

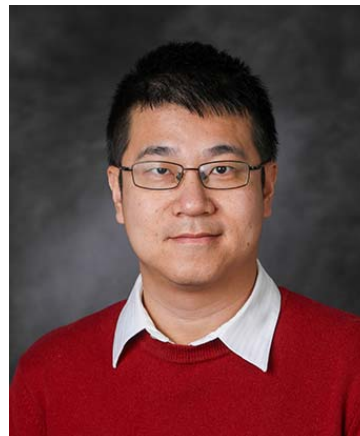
VISTA Talk

July 28, 2021

# I. Self-assembly of peptide amphiphiles

## II. Two-photon absorption with entangled photons

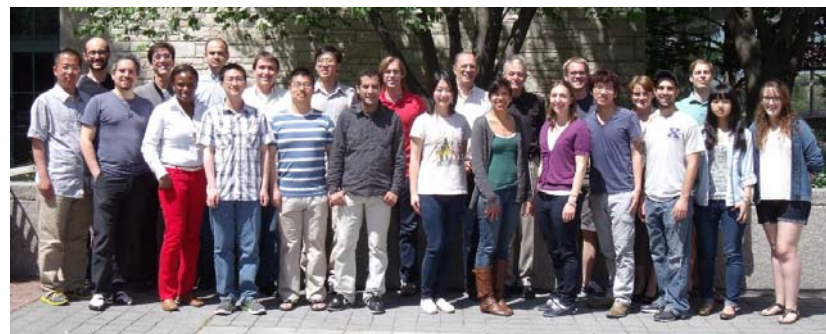
George C. Schatz  
Northwestern University



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Science | Basic Energy  
Sciences



## Tao's work as a postdoc (2012-2014) and thereafter



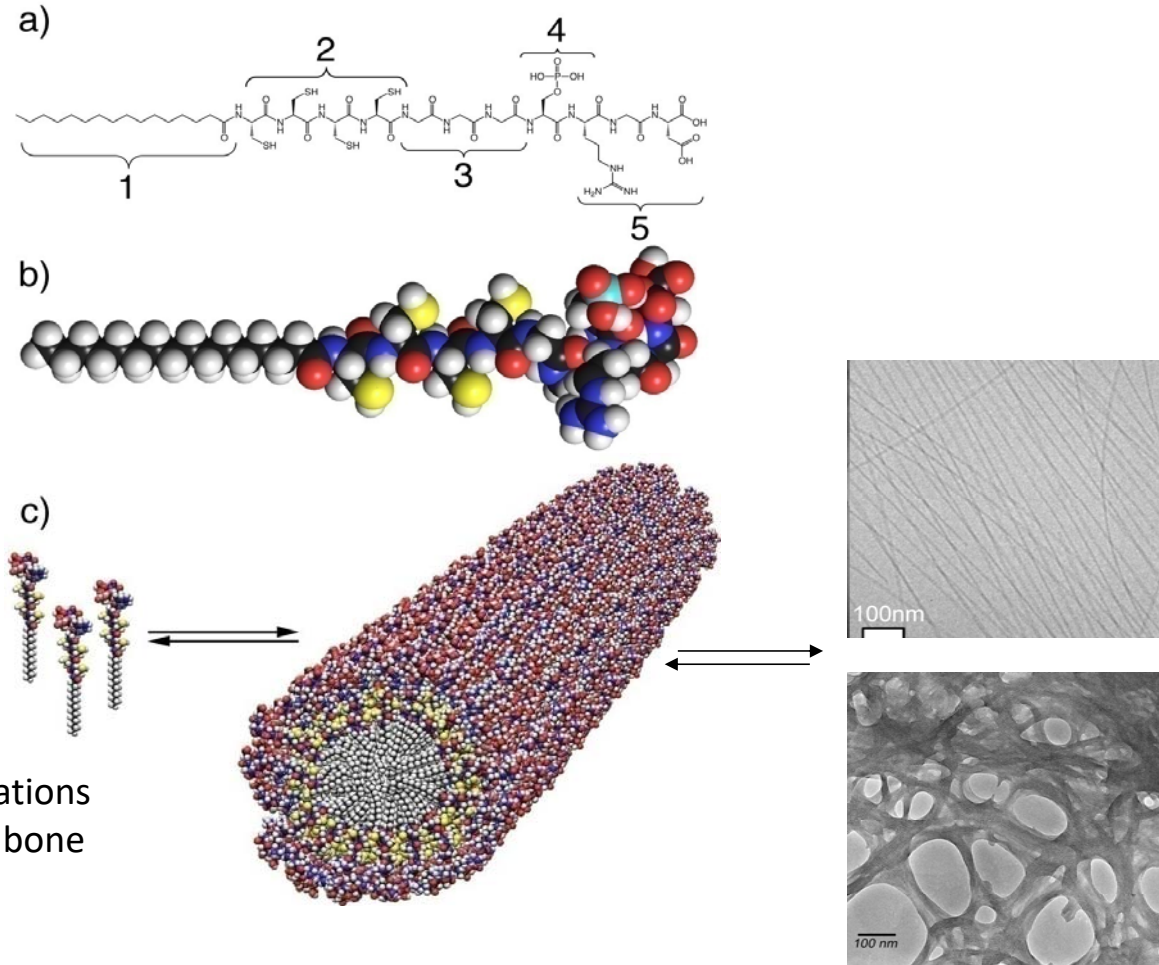
- 1) *Steered molecular dynamics studies of the potential of mean force for peptide amphiphile self-assembly into cylindrical nanofibers*, Tao Yu, One-Sun, Lee, George C. Schatz, *J. Phys. Chem. A*, 117, 9004-13 (2013)
- 2) *Free energy profile and mechanism of self-assembly of peptide amphiphiles based on a collective assembly coordinate*, Tao Yu, George C. Schatz, *J. Phys. Chem. B*, 117, 9004-9013(2013).
- 3) *Free-energy landscape for peptide amphiphile self-assembly: stepwise versus continuous assembly mechanisms*, Tao Yu, George C. Schatz, *J. Phys. Chem. B*, 117, 14059-14064 (2013).
- 4) *Molecular dynamics simulations and electronic excited state properties of a self-assembled peptide amphiphile nanofiber with metalloporphyrin arrays*, Tao Yu, One-Sun Lee and George C. Schatz, *J. Phys Chem. A*, 118, 8553-62 (2014).
- 5) *Energy landscapes and functions of supramolecular systems*, Faifan Tantakitti, Job Boekhoven, Xin Wang, Roman V. Kazantsev, Tao Yu, Jiahe Li, Ellen Zhuang, Roya Zandi, Julia H. Ortony, Christina J. Newcomb, Liam C. Palmer, Gajendra S. Shekhawat, Monic Olvera de la Cruz, George C. Schatz, Samuel I. Stupp, *Nature Materials*, 15, 469-477 (2016)
- 6) *Simultaneous covalent and noncovalent hybrid polymerizations*, Zhilin Yu, Faifan Tantakitti, Tao Yu, Liam C. Palmer, George C. Schatz, Samuel I. Stupp, *Science* 351, 497-502 (2016)
- 7) *A Mutation in Histone H2B Represents a New Class of Oncogenic Driver*, Richard L Bennett, Aditya Bele, Eliza C Small, Christine M Will, Behnam Nabet, Jon A Oyer, Xiaoxiao Huang, Rajarshi P Ghosh, Adrian T Grzybowski, Tao Yu, Qiao Zhang, Alberto Riva, Tanmay P Lele, George C Schatz, Neil L. Kelleher, Alexander J Ruthenburg, Jan Liphardt and Jonathan D Licht, *Cancer Discovery*, 9, 1438-51 (2019)

# Self-Assembly of Peptide Amphiphiles

J.D. Hartgerink, E. Beniash, and S.I. Stupp, *Science*, **294**, 1605 (2001).



Sam Stupp



Fibers have applications in wound healing, bone growth, nerve regeneration, etc

**Methods for describing the assembly of structures driven by hydrogen bonds and other weak interactions:**

Monte Carlo, Molecular Dynamics, Molecular Mechanics

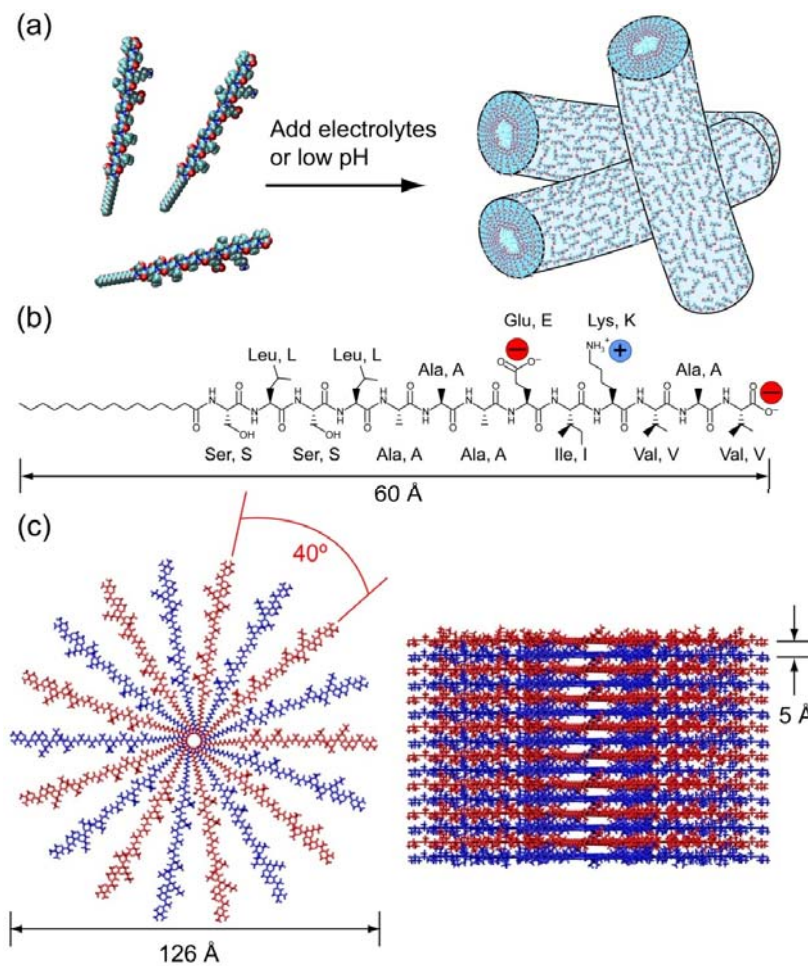
Atomistic, Coarse-grained force fields

Structural Models and Self-Assembly Tricks

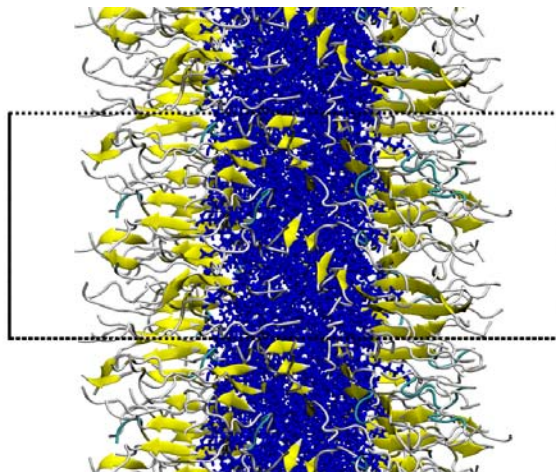


# Atomistic Modeling with Seeded Structures

One-Sun Lee, Samuel I. Stupp and George C. Schatz, J. Am. Chem. Soc. 113, 3677-3683 (2011)

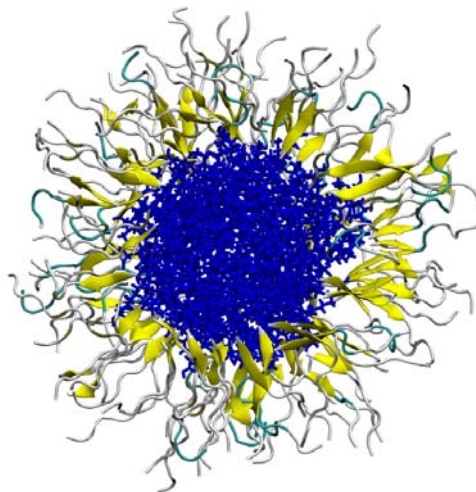


## Self-assembly stabilizes after a 40 ns simulation

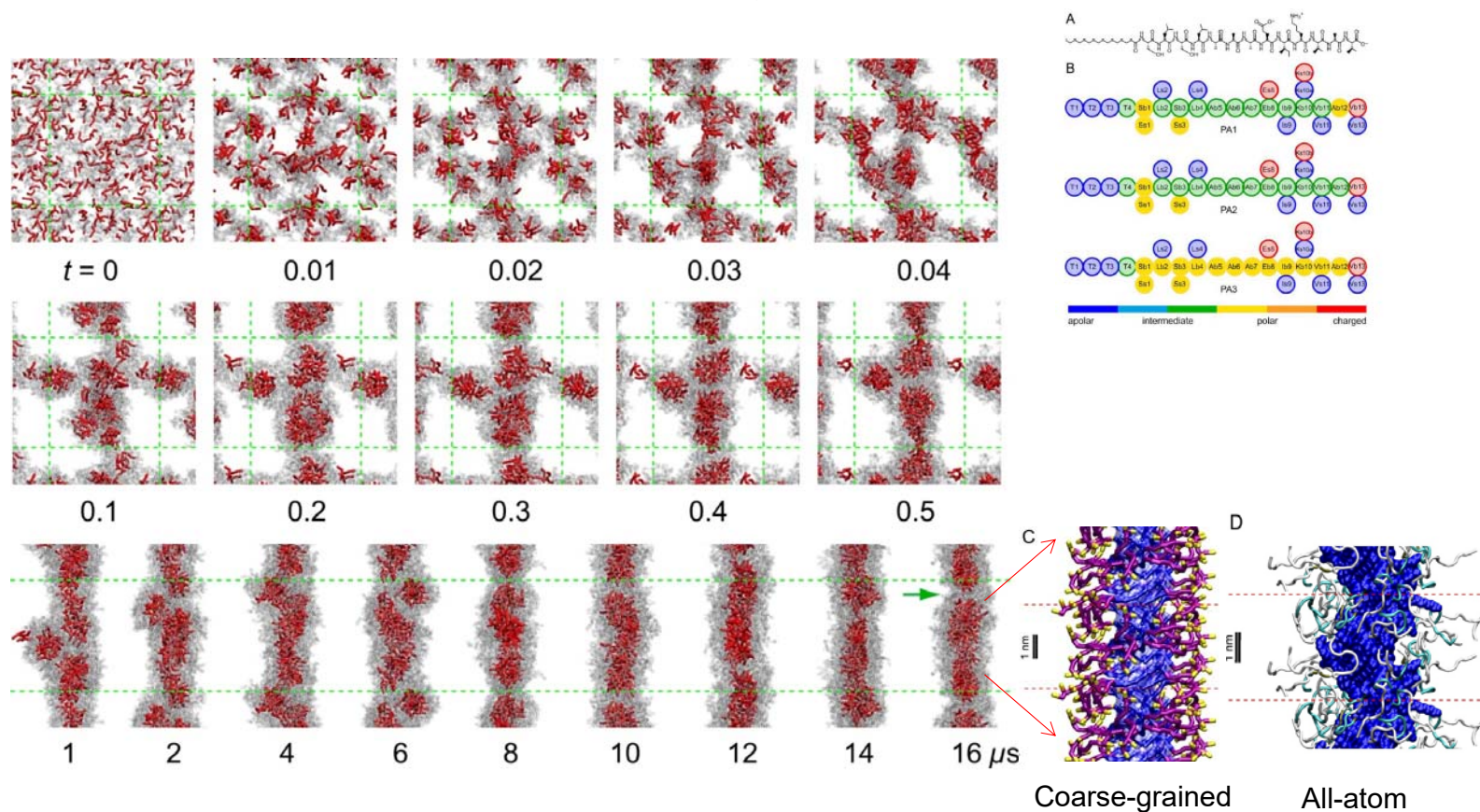


One-Sun Lee, Samuel I. Stupp and George C. Schatz, *J. Am. Chem. Soc.* 113, 3677-3683 (2011)

Aysenur Iscen and George C. Schatz 2017 *EPL* 119 38002

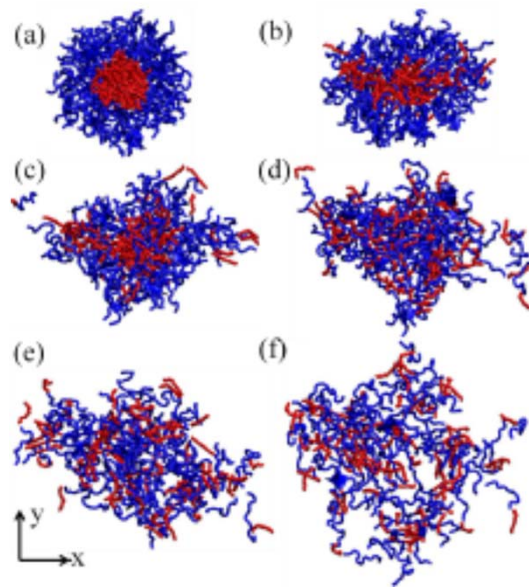


# Coarse-grained models enable simulation of the complete assembly process

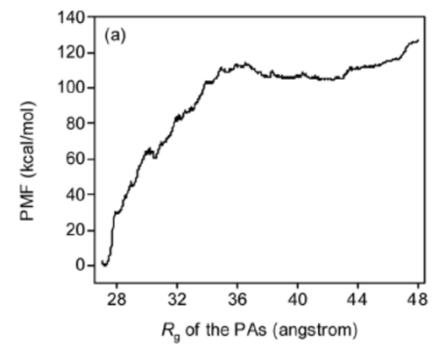


O.-S. Lee, V. Cho, G. C. Schatz, Nano Lett. 12, 4907-4913 (2012).

# Self-assembly mechanism



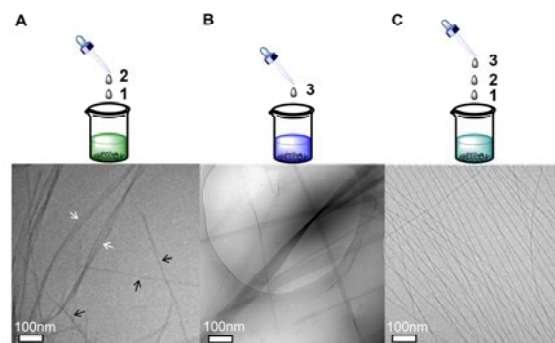
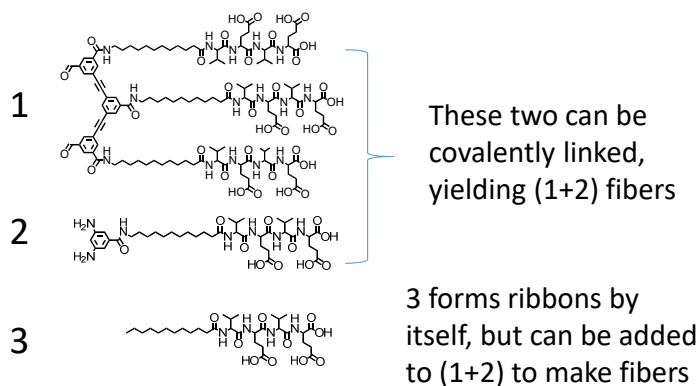
Potential of mean force vs radius of gyration



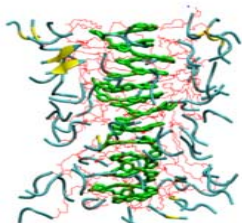


# Simultaneous Covalent and Noncovalent Hybrid Polymerizations

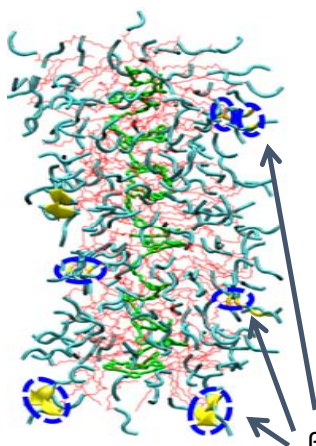
Z. Yu, F. Tantakitti, T. Yu, L. C. Palmer, G. C. Schatz and S. I. Stupp, Science, 351, 497, 2016



Theory confirms the structures, showing how beta sheets link (3) to (1+2)

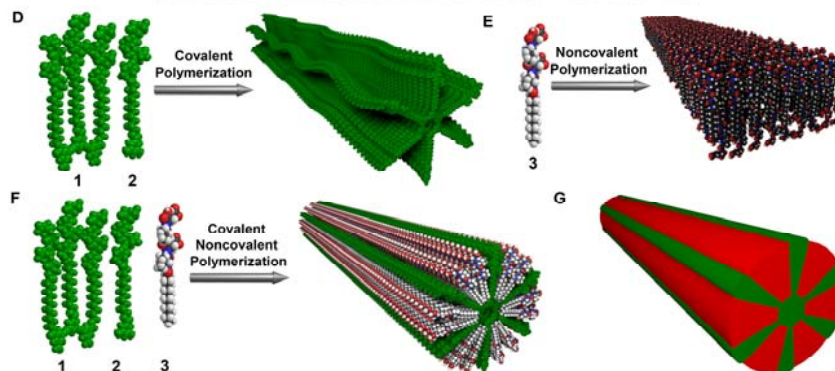


1+2



1+2+3

$\beta$  sheets linking 3 to 1+2



# A Mutation in Histone H2B Represents a New Class of Oncogenic Driver

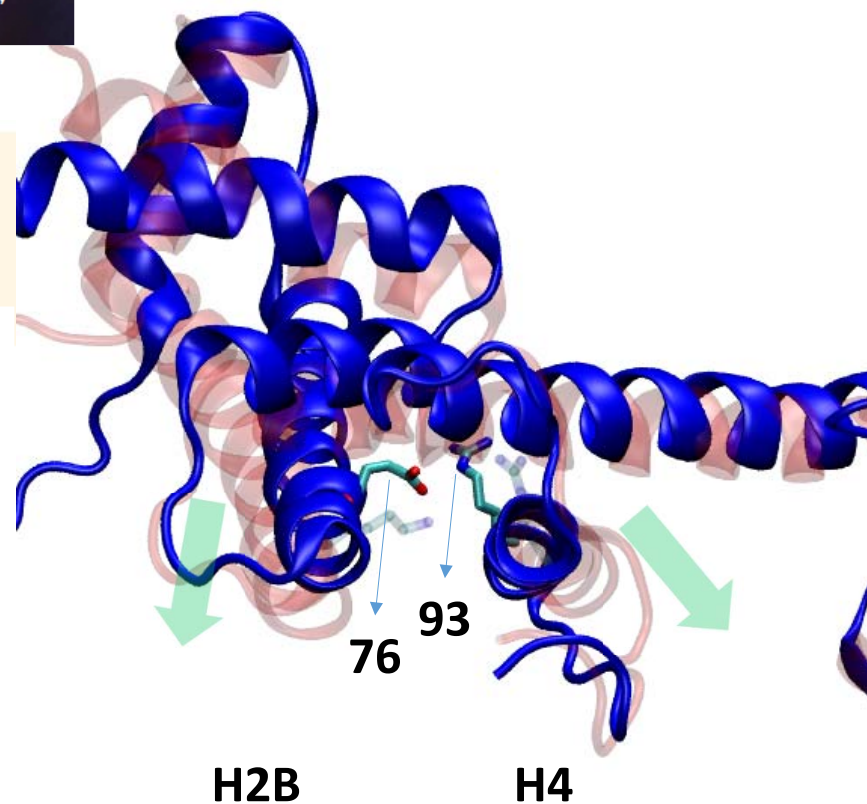
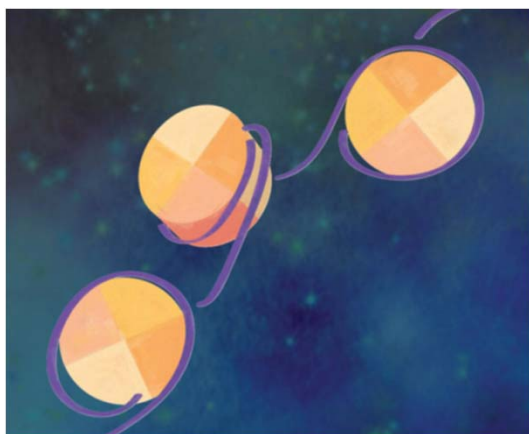


Richard L. Bennett<sup>1</sup>, Aditya Bele<sup>1</sup>, Eliza C. Small<sup>2</sup>, Christine M. Will<sup>2</sup>, Behnam Nabet<sup>3</sup>, Jon A. Oyer<sup>2</sup>, Xiaoxiao Huang<sup>1,4</sup>, Rajarshi P. Ghosh<sup>5</sup>, Adrian T. Grzybowski<sup>6</sup>, Tao Yu<sup>7</sup>, Qiao Zhang<sup>8</sup>, Alberto Riva<sup>9</sup>, Tanmay P. Lele<sup>8</sup>, George C. Schatz<sup>4</sup>, Neil L. Kelleher<sup>4</sup>, Alexander J. Ruthenburg<sup>6</sup>, Jan Liphardt<sup>5</sup>, and Jonathan D. Licht<sup>1</sup>

Cancer Discovery, 9, 1438-51 (2019)

## ABSTRACT

By examination of the cancer genomics database, we identified a new set of mutations in core histones that frequently recur in cancer patient samples and are predicted to disrupt nucleosome stability. In support of this idea, we characterized a glutamate to lysine mutation of histone H2B at amino acid 76 (H2B-E76K), found particularly in bladder and head and neck cancers, that disrupts the interaction between H2B and H4.



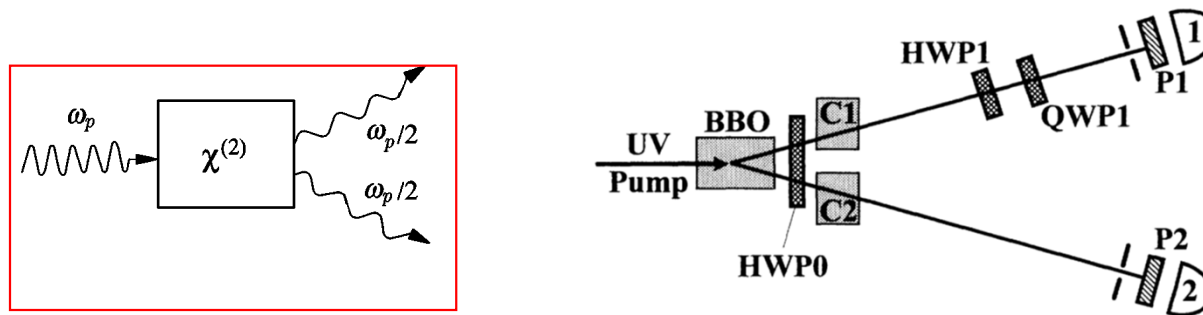
# Entangled two-photon absorption

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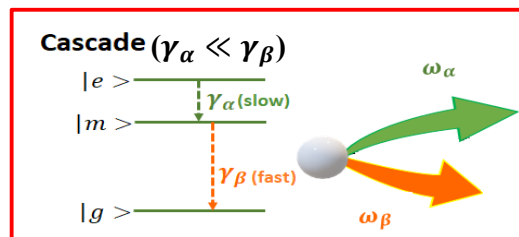
# Entangled photon sources

1. **Spontaneous parametric down conversion SPDC** (nonlinear crystal produces two photons)



$$|\text{twin}\rangle = Nl \int \int d\omega_1 d\omega_2 \exp\left[-\frac{(\omega_1 + \omega_2 - \omega_p)^2}{\Delta\omega_p^2}\right] \text{sinc}\left[\frac{l}{2\pi}(k_1 + k_2 - k_p)\right] |\omega_1, \omega_2\rangle.$$

2. **Quantum cascade** (emission by a doubly excited species)

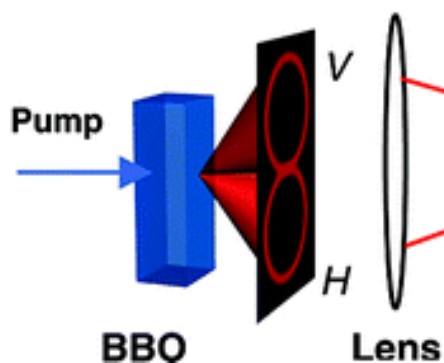




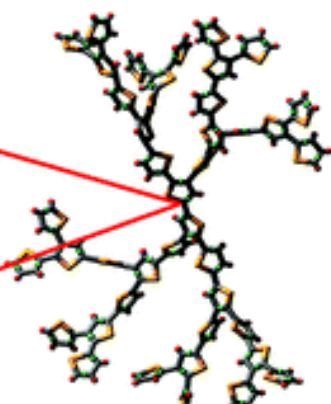
# Entangled Two-photon Absorption (ETPA)



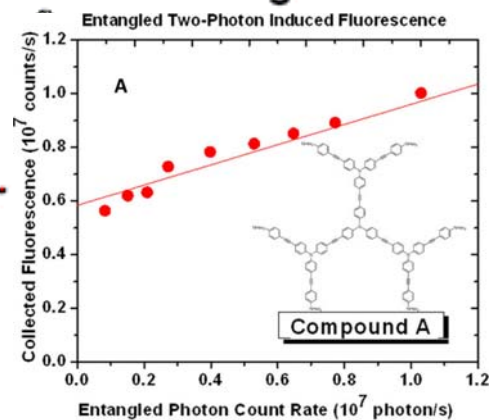
Entangled Photons  
from Downconversion



NLO Material



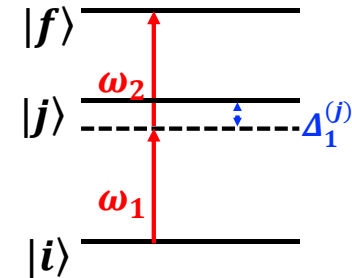
Entangled TPA



- ❑ Type II Spontaneous parametric down-conversion (SPDC) is used to generate entangled photon pairs ( $\vec{p}_1 \perp \vec{p}_2$ )
- ❑ Quadratic (Random TPA) vs Linear (ETPA) dependence on input flux

Goodson et al. *J. Am. Chem. Soc.* **2009**, 131, 973-979;  
Goodson, Mukamel, *J. Phys. Chem. Lett.* **2013**, 4, 2046

# Entangled and Unentangled two photon absorption



unentangled  
TPA rate:

$$\text{Rate} = \delta_R \Phi^2$$

$\Phi^2 = \text{square of flux (cm}^{-4}\text{s}^{-2}\text{)}$

Turns into a lineshape function  $g(\omega)$  when factoring in final state density

**TPA cross section:**

$$\delta_R = 2\pi^3 \frac{(\hbar r_0 c)^2}{(\hbar \omega)^2} \delta\left(\frac{E_f - E_i - 2\hbar\omega}{\hbar}\right) \left[ \sum_j \frac{2 \langle f | \mathbf{p}_x | j \rangle \langle j | \mathbf{p}_x | i \rangle}{m(E_j - E_i - \hbar\omega - \hbar\kappa_j / 2)} \right]^2 \quad (r_0 = e^2 / mc^2)$$

ETPA rate:

$$\text{Rate} = \sigma_e \Phi$$

**ETPA cross section:**

$$\sigma_e = \frac{\pi}{4A_e T_e} (2\pi)^2 \frac{(r_0 c)^2}{(\omega_1^0 \omega_2^0)} \delta\left(\frac{E_f - E_i - 2\hbar\omega}{\hbar}\right)$$

$A_e$ : Entanglement Area

$T_e$ : Entanglement Time

$$\times \left| \sum_j \left\{ \frac{\langle f | \mathbf{p}_{x,2} | j \rangle \langle j | \mathbf{p}_{x,1} | i \rangle}{m(E_j - E_i - \hbar\omega_1^0 - \hbar\kappa_j / 2)} (1 - e^{-iT_e(E_j - E_i - \hbar\omega_1^0) / \hbar - T_e \kappa_j / 2}) + \frac{\langle f | \mathbf{p}_{x,1} | j \rangle \langle j | \mathbf{p}_{x,2} | i \rangle}{m(E_j - E_i - \hbar\omega_2^0 - \hbar\kappa_j / 2)} (1 - e^{-iT_e(E_j - E_i - \hbar\omega_2^0) / \hbar - T_e \kappa_j / 2}) \right\} \right|^2$$

## Relationship between TPA and ETPA

TPA cross section:

$$\delta_R = 2\pi^3 \frac{(\hbar r_0 c)^2}{(\hbar \omega)^2} \delta\left(\frac{E_f - E_i - 2\hbar\omega}{\hbar}\right) \left[ \left| \sum_j \frac{2 \langle f | p_x | j \rangle \langle j | p_x | i \rangle}{m(E_j - E_i - \hbar\omega - \hbar\kappa_j/2)} \right|^2 \right]$$

ETPA rate:

$$\sigma_e = \frac{\pi}{4A_e T_e} (2\pi)^2 \frac{(r_0 c)^2}{(\omega_1^0 \omega_2^0)} \delta\left(\frac{E_f - E_i - 2\hbar\omega_p}{\hbar}\right) \times \left[ \sum_j \left\{ \frac{\langle f | p_{x,2} | j \rangle \langle j | p_{x,1} | i \rangle}{m(E_j - E_i - \hbar\omega_1^0 - \hbar\kappa_j/2)} \left(1 - e^{-iT_e(E_j - E_i - \hbar\omega_1^0)/\hbar - T_e\kappa_j/2}\right) + \frac{\langle f | p_{x,1} | j \rangle \langle j | p_{x,2} | i \rangle}{m(E_j - E_i - \hbar\omega_2^0 - \hbar\kappa_j/2)} \left(1 - e^{-iT_e(E_j - E_i - \hbar\omega_2^0)/\hbar - T_e\kappa_j/2}\right) \right\} \right]^2$$

If you ignore the oscillatory terms, then taking the ratio of these equations gives:

$$\sigma_e = \frac{\delta_R}{2T_e A_e}$$

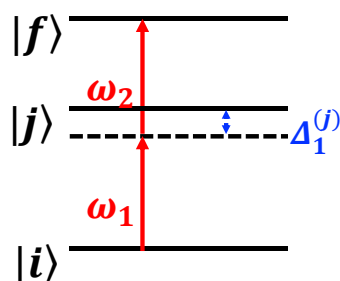
This means that ETPA is like TPA with a flux  $\Phi = 1/(2T_e A_e) = 5 \times 10^{18}$  photons/cm<sup>2</sup>/s for reasonable estimates of  $T_e A_e$

## ETPA Cross-section

SPDC sources have been used to study entangled two-photon absorption, leading to cross sections that are many orders of magnitude larger than for TPA.

$$\sigma_e = \frac{\pi}{4A_e T_e} \omega_1^0 \omega_2^0 \delta(\epsilon_f - \epsilon_i - \omega_1^0 - \omega_2^0) \quad \text{Goodson, Mukamel J. Phys. Chem. Lett. 2013, 4, 2046}$$

$$\times \left| \sum_j \left[ D_{21}^{(j)} \frac{1 - \exp[-iT_e \Delta_1^{(j)} - T_e \kappa_j/2]}{\Delta_1^{(j)} - i\kappa_j/2} + D_{12}^{(j)} \frac{1 - \exp[-iT_e \Delta_2^{(j)} - T_e \kappa_j/2]}{\Delta_2^{(j)} - i\kappa_j/2} \right] \right|^2$$



$D_{kl}^{(j)} = \langle f | \mu_k | j \rangle \langle j | \mu_l | i \rangle$  (transition matrix element)

$\omega_1 = \omega_2 = \omega_p/2$  (signal and idler photon energies)

$\Delta_k^{(j)} = \epsilon_j - \epsilon_i - \omega_k$  (energy mismatch)

$\kappa_j$ : Linewidth of state  $j$

These values are calculated with 2<sup>nd</sup> linear response TDDFT (SLR-TDDFT)

Martin A. Mosquera, Lin X. Chen, Mark A. Ratner, George C. Schatz, J. Chem. Phys. 144, 204105/1-11, (2016).



## Problem with reported values of $\sigma_e$ and $\delta_R$ given the “experimental parameters” $T_e$ and $A_e$

For spdc, previous estimates of  $T_e$  are usually about 0.1 ps, corresponding to the time it takes photons to traverse the BBO crystal.

$A_e$  has been estimated as  $10^{-6} \text{ cm}^2$ , as determined by the  $10 \text{ }\mu\text{m}$  diameter of the beam that irradiates the BBO crystal.

However if  $\sigma_e = 10^{-18} \text{ cm}^2$  (typical) and  $\delta_R = 100 \text{ GM (} \times 10^{-50} \text{ cm}^4 \text{ s)} = 10^{-48} \text{ cm}^4 \text{ s}$

and 
$$A_e = \frac{\delta_R}{2\sigma_e T_e}$$

Then this predicts  $A_e = 10^{-48} / (10^{-13} 10^{-18}) = 10^{-17} \text{ cm}^2$

Ouch!, off by  $10^{11}$

## Problem can be fixed by realizing that the Dirac Delta functions in the TPA and ETPA expressions don't refer to the same state density

In this case, the delta functions become different lineshape functions:

$$\begin{aligned}\delta(2\omega - \omega_0) &\rightarrow g_R && \text{TPA} \\ \delta(2\omega - \omega_0) &\rightarrow g_e && \text{ETPA}\end{aligned}$$

that describe the density of states for entangled (E) and unentangled (R) states of the molecule in the presence of the radiation field.

If the lineshape is Lorentzian:

$$g_{e,R} = \frac{1}{\pi} \frac{\Gamma_{e,R}}{(2\omega - \omega_0)^2 + \Gamma_{e,R}^2}$$

where  $\Gamma_{e,R}$  is the HWHM of the lineshape for unentangled and entangled absorption. Typically it is only the peak cross section that is studied, in which case

$$g_{e,R} = \frac{1}{\pi\Gamma_{e,R}} = \frac{2}{\pi}\tau$$

## Calculating TPA (unentangled cross sections)

There are many papers but a recent paper that provides exhaustive details is:



Cite this: *Phys. Chem. Chem. Phys.*,  
2015, 17, 19306

### Benchmarking two-photon absorption cross sections: performance of CC2 and CAM-B3LYP

Maarten T. P. Beerepoot,<sup>\*a</sup> Daniel H. Friese,<sup>a</sup> Nanna H. List,<sup>b</sup> Jacob Kongsted<sup>b</sup> and Kenneth Ruud<sup>a</sup>

$$\sigma^{\text{TPA}} = \frac{4\pi^3 \alpha a_0^5 \omega^2}{15c} \langle \delta^{\text{TPA}} \rangle g(2\omega, \omega_0, \Gamma). \quad \langle \delta^{\text{TPA}} \rangle = \frac{1}{30} \sum_{ab} (FS_{aa}S_{bb} + GS_{ab}S_{ab} + HS_{ab}S_{ba}),$$

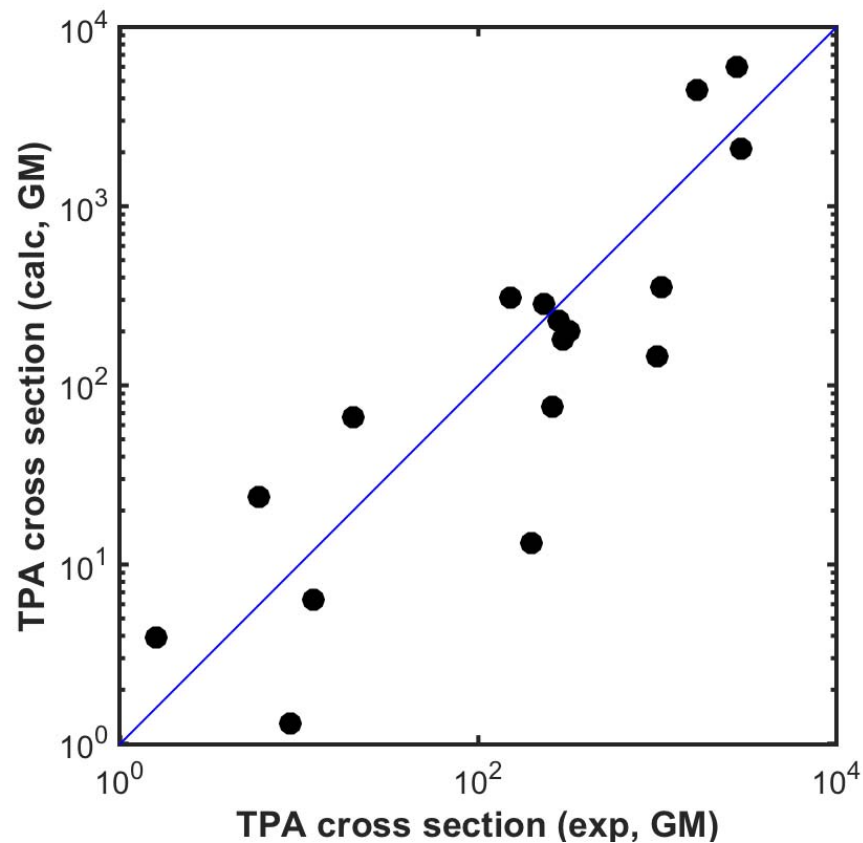
$$S_{ab}^{ij}(\omega_1, \omega_2) = \frac{1}{\hbar} \sum_{n \neq i} \left\{ \frac{\langle i|\mu_a|n\rangle \langle n|\mu_b|f\rangle}{\omega_{ni} - \omega_1} + \frac{\langle i|\mu_b|n\rangle \langle n|\mu_a|f\rangle}{\omega_{ni} - \omega_2} \right\}$$

The broadening effects are different for each excited state,<sup>26</sup> which can also be taken into account in theoretical work.<sup>30,32,36,39</sup>

What is usually done, however, is choosing a single empirical parameter for  $\Gamma$ , often chosen to be 0.1 eV,<sup>13,14,18,21,22,28,29,37,38</sup>

but other broadening functions have been used, such as the Lorentzian

Our calculations (width = 0.1 eV) give results that are consistent with a factor of 10 variation between theory and expt for TPA

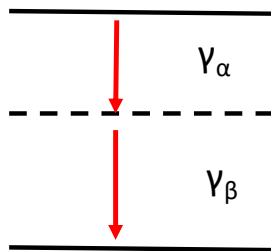


	exp, GM	calc, GM
6-Thiophene	6	24
18-Thiophene	230	288
ZnTPP	20	66
T161B	280	228
T161D	1040	358
bisstyrylbenzene	300	182
Bisannulene	150	305
Triannulene	1650	4446
Tetraannulene	2960	2105
Flavin		
Mononucleotide	1.6	3.9
R6G	9	1.3
stilbene1	12	6.3
stilbene2	200	13
stilbene3	995	144
stilbene12	260	76
stilbene13	320	199
dibenzoheptazethrene	2800	5959



## What is the width or lifetime that is relevant to ETPA?

If only radiative effects are important, then we need the *radiative* lifetime of the two photon state. This can be calculated if we know the excited state properties (energies and transition moments to lower state). Typically it leads to lifetimes of  $\mu\text{s}$  to  $\text{ms}$  that are dominated by the lifetime for transition to the dominant intermediate state ( $\Gamma = \gamma_\alpha$ ).



$$\tau_r = \left[ \sum_{j < f} \frac{4}{3\hbar} \left( \frac{|\omega_{jf}|}{c} \right)^3 |\mu_{jf}|^2 \right]^{-1}$$

Further, if  $\gamma_\alpha \ll \gamma_\beta$ , photon is highly entangled, with a big Schmidt number.

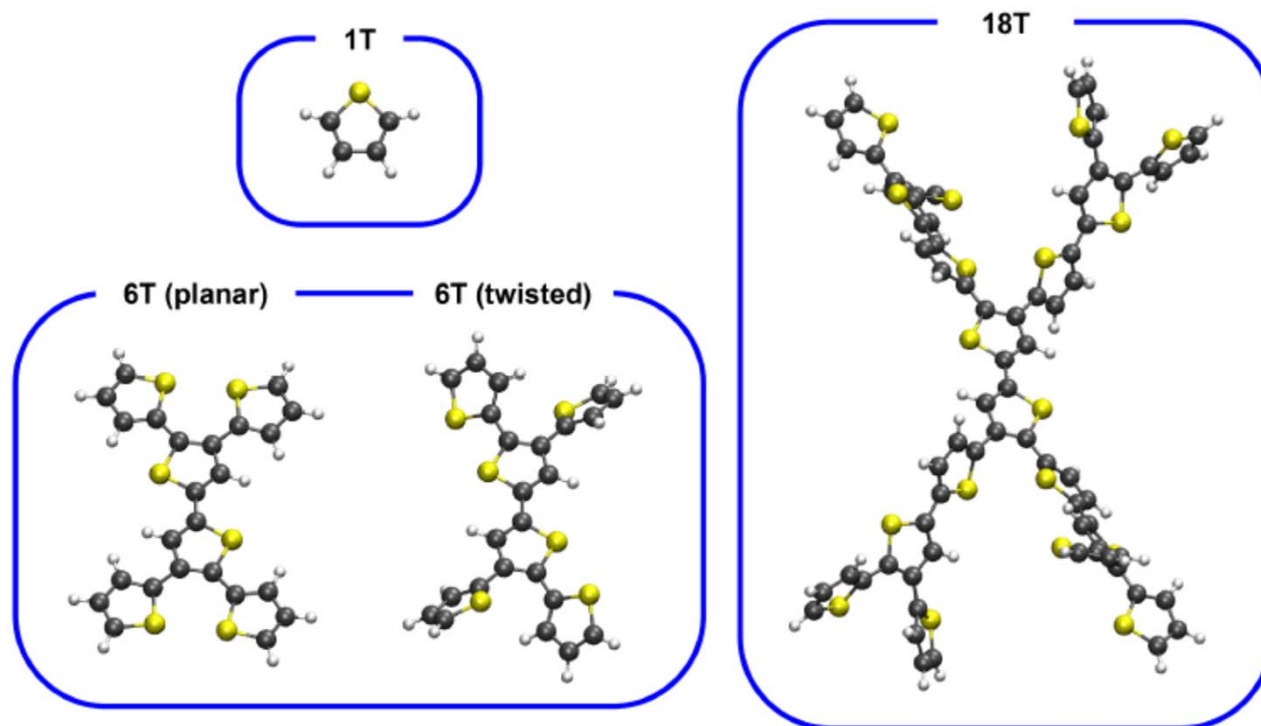
### Criticisms:

1. Assumes that entangled and unentangled states are not coupled. Are there proofs of this? There are theorems related to entanglement monogamy/faithfulness. Also, there are experiments (see later)
2. Only considers radiative contribution
3. Not clear that the two photon lifetime is measurable

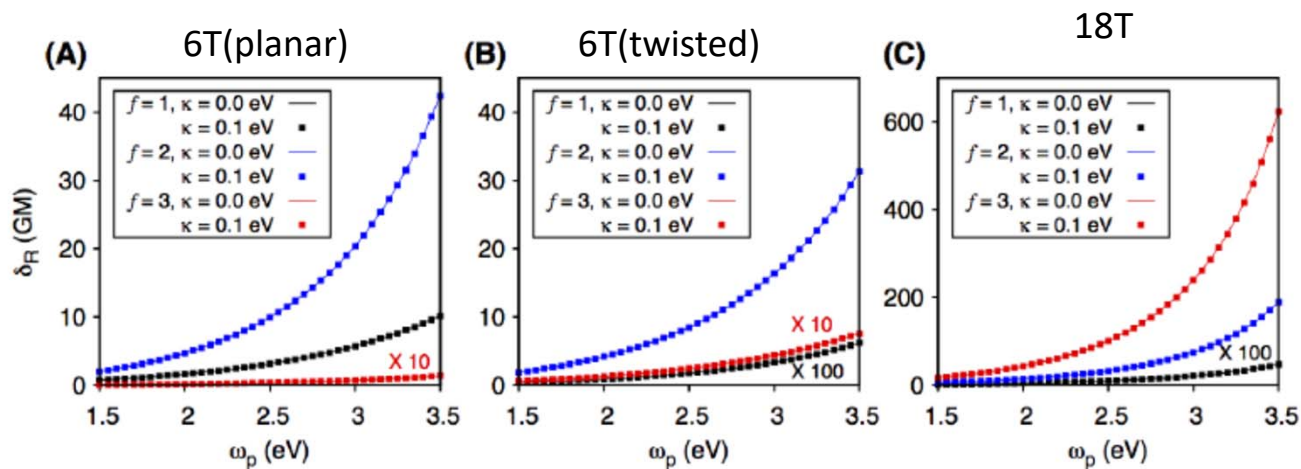
“geometrically faithful” entanglements quantitatively capture the fact that entangled states are very different than unentangled states.

# Application of ETPA theory that uses SLR to calculate cross section including for radiative lifetime.

Thiophene dendrimers



## Calculated TPA cross sections



Molecule	Experiment <sup>20</sup>	Theory	
	$\delta_R$ ( $10^{-50} \text{ cm}^4 \text{ s}$ )	$f$	$\delta_R$ ( $10^{-50} \text{ cm}^4 \text{ s}$ )
6T (planar)	6	ES2	24
6T (twisted)		ES2	19
18T	230	ES2	90
		ES3	288

G. Kang, K. Avanaki, M. Mosquera, R. K. Burdick, J. Monsalve, T. Goodson, G. Schatz, JACS142, 10446-10458 (2020).

Two possible choices of the 2-photon excited state



Goodson *et al.* *J. Am. Chem. Soc.* 131, 3, 2009

## ETPA cross sections

Molecule	Experiment <sup>20</sup>	Theory			
	$\sigma_e$ ( $10^{-19}$ cm <sup>2</sup> )	$f$	$\sigma_e$ ( $10^{-19}$ cm <sup>2</sup> )	$g_e$ ( $\mu$ s)	$\gamma_f/\gamma_1$
6T (planar)	1.3	ES2	1.7	988	$6 \times 10^{-4}$
6T (twisted)		ES2	0.18	175	$3 \times 10^{-3}$
18T	7.1	ES2	5.3	853	$6 \times 10^{-4}$
		ES3	1.2	84	$6 \times 10^{-3}$



Goodson *et al.* *J. Am. Chem. Soc.* 131, 3, **2009**



G. Kang, K. Avanaki, M. Mosquera,  
R. K. Burdick, J. Monsalve, T.  
Goodson, G. Schatz, *JACS* 142,  
10446-10458 (2020).

## Conclusions

1. We can calculate TPA cross sections that are good to a factor of 10 using 0.1 eV for the width
2. ETPA cross sections are systematically shifted by a factor of  $\sim 10$ , and then the factor of 10 uncertainty also applies.