

Real-Time GW-BSE Investigation on Spin-Valley Exciton Dynamics

Jin Zhao

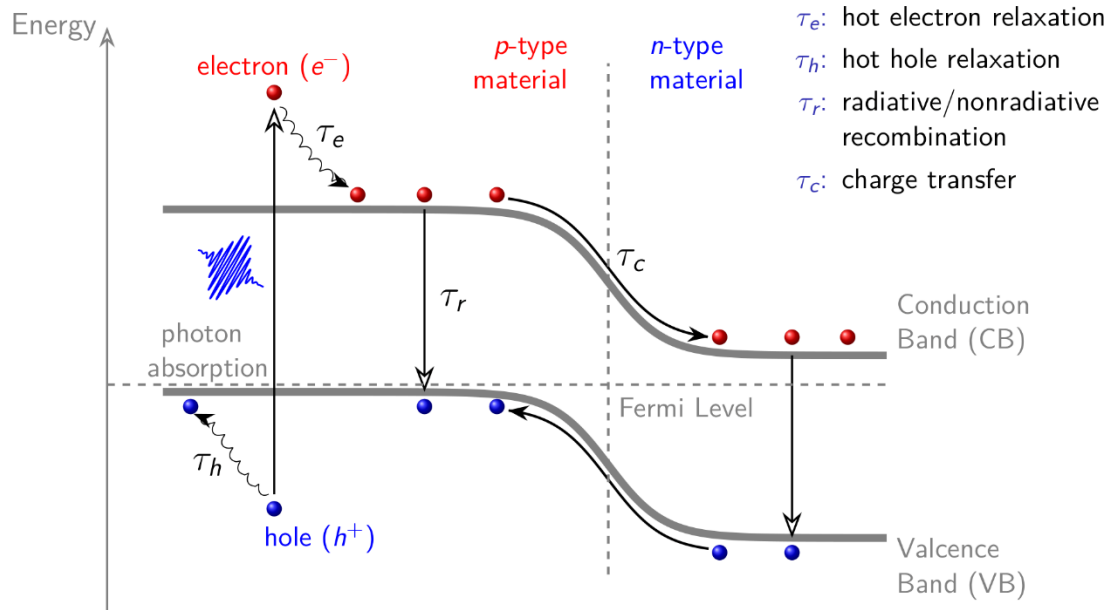
**University of Science and
Technology, China (USTC)**

VISTA Seminar

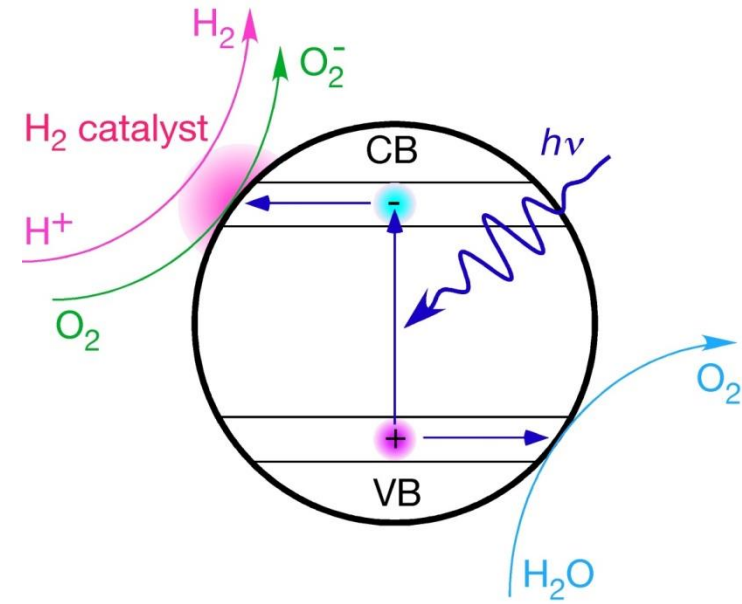
Seminar 25

10/14/2021

Excited Carrier Dynamics in Solar Energy Conversion



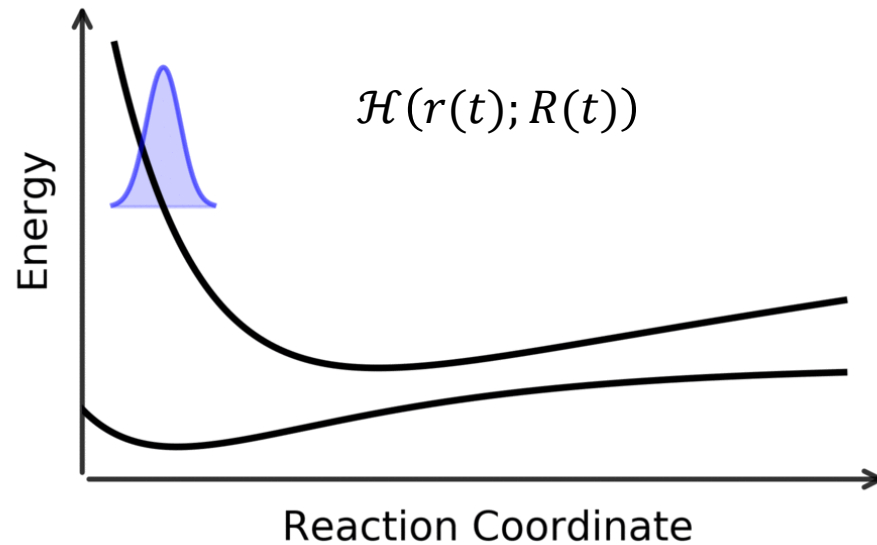
Photovoltaics



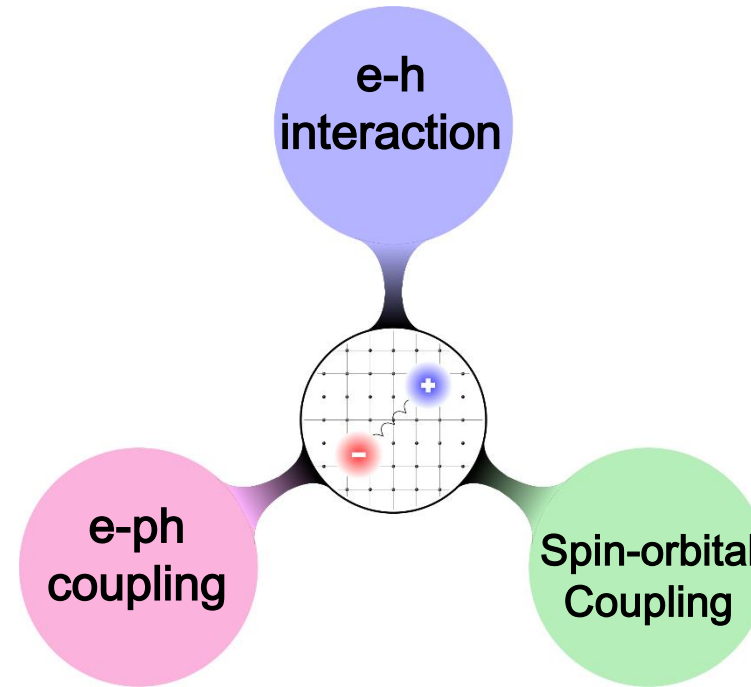
Photocatalysis

- hot carrier relaxation
- electron-hole recombination
- interfacial charge transfer

Challenges for *Ab Initio* Calculations

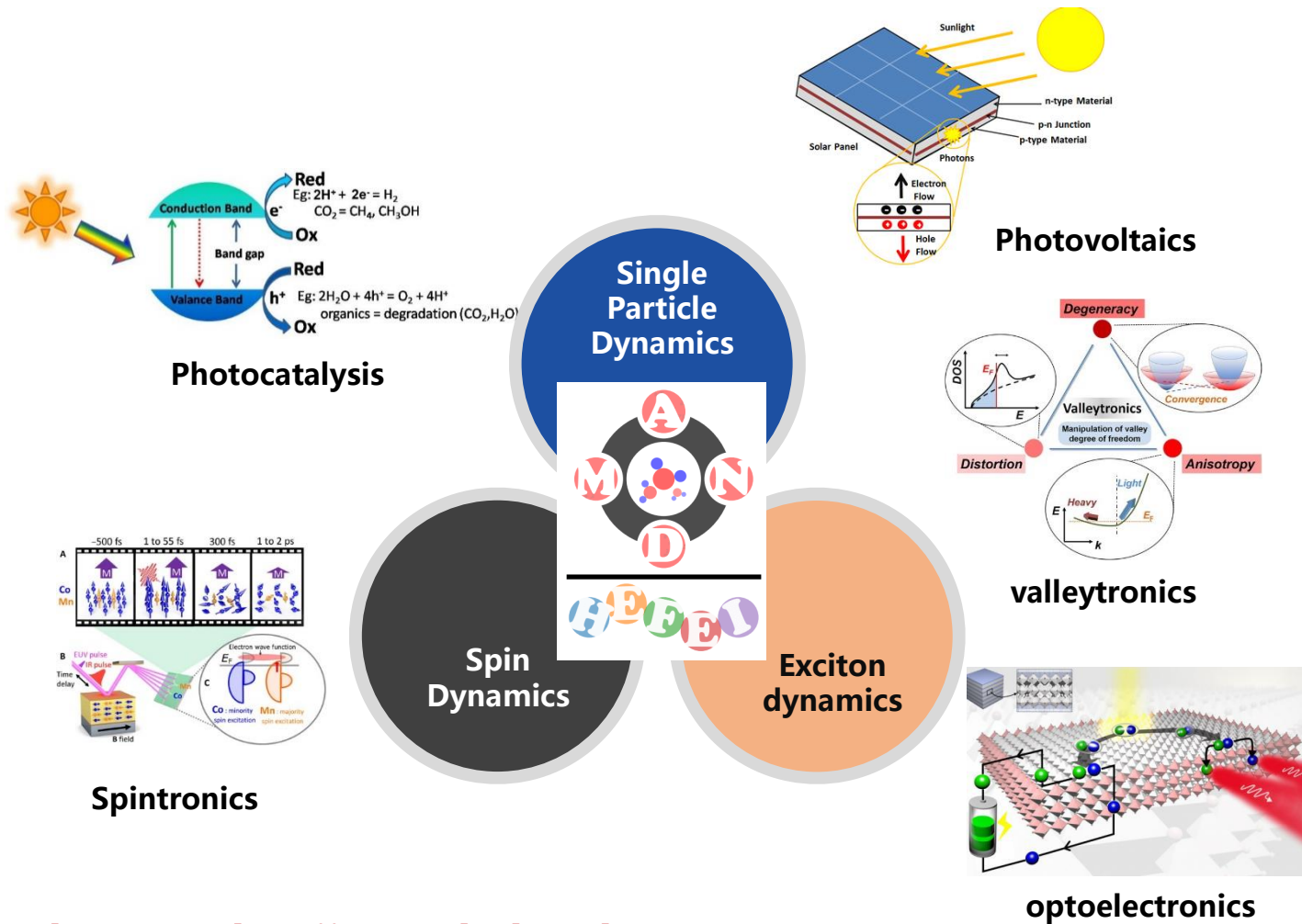


Time-dependent dynamics
Nonadiabatic effects



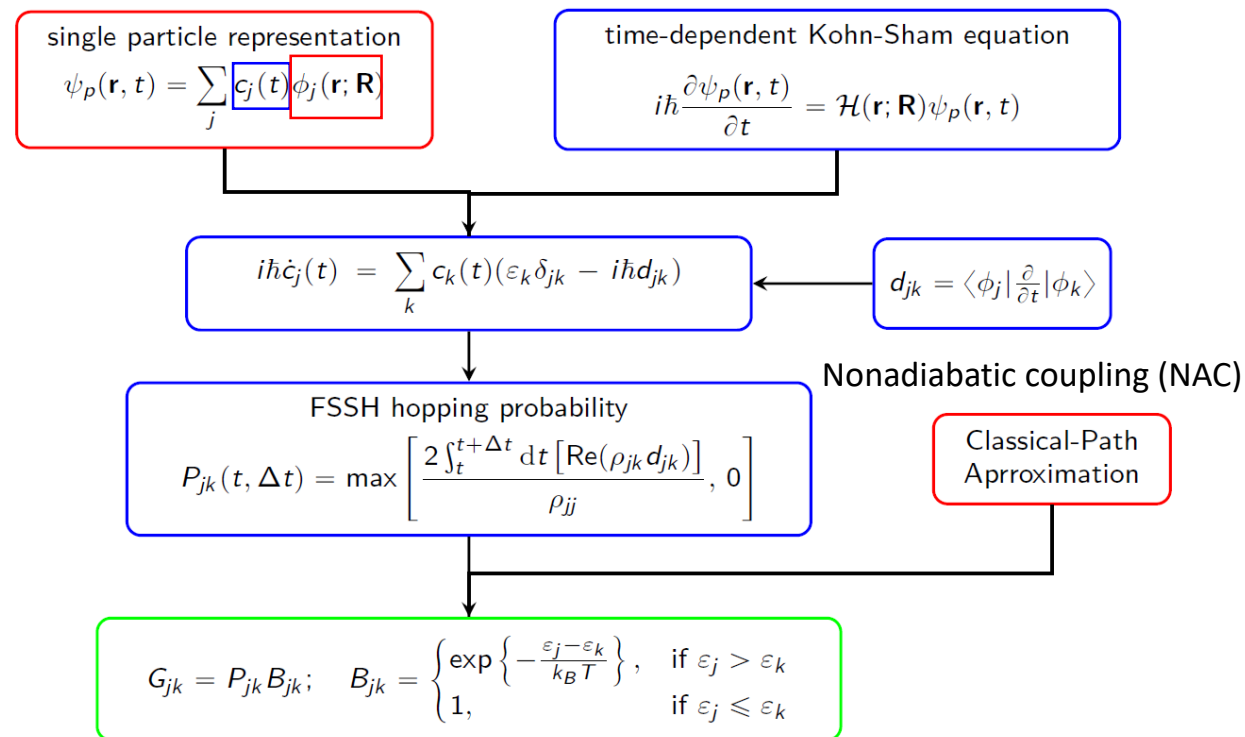
Different interactions

Hefei – Nonadiabatic Molecular Dynamics (Hefei-NAMD)



<http://staff.ustc.edu.cn/~zhaojin/code.html>

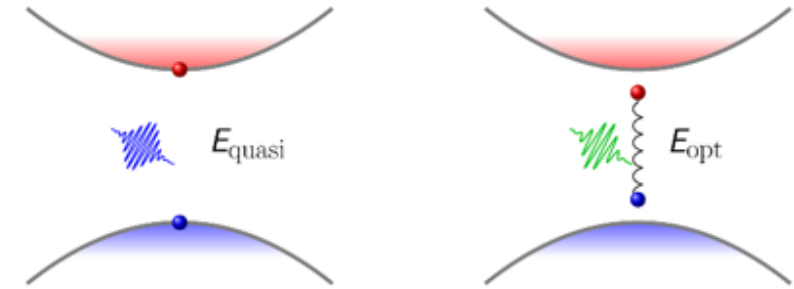
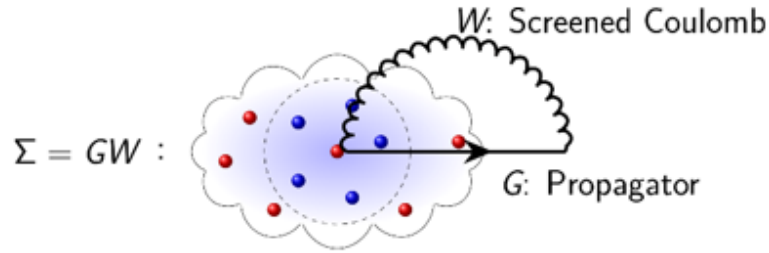
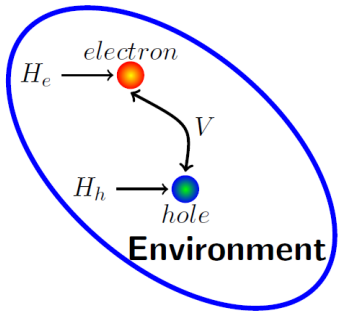
Surface Hopping Combined with TDKS



A. Akimov and O. Prezhdo: Pyxaid

Q. Zheng, X. Jiang, et. al. J. Zhao: Hefei-NAMD

GW + BSE to Describe the Exciton



$$\Sigma^{\text{GW}}(\mathbf{r}, \mathbf{r}', \omega) = -\frac{i}{2\pi} \int d\omega' e^{i\omega'\eta} G(\mathbf{r}, \mathbf{r}', \omega + \omega') W(\mathbf{r}, \mathbf{r}', \omega')$$

$$H_{c'v'k'}^{cvk} = [E_{ck}^{\text{QP}} - E_{vk}^{\text{QP}}] \delta_{cc'} \delta_{vv'} \delta_{kk'} - W_{c'v'k'}^{cvk} + 2v_{c'v'k'}$$

GW: self-energy take place of exchange correlation potential

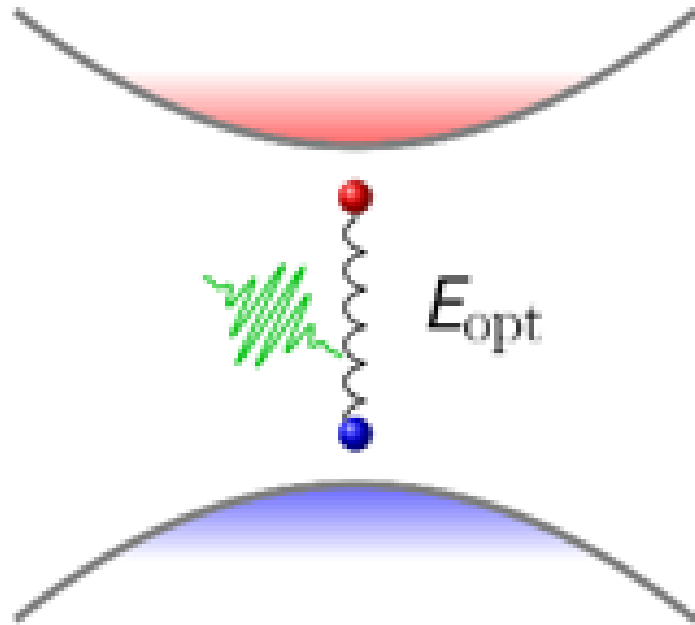
Accurate Quasi-particle energies



Bethe-Salpeter Equation (BSE): Screened Coulomb (*W*) and exchange (*v*) interaction of electron and hole

Exciton binding energy and wave function

Exciton Dynamics is Important



Exciton dynamics:

Exciton relaxation

Bright-to-dark transition

Single-to-multi transition

Singlet fission

Exciton annihilation (radiative and nonradiative)

...

10,000 times of GW+BSE calculations are too expensive!

Realization of *GW* + Real-Time BSE

BSE Hamiltonian:

From GW

$$\langle \mathbf{k}c\nu | H | \mathbf{k}'c'\nu' \rangle = \left(E_{\mathbf{k}c}^{QP} - E_{\mathbf{k}\nu}^{QP} \right) \delta_{\mathbf{k}\mathbf{k}'} \delta_{cc'} \delta_{\nu\nu'} - W_{\mathbf{k}'c'\nu'}^{\mathbf{k}c\nu} + v_{\mathbf{k}'c'\nu'}^{\mathbf{k}c\nu}.$$

Coulomb Interaction:

$$W_{\mathbf{k}'c'\nu'}^{\mathbf{k}c\nu} = \frac{1}{\Omega} \sum_{\mathbf{G}\mathbf{G}'} \frac{4\pi \varepsilon_{\mathbf{G}\mathbf{G}'}^{-1}(\mathbf{k} - \mathbf{k}')}{|\mathbf{k} - \mathbf{k}' + \mathbf{G}| |\mathbf{k} - \mathbf{k}' + \mathbf{G}'|} \left(B_{\uparrow\mathbf{k}'c'}^{\uparrow\mathbf{k}c}(\mathbf{G}) + B_{\downarrow\mathbf{k}'c'}^{\downarrow\mathbf{k}c}(\mathbf{G}) \right) \left(B_{\uparrow\mathbf{k}'\nu'}^{\uparrow\mathbf{k}\nu*}(\mathbf{G}') + B_{\downarrow\mathbf{k}'\nu'}^{\downarrow\mathbf{k}\nu*}(\mathbf{G}') \right)$$

Exchange Interaction:

$$v_{\mathbf{k}'c'\nu'}^{\mathbf{k}c\nu} = \frac{1}{\Omega} \sum_{\mathbf{G} \neq 0} \frac{4\pi}{|\mathbf{G}|^2} \left(B_{\uparrow\mathbf{k}\nu}^{\uparrow\mathbf{k}c}(\mathbf{G}) + B_{\downarrow\mathbf{k}\nu}^{\downarrow\mathbf{k}c}(\mathbf{G}) \right) \left(B_{\uparrow\mathbf{k}'\nu'}^{\uparrow\mathbf{k}'c'*}(\mathbf{G}) + B_{\downarrow\mathbf{k}'\nu'}^{\downarrow\mathbf{k}'c'*}(\mathbf{G}) \right),$$

from time-dependent
Kohn-Sham basis sets

Rigid dielectric function during MD

QP energy: Rigid shift from KS energy

10,000 real-time *GW* + BSE
– 1 *GW* + real-time BSE

Diabatic Representation

Single-particle

$$\text{TDDFT}$$

$$i\hbar \frac{\partial \psi(r, t)}{\partial t} = \mathcal{H}(r; R) \psi(r, t)$$



$$\text{Real-time - BSE}$$

$$i\hbar \frac{\partial \psi(r_e, r_h, t)}{\partial t} = \mathcal{H}(r; R) \psi(r_e, r_h, t)$$

two-particle

Diabatic:

$$H_{c'v'k'}^{cvk} = \left(E_{ck}^{QP} - E_{vk}^{QP} \right) \delta_{cc'} \delta_{vv'} \delta_{kk'} - W_{c'v'k'}^{cvk} + 2v_{c'v'k'}^{cvk}$$

$$- i\hbar \left(\langle ck | \frac{\partial}{\partial t} | c'k \rangle \delta_{vv'} + \langle vk | \frac{\partial}{\partial t} | v'k \rangle^* \delta_{cc'} \right) \delta_{kk'}$$

Adiabatic Representation

$$H_{c'v'k'}^{cvk} = \left(E_{ck}^{QP} - E_{vk}^{QP} \right) \delta_{cc'} \delta_{vv'} \delta_{kk'} - W_{c'v'k'}^{cvk} + 2v_{c'v'k'}^{cvk}$$

Diagonalize the BSE Hamiltonian to get the exciton basis sets

Adiabatic:

Expand the state ket $|\Psi\rangle$ in time-dependent two-particle Schrödinger equation

$$i\hbar \frac{\partial |\Psi(r_e, r_h, t)\rangle}{\partial t} = H |\Psi(r_e, r_h, t)\rangle$$

using the following basis:

$$|x_I\rangle = \sum_{cvk} f_{cvk}^I |cvk\rangle$$

where $|x_I\rangle$ is the eigenstate of exciton:

$$H |x_I\rangle = E_I^x |x_I\rangle$$

Adiabatic:

$$|\Psi\rangle = \sum_I A_I |x_I\rangle$$

$$H_{ki} = \epsilon_k \delta_{ik} - \boxed{i\hbar \left\langle k \left| \frac{\partial}{\partial t} \right| i \right\rangle}$$

$$i\hbar \dot{A}_K = \sum_I H_{KI} A_I$$

Single particle

$$H_{KI} = E_K^x \delta_{KI} - \boxed{i\hbar \left\langle x_K \left| \frac{\partial}{\partial t} \right| x_I \right\rangle}$$

$$\boxed{\left\langle x_K \left| \frac{\partial}{\partial t} \right| x_I \right\rangle}$$

Exciton-phonon interaction

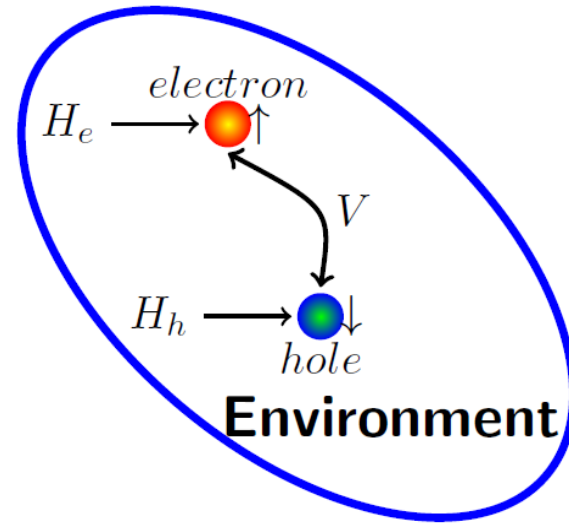
Non-adiabatic coupling for exciton

$$\begin{aligned} &= \sum_{\substack{cvk \\ c'v'k'}} f_{c'v'k'}^I f_{cvk}^{*K} \left\langle cvk \left| \frac{\partial}{\partial t} \right| c'v'k' \right\rangle + \sum_{cvk} \dot{f}_{cvk}^I f_{cvk}^{*K} \\ &= \sum_{\substack{cvk \\ c'v'}} f_{c'v'k}^I f_{cvk}^{*K} \left(\left\langle ck \left| \frac{\partial}{\partial t} \right| c'k \right\rangle \delta_{vv'} + \left\langle vk \left| \frac{\partial}{\partial t} \right| v'k \right\rangle^* \delta_{cc'} \right) \\ &\quad + \sum_{cvk} \left\langle f_{cvk}^K \left| \frac{d}{dt} \right| f_{cvk}^I \right\rangle \\ &= \sum_{cc'vk} f_{cvk}^{*K} \left\langle ck \left| \frac{\partial}{\partial t} \right| c'k \right\rangle f_{c'vk}^I + \sum_{vv'ck} f_{cvk}^{*K} \left\langle vk \left| \frac{\partial}{\partial t} \right| v'k \right\rangle^* f_{cv'k}^I \\ &\quad + \frac{1}{2dt} \sum_{cvk} \left(\left\langle f_{cvk}^K(t) \left| f_{cvk}^I(t+dt) \right\rangle - \left\langle f_{cvk}^K(t+dt) \left| f_{cvk}^I(t) \right\rangle \right) \right) \end{aligned}$$

Exciton with Spin

Hamiltonian:

$$H = H_e + H_h + V_{e-h} + H^{SO}$$



Exciton dynamics with SOC

$$\begin{aligned}
 & \langle c\sigma_c v\sigma_v k | H | c'\sigma_{c'} v'\sigma_{v'} k' \rangle \\
 = & \left(E_{c\sigma_c k}^{QP} \delta_{cc'} \delta_{\sigma_c \sigma_{c'}} + \langle c\sigma_c k | H^{SO} | c'\sigma_{c'} k \rangle \right) \delta_{vv'} \delta_{\sigma_v \sigma_{v'}} \delta_{kk'} \\
 & - \left(E_{v\sigma_v k}^{QP} \delta_{vv'} \delta_{\sigma_v \sigma_{v'}} - \langle v'\sigma_{v'} k | H^{SO} | v\sigma_v k \rangle \right) \delta_{cc'} \delta_{\sigma_c \sigma_{c'}} \delta_{kk'} \\
 & - W_{c'\sigma_{c'} v'\sigma_{v'} k'}^{c\sigma_c v\sigma_v k} \delta_{\sigma_c \sigma_{c'}} \delta_{\sigma_v \sigma_{v'}} + v_{c'\sigma_{c'} v'\sigma_{v'} k'}^{c\sigma_c v\sigma_v k} \delta_{\sigma_c \sigma_v} \delta_{\sigma_{c'} \sigma_{v'}}
 \end{aligned}$$

Screened coulomb interaction

$$W_{c'\sigma_{c'} v'\sigma_{v'} k'}^{c\sigma_c v\sigma_v k} = \frac{1}{\Omega} \sum_{GG'} \frac{4\pi \epsilon_{GG'}^{-1} (k - k')}{|k - k' + G| |k - k' + G'|} B_{c'\sigma_{c'} k'}^{c\sigma_c k}(G) B_{v'\sigma_{v'} k'}^{v\sigma_v k*}(G')$$

Exchange interaction

$$v_{c'\sigma_{c'} v'\sigma_{v'} k'}^{c\sigma_c v\sigma_v k} = \frac{1}{\Omega} \sum_{G \neq 0} \frac{4\pi}{|G|^2} B_{v\sigma_v k}^{c\sigma_c k}(G) B_{v'\sigma_{v'} k'}^{c'\sigma_{c'} k'*}(G)$$

$$\langle n\sigma k | H^{SO} | n'\sigma' k \rangle = \sum_{\alpha, i, j} \langle \tilde{\psi}_{n\sigma k} | \tilde{p}_{i\sigma k}^\alpha \rangle \langle \phi_{i\sigma k}^\alpha | H^{SO} | \phi_{j\sigma' k'}^\alpha \rangle \langle \tilde{p}_{j\sigma' k'}^\alpha | \tilde{\psi}_{n'\sigma' k'} \rangle$$

Realization of *GW* + Real-Time BSE

Single-particle

$$\text{TDDFT}$$
$$i\hbar \frac{\partial \psi(r, t)}{\partial t} = \mathcal{H}(r; R) \psi(r, t)$$

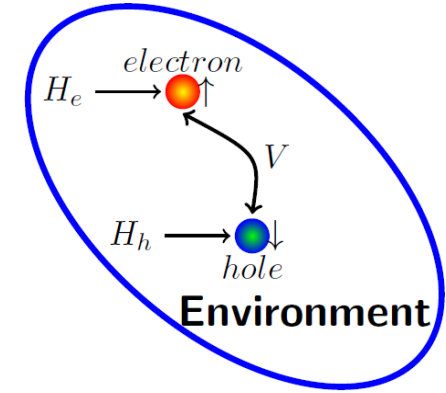


$$\text{Real-time - BSE}$$
$$i\hbar \frac{\partial \psi(r_e, r_h, t)}{\partial t} = \mathcal{H}(r; R) \psi(r_e, r_h, t)$$

two-particle

Hamiltonian:

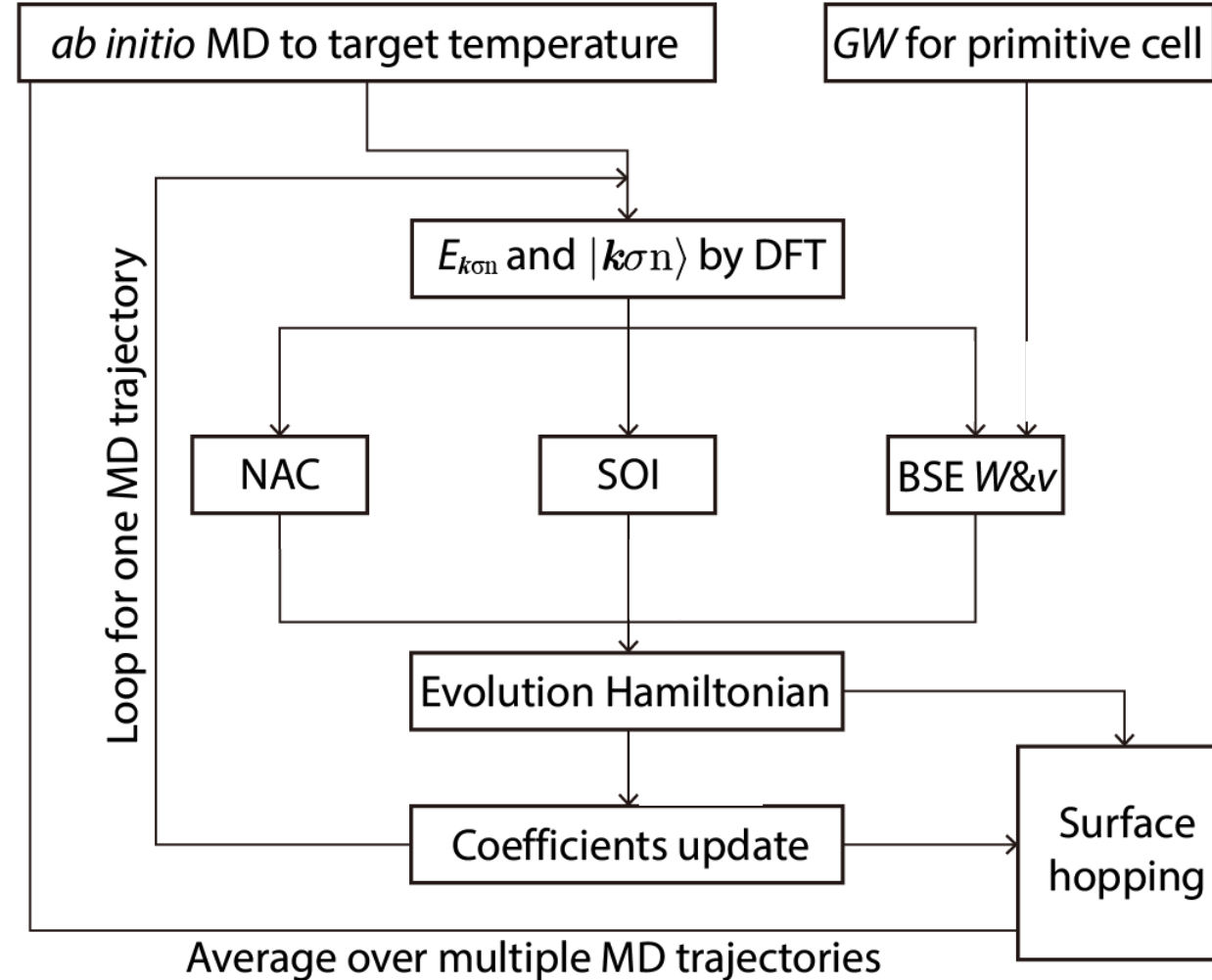
$$H = H_e + H_h + V_{e-h} + H^{SO}$$



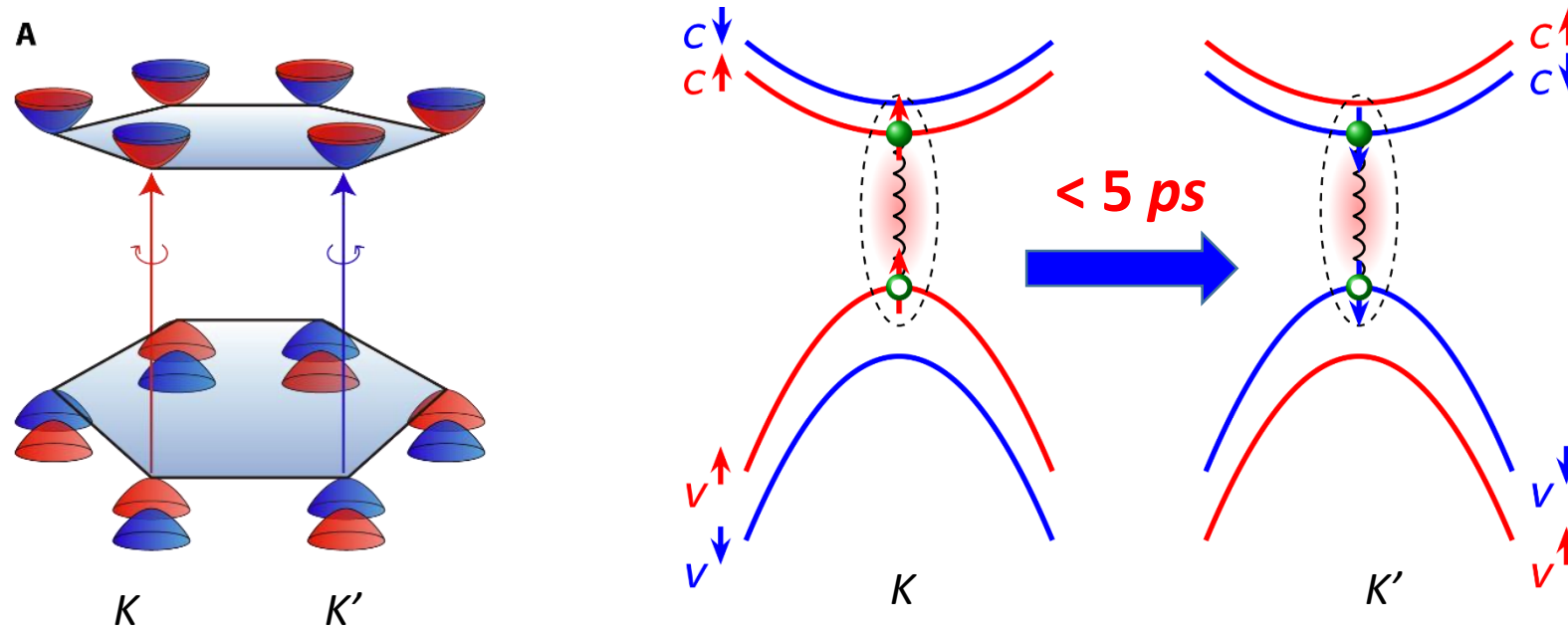
Spin orbital coupling

- ✓ **Many-body interaction:** Coulomb and exchange
- ✓ **Exciton-phonon interaction:** real-time BSE + molecular dynamics
- ✓ **Spin orbital coupling:** adiabatic and diabatic representation
- ✓ **Nonadiabatic:** surface hopping

Workflow of Real-Time GW-BSE NAMD



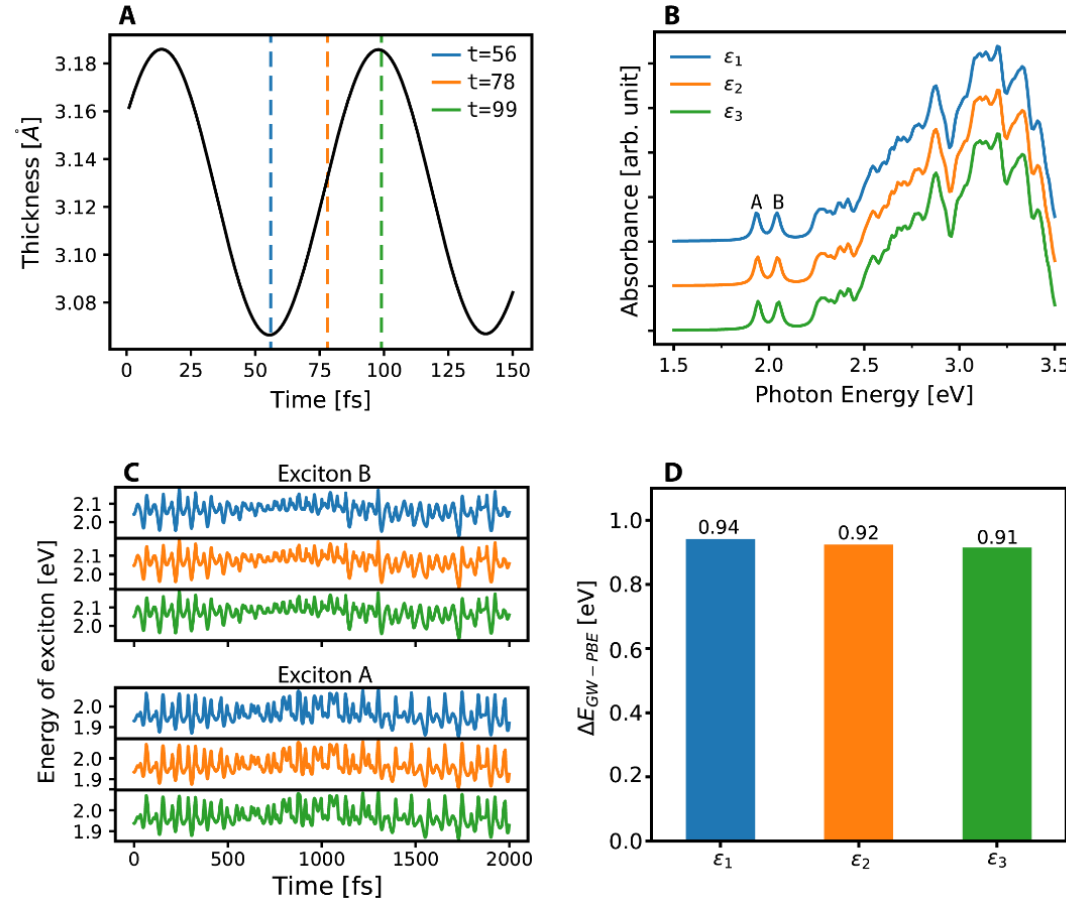
Fast Intervalley Bright Exciton Scattering in Transition Metal Dichalcogenide



Intervalley bright exciton scattering requires the **spin flip** and **momentum transition** of **both electron and hole**

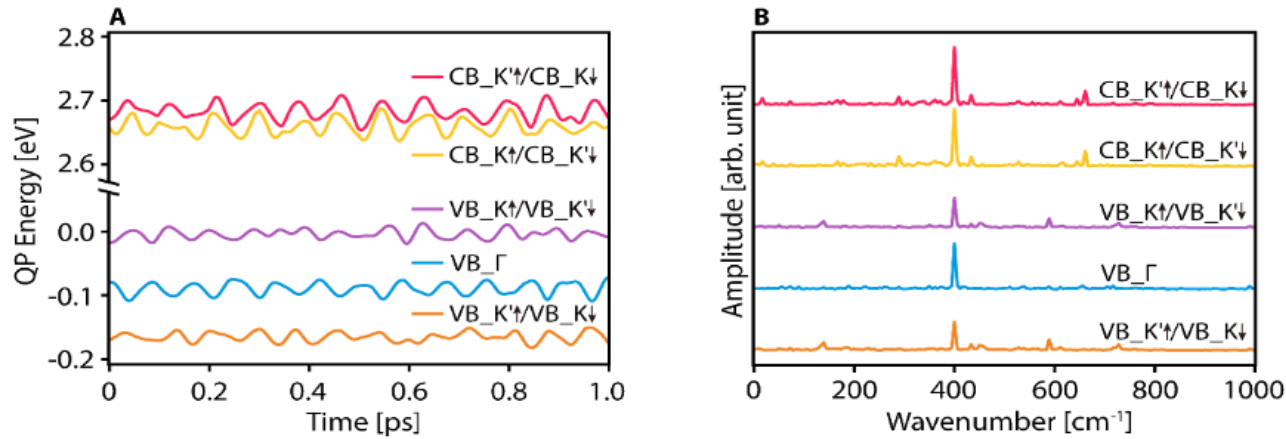
Puzzle: how can such intervalley bright exciton happen within **several picoseconds**?

Test of the Dielectric Function Approximation

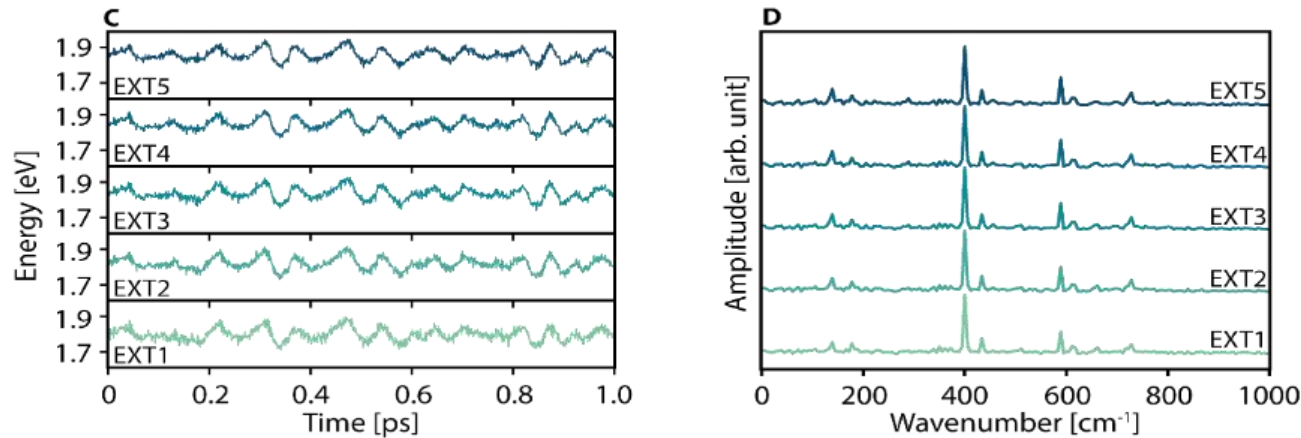
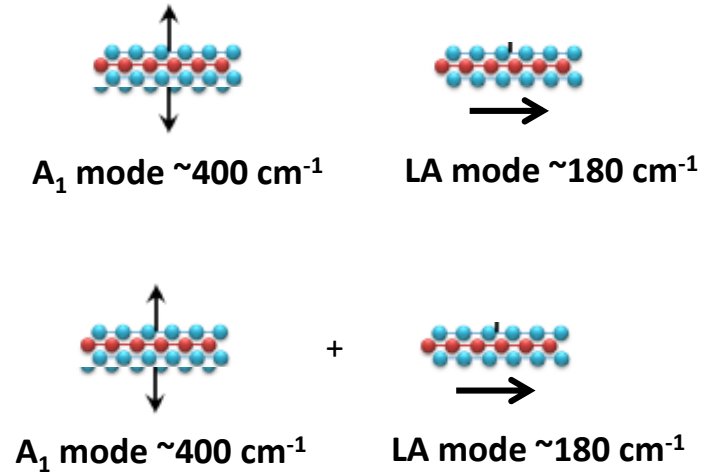


The **dielectric function** and ***GW*QP correction** almost **does not change** with the structure

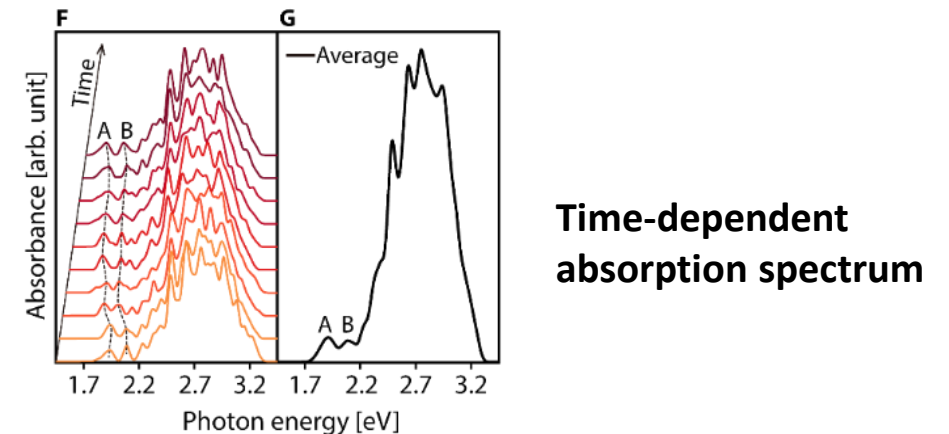
Time Dependent Quasi-Particle Energies and Optical Band Gap of MoS₂



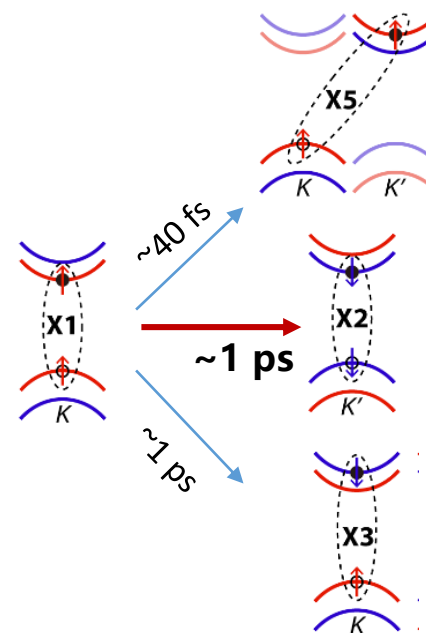
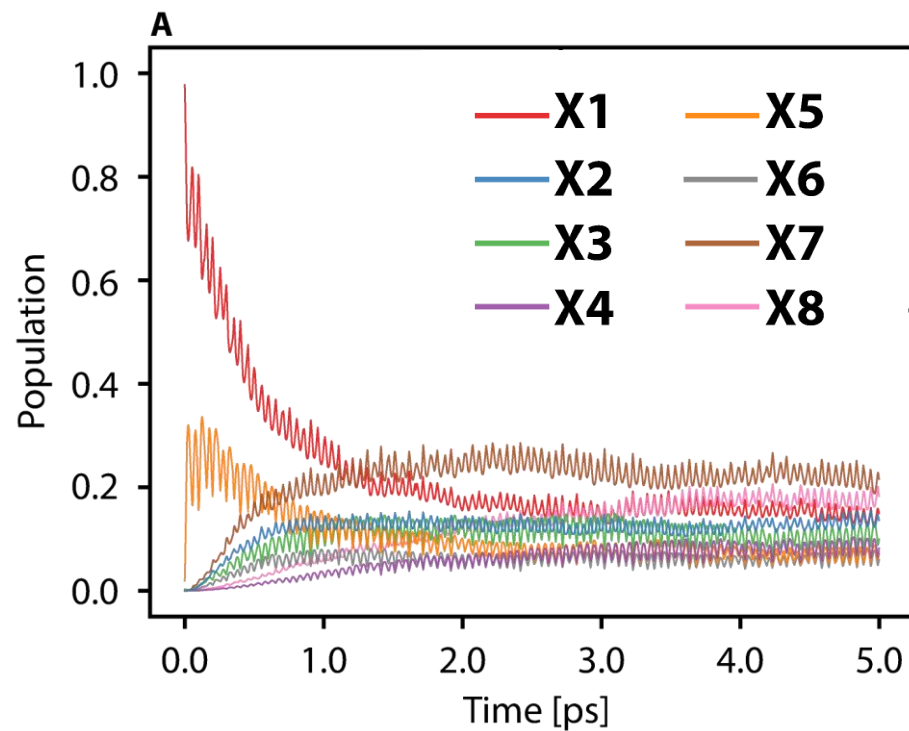
Electron-phonon interaction



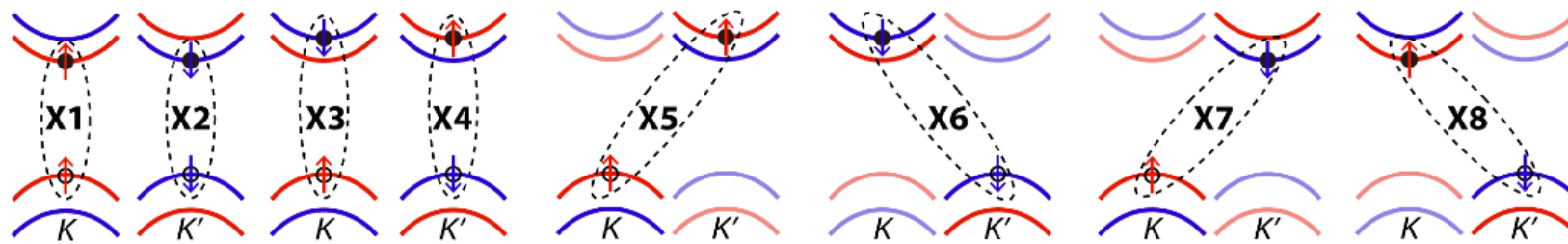
Exciton-phonon interaction



Exciton Dynamics in MoS₂



Bright Exciton transition happens in several ps

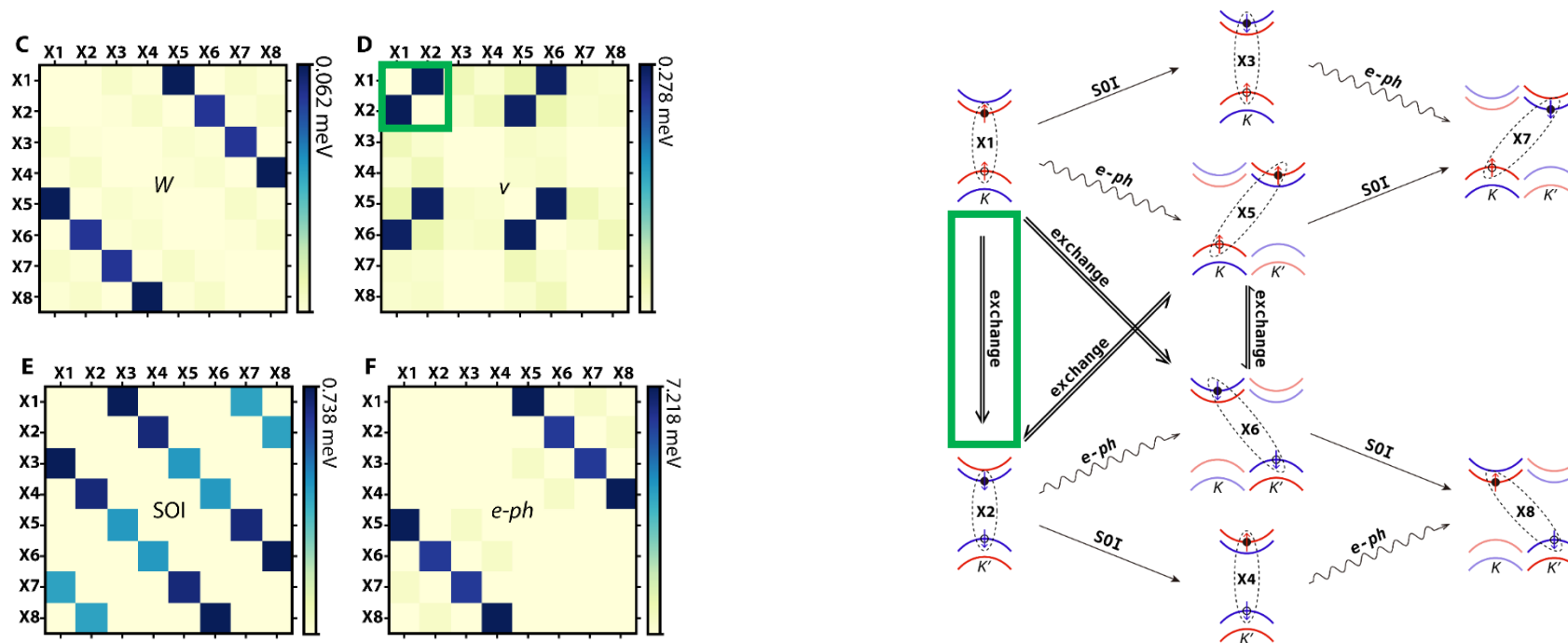


Exchange Interaction Induced Bright Exciton Scattering

Nonadiabatic Coupling Elements:

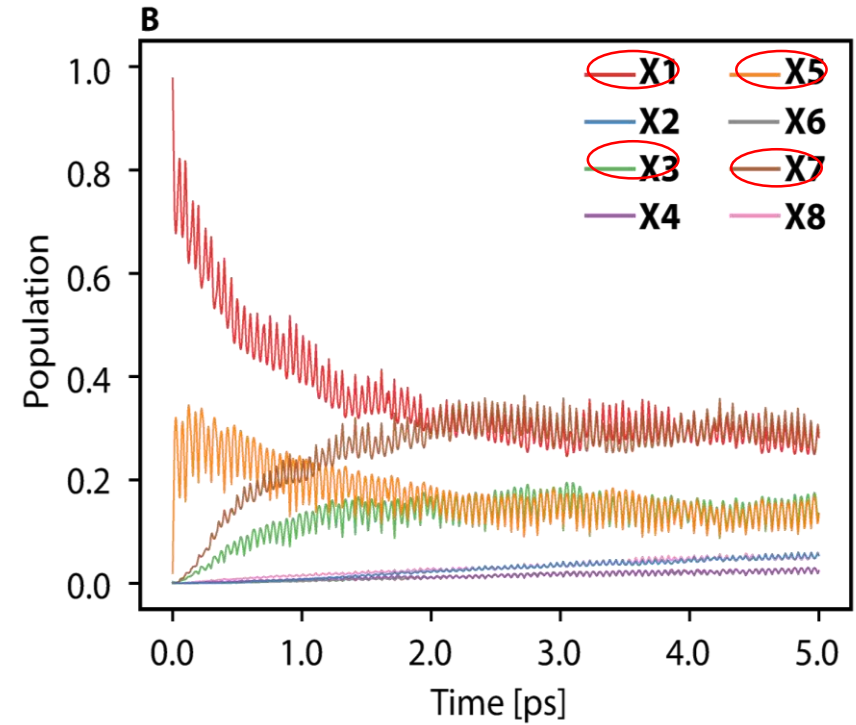
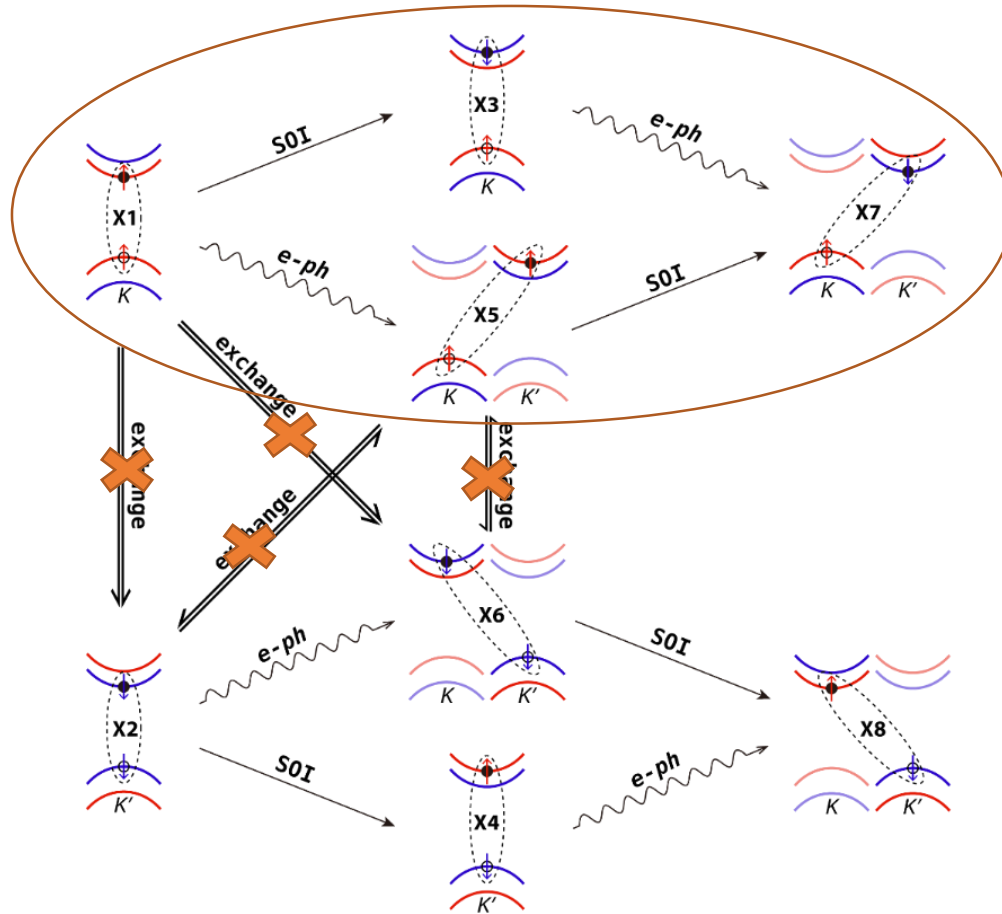
Single particle dynamics: $e-ph$

Exciton dynamics: $e-ph + W(e-h \text{ Coulomb}) + v(e-h \text{ exchange}) + \text{SOC}$



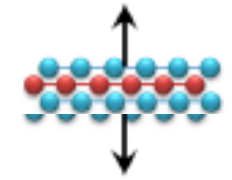
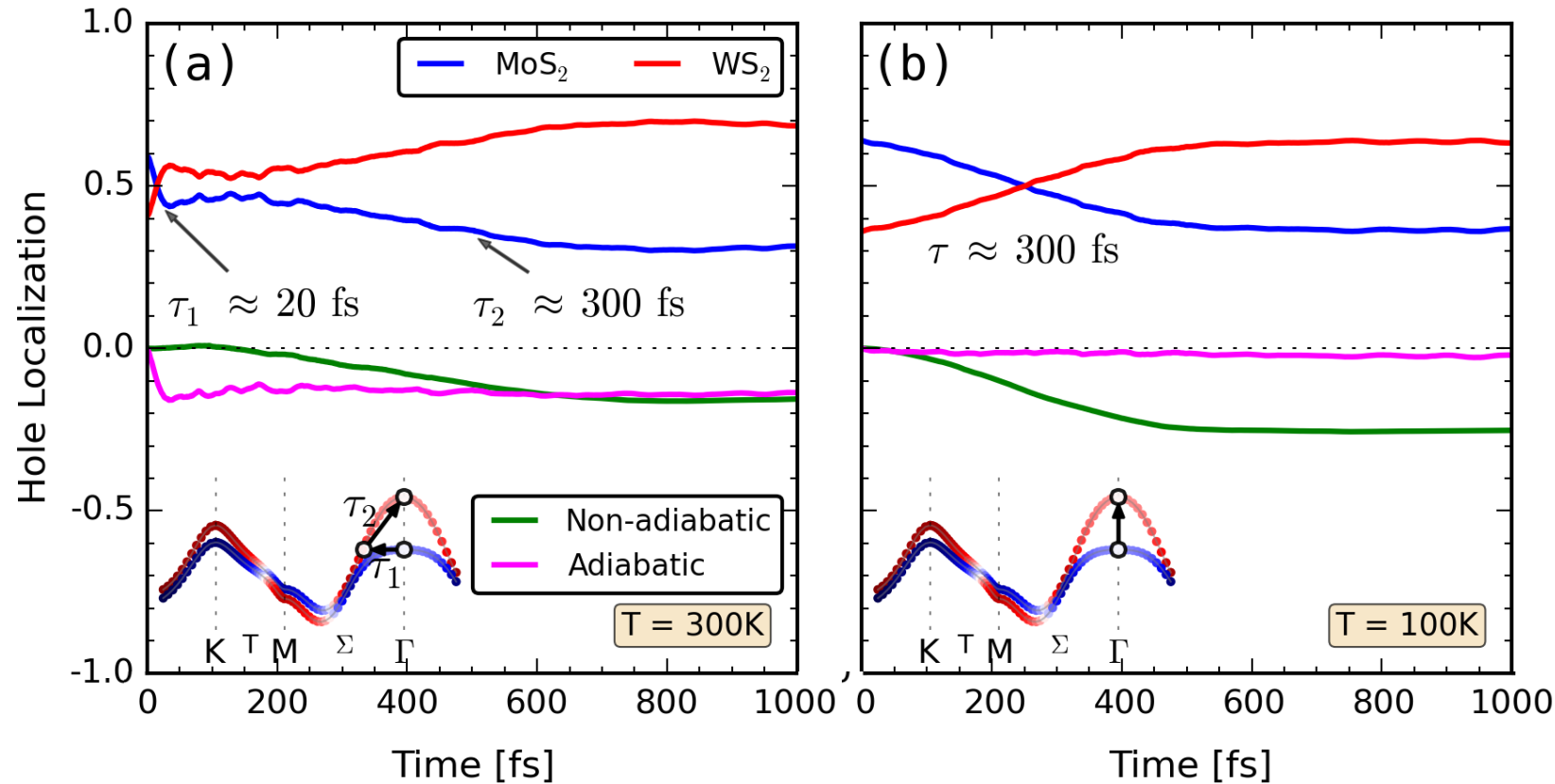
Bright Exciton transition is induced by $e-h$ exchange interaction

Single Particle Picture

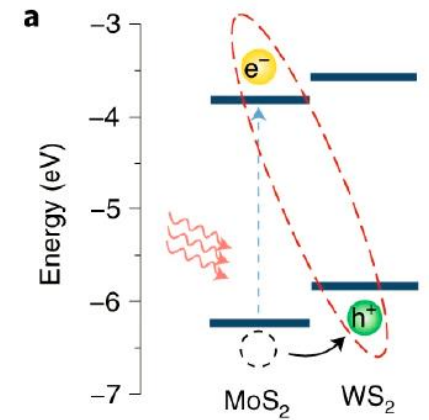


Spin particle picture: photoexcited hole keeps in K valley

Ultrafast Charge Transfer at TMD Heterostructure

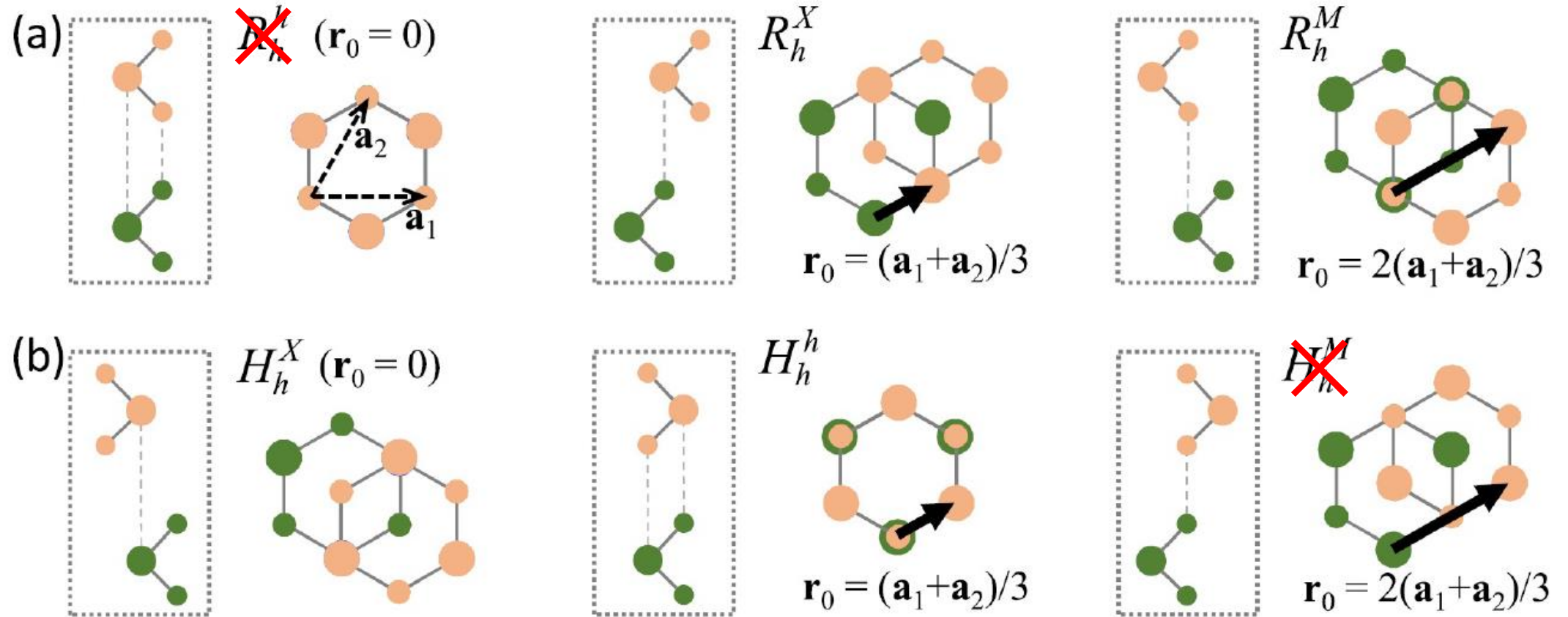


A₁ mode
~400 cm⁻¹
T=83 fs

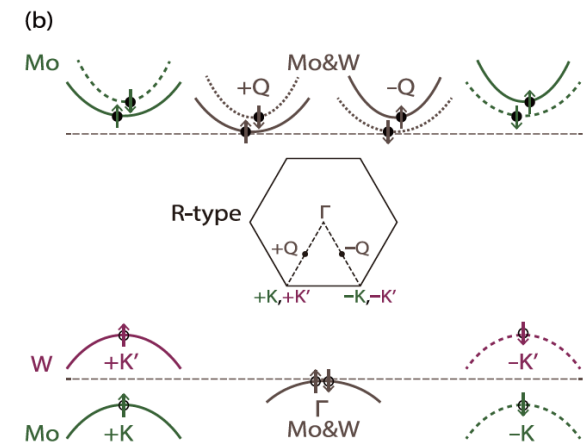
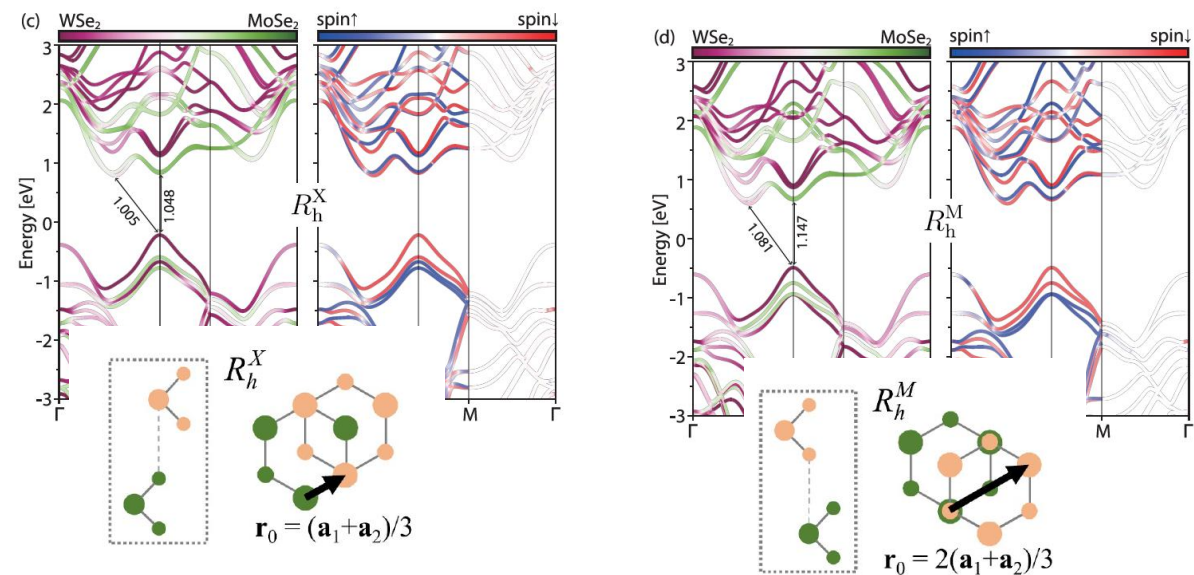
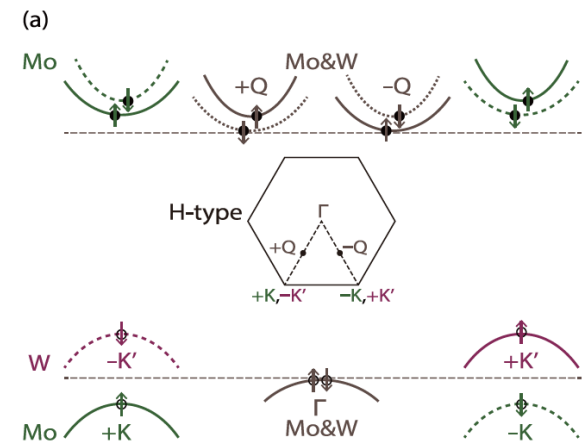
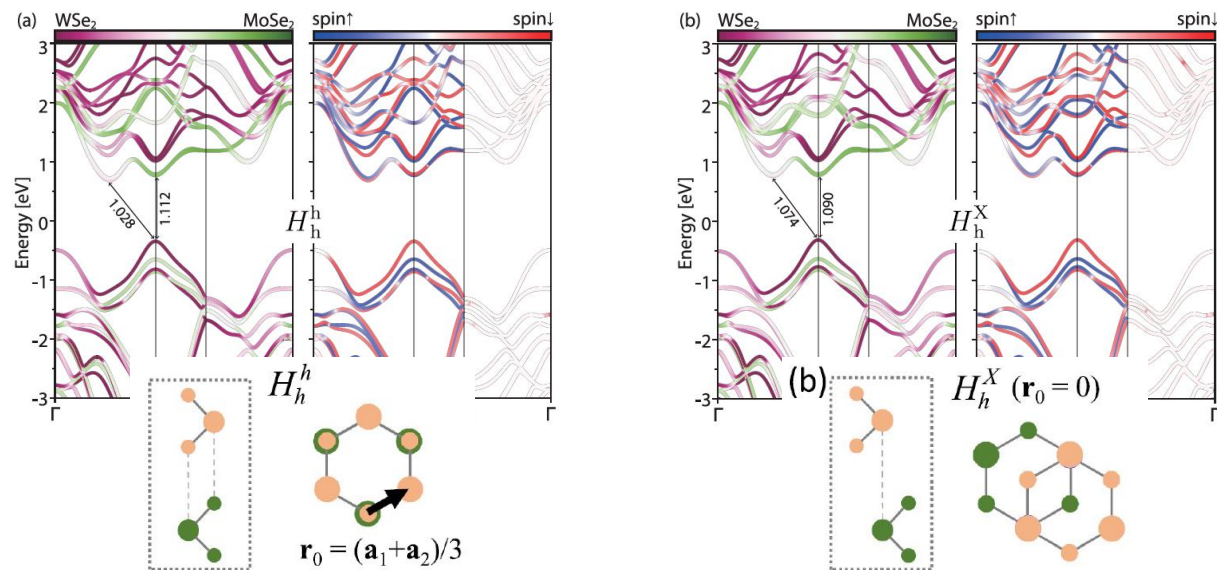


Q. Zheng, et.al. *Nano Lett.* **17**, 6435-6442 (2017)

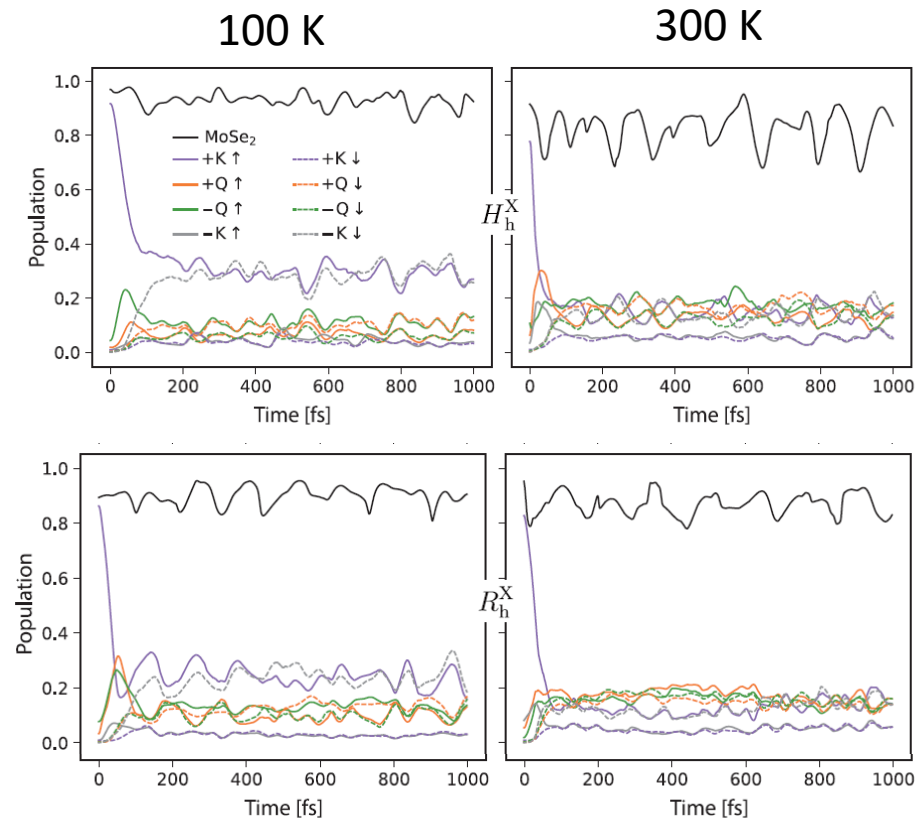
Different Stacking of MoSe₂/WSe₂



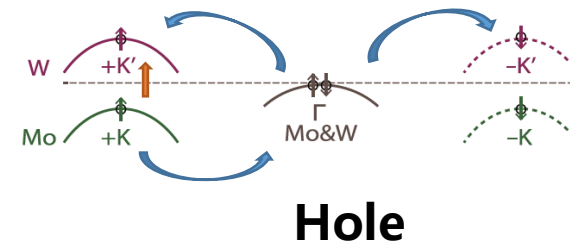
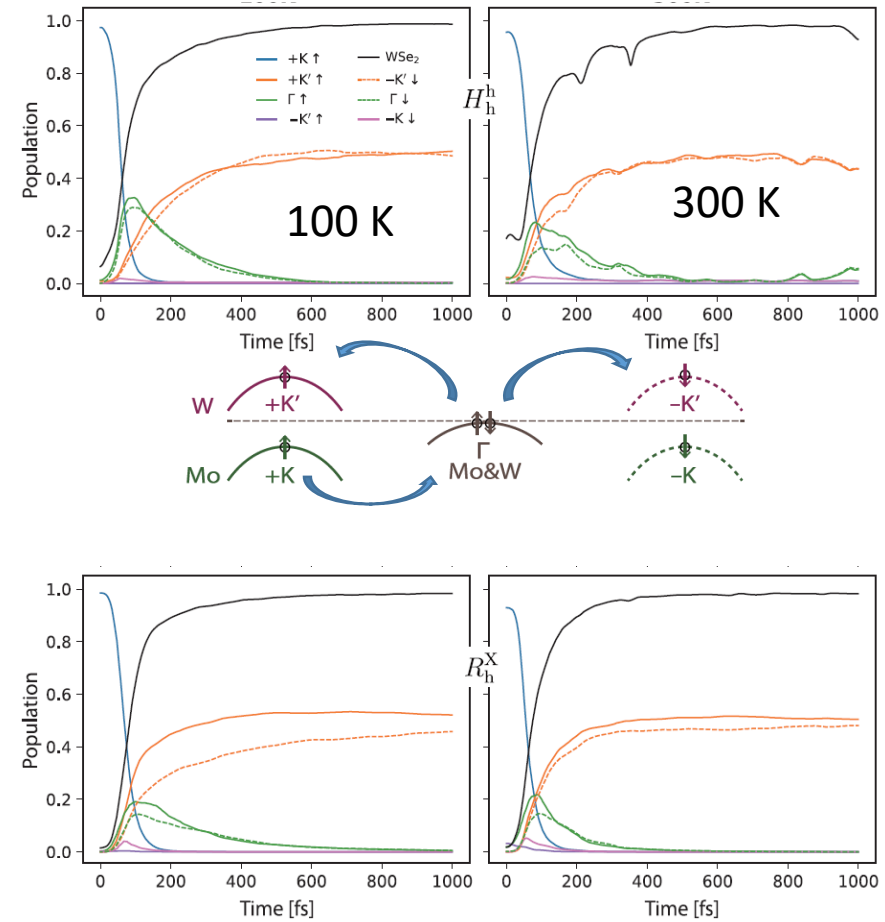
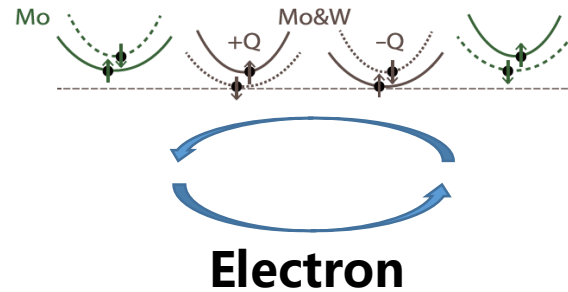
Band Structure of MoSe₂/WSe₂



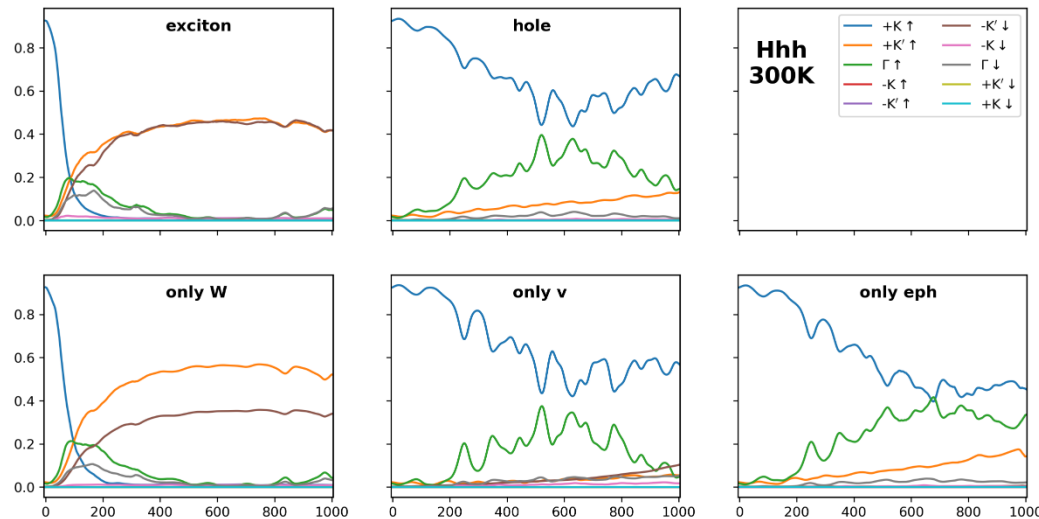
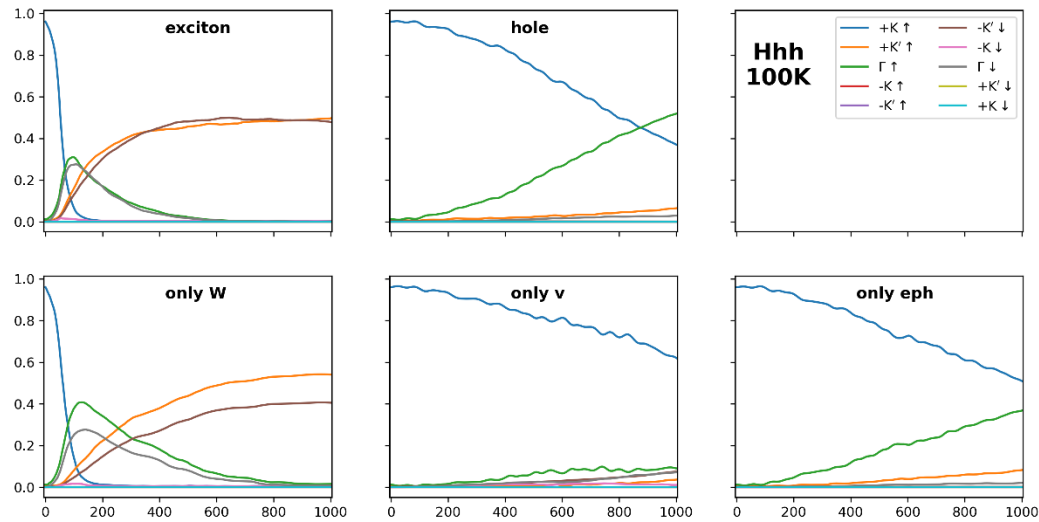
Exciton Dynamics at MoSe₂/WSe₂



(a)



Mechanism of Exciton Dynamics at MoSe₂/WSe₂



E-ph and SOC

Coulomb Interaction (W)

Exchange Interaction (v)

Summary

Single-particle

$$\text{TDDFT}$$
$$i\hbar \frac{\partial \psi(r, t)}{\partial t} = \mathcal{H}(r; R) \psi(r, t)$$



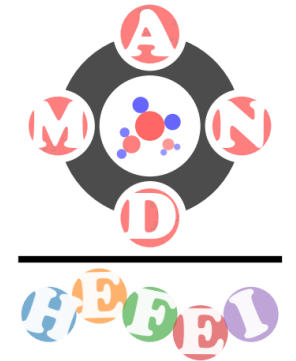
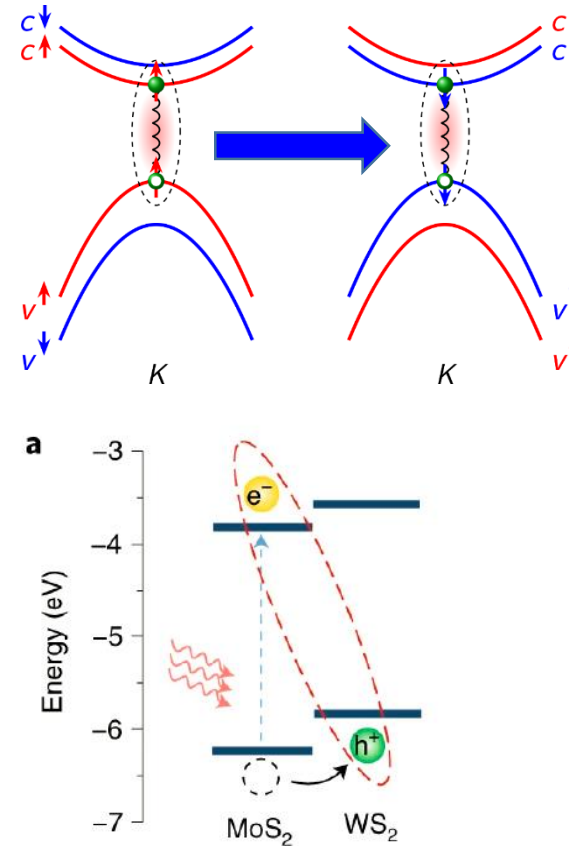
Real-time - BSE

$$i\hbar \frac{\partial \psi(r_e, r_h, t)}{\partial t} = \mathcal{H}(r; R) \psi(r_e, r_h, t)$$

two-particle

- ✓ Many-body interaction
- ✓ Exciton-phonon interaction
- ✓ Spin orbital coupling
- ✓ Nonadiabatic effects

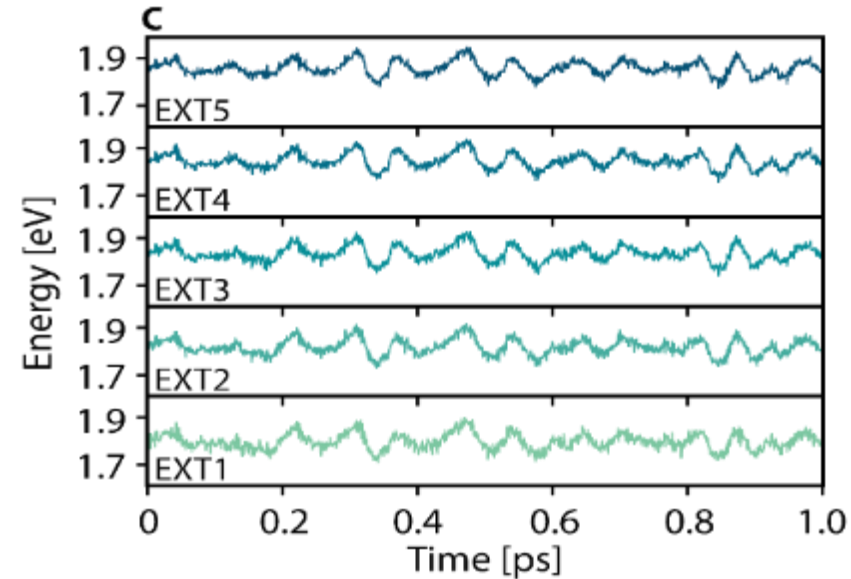
X. Jiang, Q. Zheng, Z. Lan, W. A. Saidi, X. Ren and J. Zhao*
Sci. Adv., **7**, eabf3759, (2021)



Xiang Jiang
蒋翔

Future work based on *GW* + rt-BSE NAMD Simulation

- ✓ Exciton Lifetime
- ✓ Hot exciton Relaxation
- ✓ Exciton transition
at interface
via spin valley
bright-to-dark
- ...
- ✓ Exciton-Phonon Interaction
- ✓ Exciton-Polaron interaction
- ...



Several thousands of exciton energies
from *GW* + rt-BSE NAMD

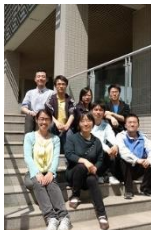


Machine learning for excited state
potential energy surface



Photocatalysis

Acknowledgement



2011



2013



2015



2016



2017



2019



2020

My Group



Zhenggang Lan
兰峥岗

South China Normal University
Surface hopping



Wissam A. Saidi
U. Pitt

DFT calculations



Xingguo Ren
任新国

USTC
GW+BSE



Hrvoje Petek

U. Pitt
Experiments



Jinlong Yang
杨金龙

USTC
Ph.D advisor
Discussion & support

Collaborators