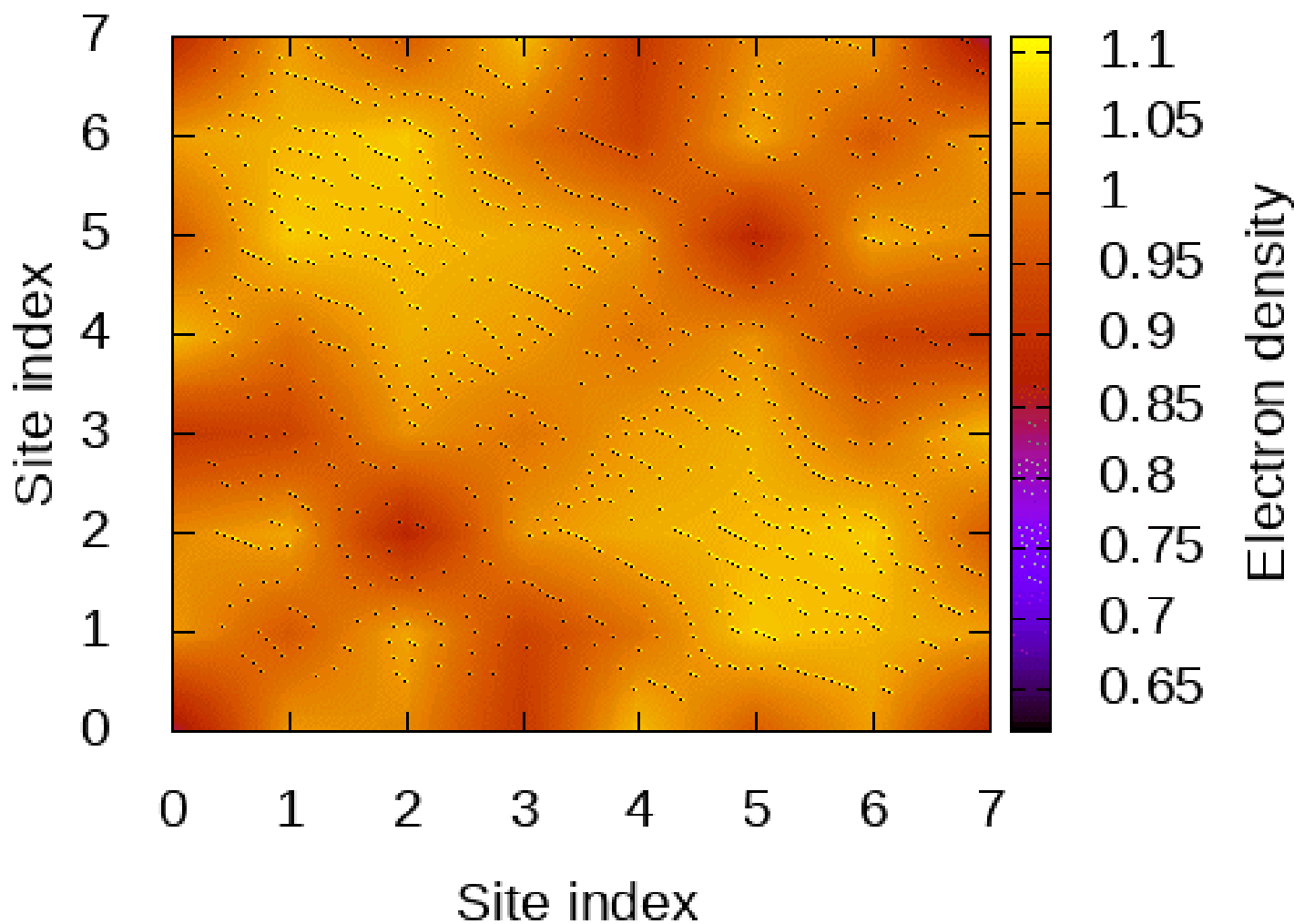
Real-time pDMET accurately treats all regions of space even for a disordered system

Laser-Driven Electron Dynamics

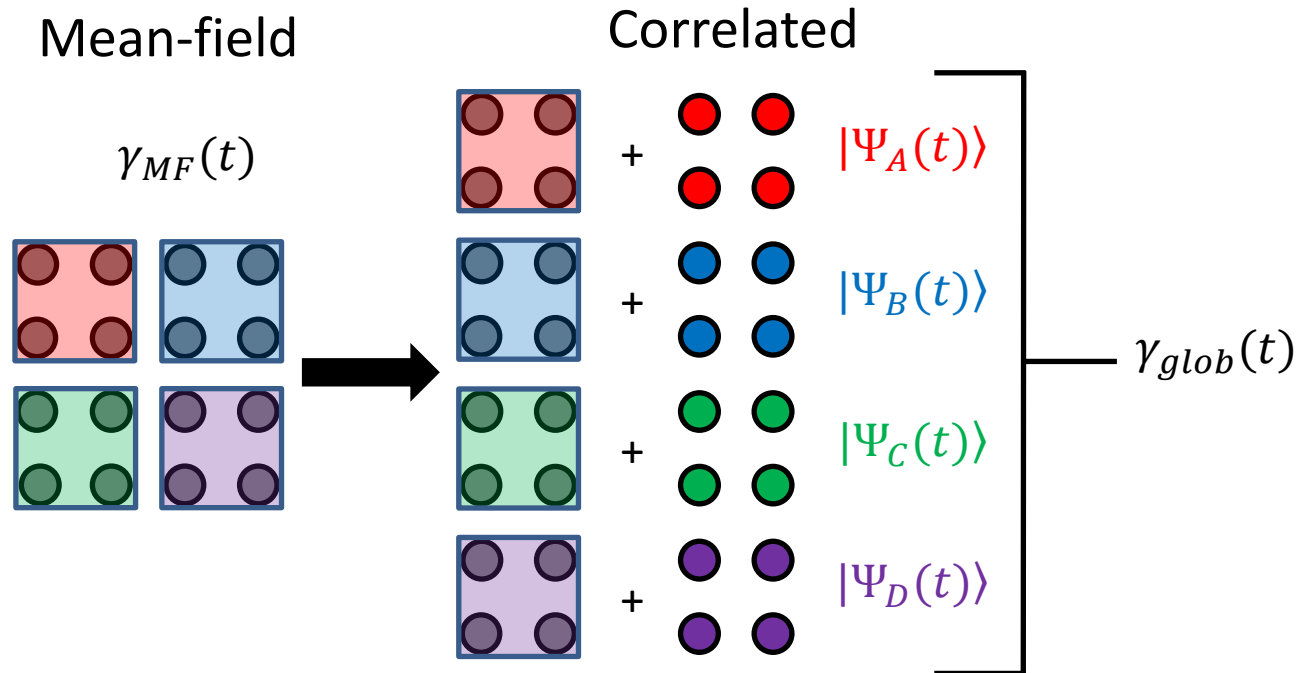


Multi-dimensional Impurity Model

Electron density change in 2D SIAM



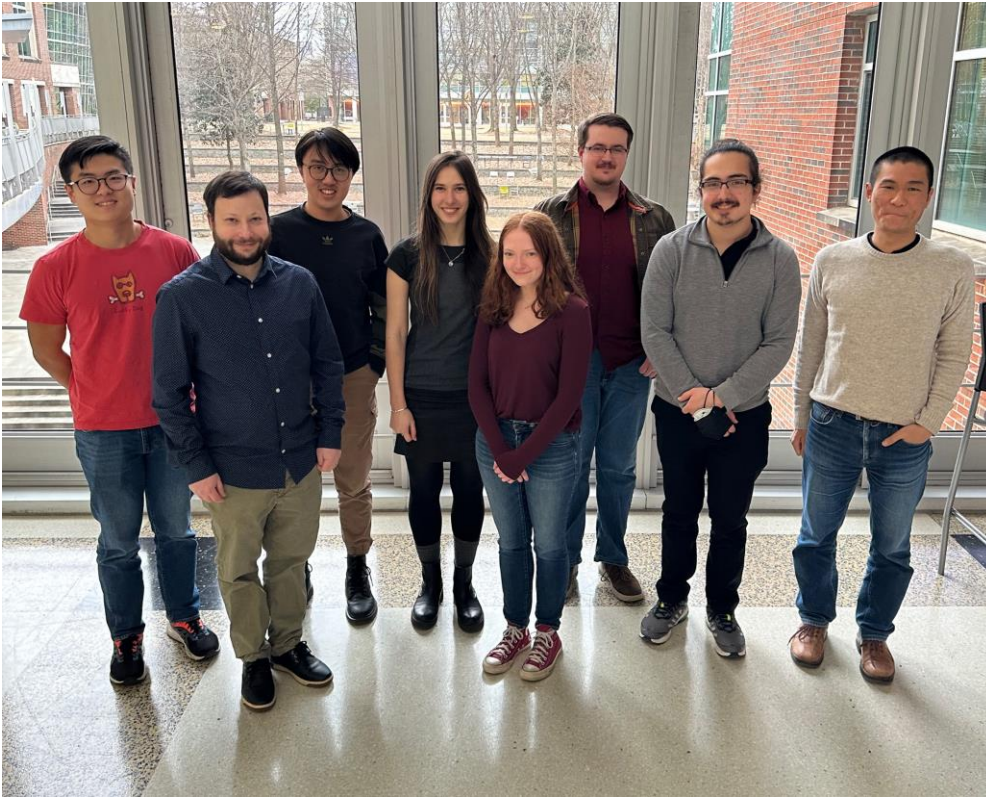
Conclusions



Real-time projected density matrix embedding theory

- Multi-fragment real-time quantum embedding method
- Non-equilibrium electron dynamics in strongly correlated systems
- Accurately treat dynamics across entire system

Acknowledgements



- Yi-Siang Wang
- Ziyang Cao
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- Victor Suarez



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