

Many Body Quantum Dynamics *of Exciton-Polaritons*

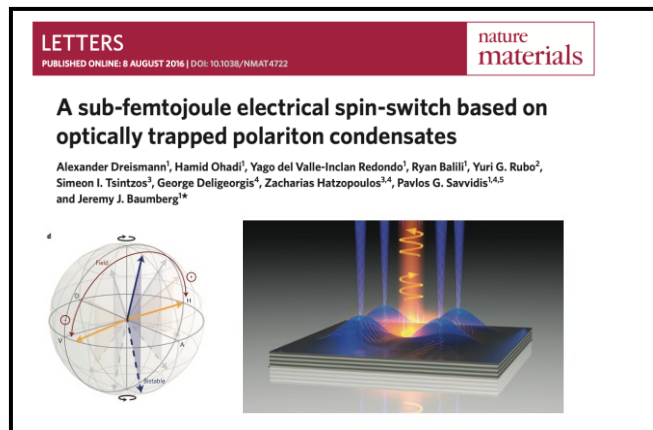
Pritha Ghosh

Mandal Group, *Department of Chemistry*

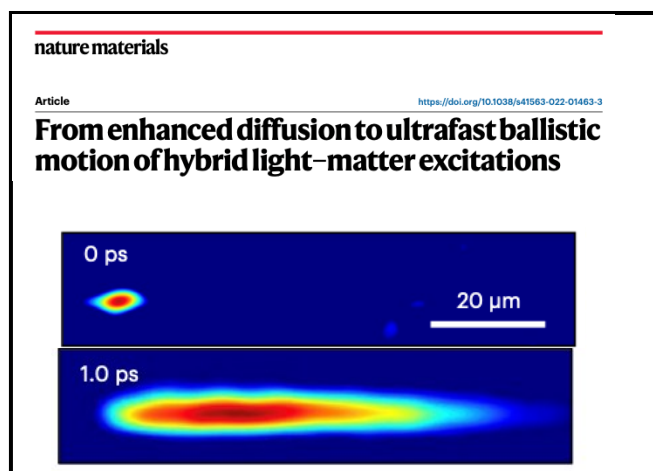
Ghosh, Manjalingal, Wickramasinghe, Rahmanian, Mandal, *Phys. Rev. B.* (2025)
Editors' Suggestion

Motivation

Polaritons Exhibit Room-temperature Quantum Coherence!

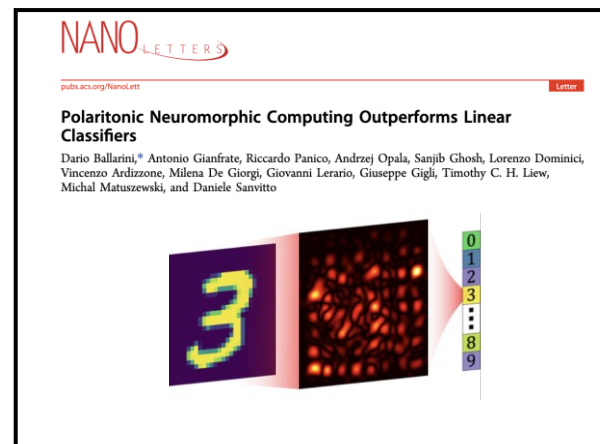


Room-temperature Bose-Einstein Condensate.

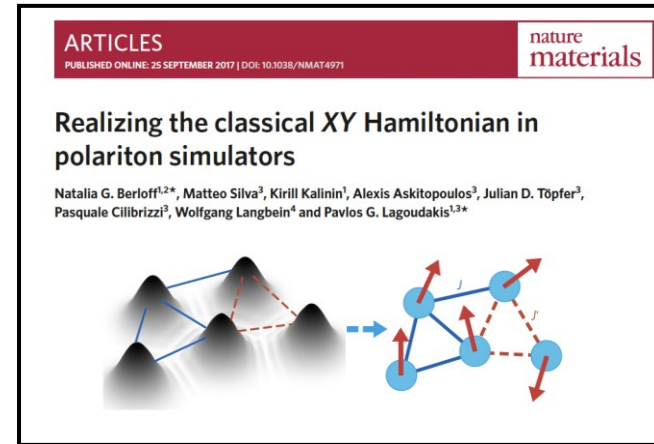


Room temperature Coherent Propagation of Excitons.

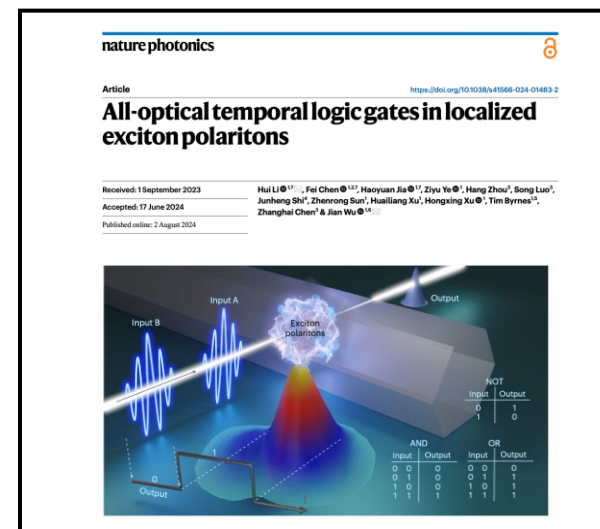
Diverse Applications!



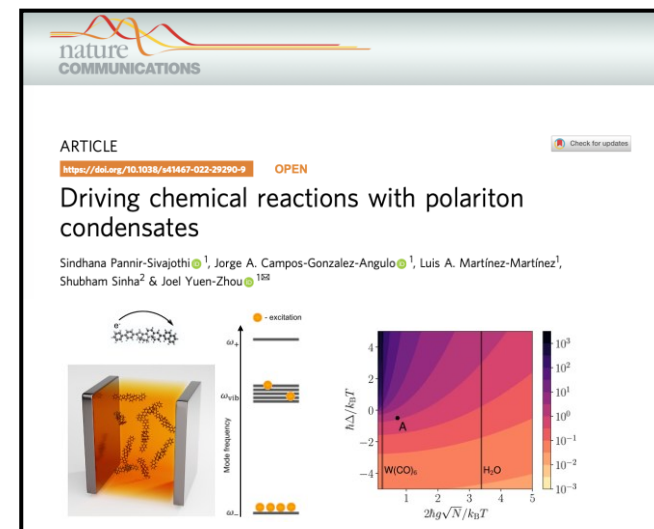
Efficient Neuromorphic Computing



Platform for Analogue Quantum Computing



Exciton-polariton Based Logic Gates



Exotic Chemistry in Cavities

Theoretical Challenges

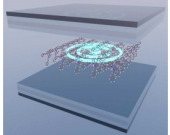
Holstein Tavis-Cummings

NANO LETTERS

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Quantum Dynamics Simulations of Exciton Polariton Transport

Benjamin X. K. Chng, M. Elious Mondal, Wenxiang Ying, and Pengfei Huo*

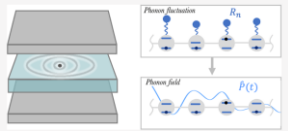


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Microscopic Theory of Polaron-Polariton Dispersion and Propagation

Logan Blackham, Arshath Manjalingal, Saeed Rahmanian Koshkaki, and Arkajit Mandal*



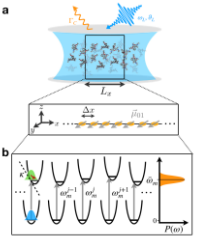
nature communications

Article <https://doi.org/10.1038/s41467-025-61298-9>

Quantum dynamics simulation of exciton-polariton transport

Received: 30 October 2024 | Accepted: 18 June 2025

Niclas Krupp¹, Gerrit Groenhof² & Oriol Vendrell¹*




Explicit Phonon Interaction



Many body interactions



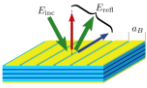
Gross-Pitaevskii

PHYSICAL REVIEW B **100**, 035305 (2019)

Generalized Gross-Pitaevskii model for intersubband polariton lasing

Jacopo Nespolo* and Jacopo Carusotto†

INO-CNR BEC Center and Dipartimento di Fisica, Università di Trento, I-38123 Povo, Italy



IOP Publishing | London Mathematical Society

Nonlinearity **32** (2019) 4317–4345

<https://doi.org/10.1088/1361-6544/ab2bc1>

On a dissipative Gross-Pitaevskii-type model for exciton-polariton condensates*

Paolo Antonelli¹, Peter Markowich², Ryan Obermeyer^{1,3}, Jesus Sierra¹ and Christof Sparber²



PHYSICAL REVIEW B **89**, 155302 (2014)

Quantum kinetic derivation of the nonequilibrium Gross-Pitaevskii equation for nonresonant excitation of microcavity polaritons

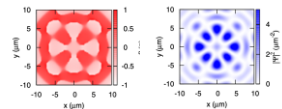
H. Haug

Institut für Theoretische Physik, Goethe Universität Frankfurt, Max-von-Laue-Straße 1, D-60438 Frankfurt a.M., Germany

T. D. Doan and D. B. Tran Thoai

Ho Chi Minh City Institute of Physics, Vietnam Center for Natural Science and Technology, 1 Mac Dinh Chi, Ho Chi Minh City, Vietnam

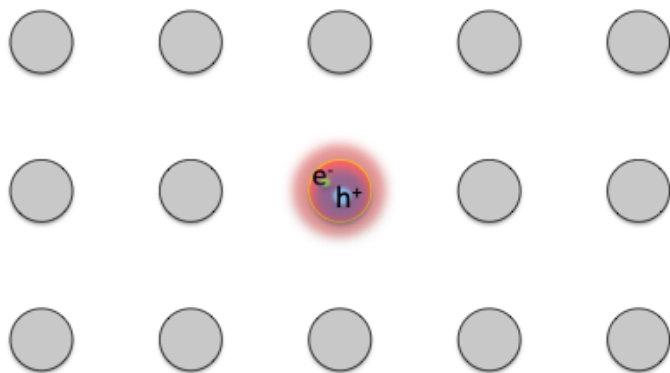
(Received 29 October 2013; revised manuscript received 17 February 2014; published 2 April 2014)



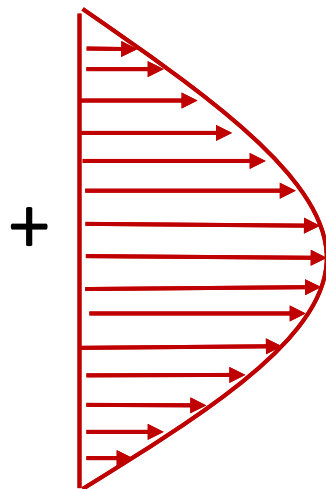
Simulating polaritons at room temperature → Mesoscale systems + many-body interactions + coupling to phonons

Exciton-polaritons

A single Frenkel exciton in an organic crystal

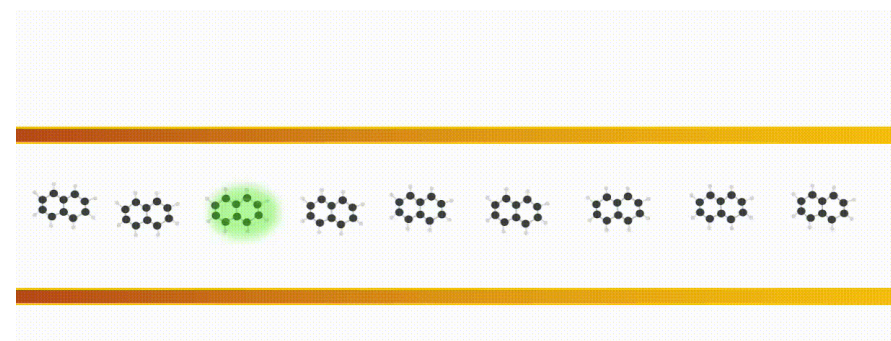
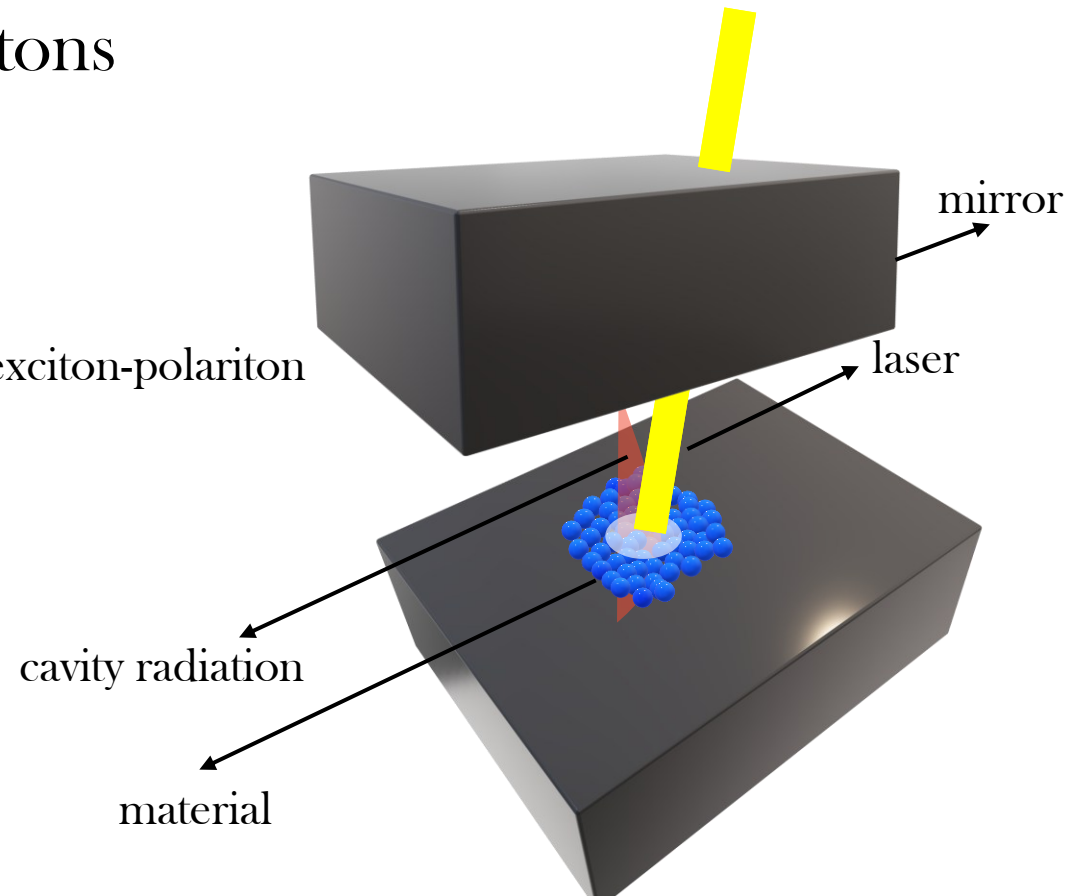


Quantized cavity radiation mode



+

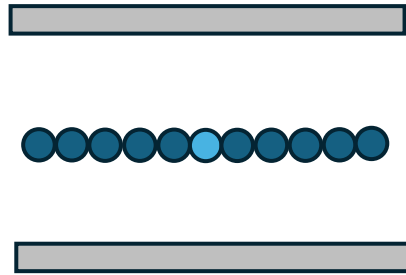
exciton-polariton



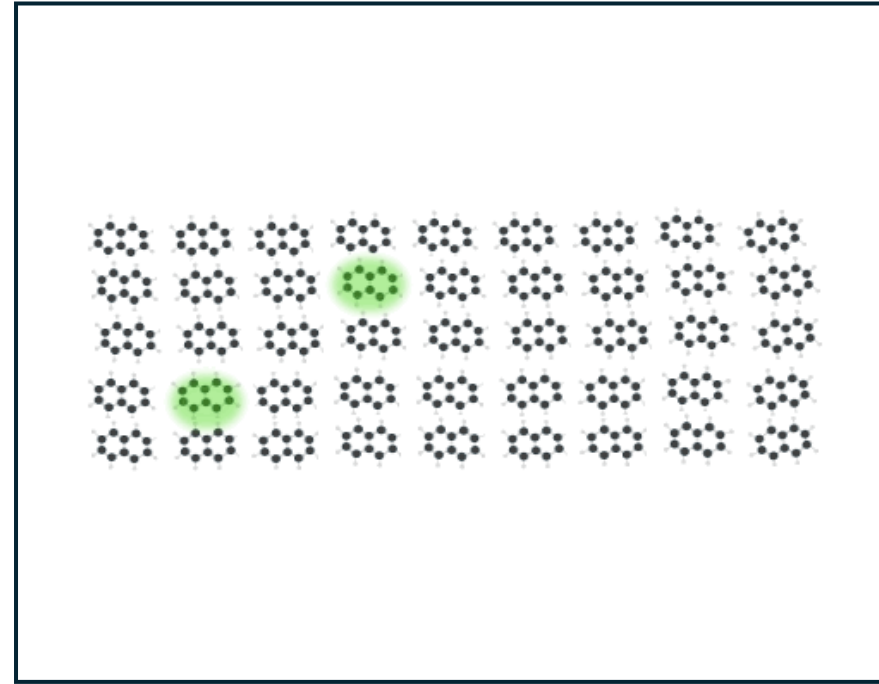
In a bare material, exciton shows (incoherent) diffusive transport.

In a cavity, exciton-polaritons show (coherent) ballistic transport.

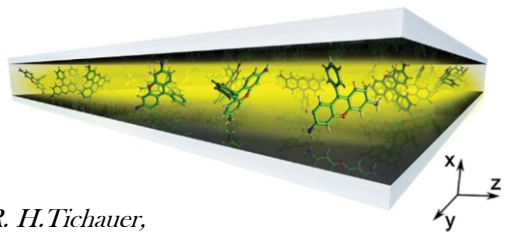
INCOMPATIBLE with experiments!



$N_{\text{ex}}=1$
Single excited subspace

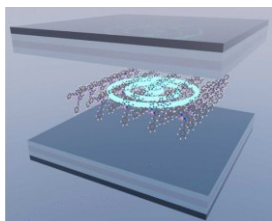


For $N_{\text{ex}} \gg 1$, many-body interactions have to be included \rightarrow dynamics become **complicated!**



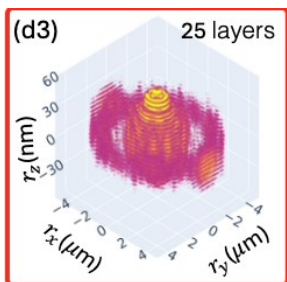
R. H. Tichauer,
I. Sokolovskii, G. Groenhof,
Adv. Sci. 2023

~ 500 molecules



Chng, *Nano Lett.* 2025

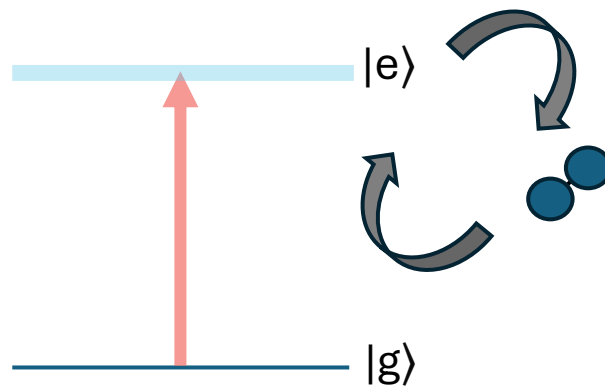
~ 10^4 molecules



> 10^6 molecules

Koshkaki,
arXiv:2502.12933v2

Mixed quantum classical simulation



Multitrajectory
Ehrenfest

Fewest Switches
Surface Hopping

Path Integral
Based Approach

Huge numbers of classical degrees of freedom

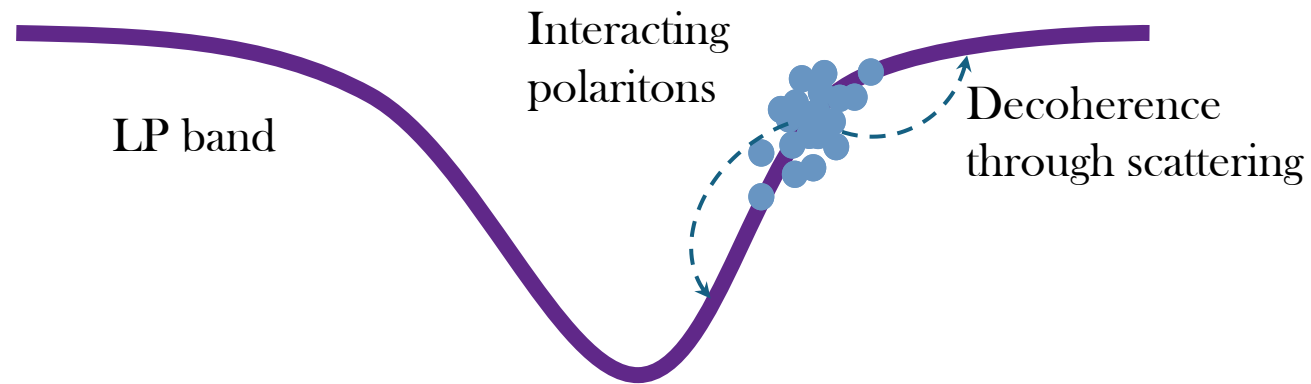
Computationally intractable!

A novel approach

Multitrajectory mean field **Ehrenfest** for the phonons (quasi-classical)

+

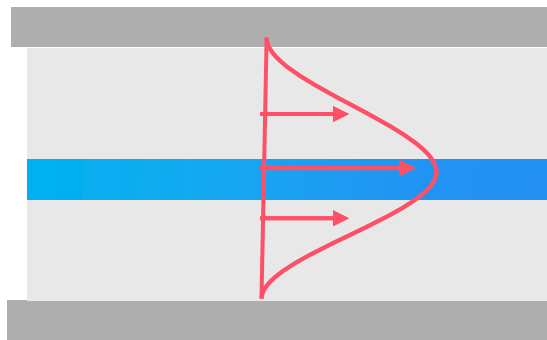
Mean field **Gross-Pitaevskii** for the exciton-polaritons (quantum)



Light-Matter Hamiltonian

$$\hat{H}_{\text{LM}} = \hat{H}_e + \hat{H}_p + \hat{H}_{e-p} + \hat{H}_c + \hat{H}_{e-c}.$$

$$= \sum_{n,m} \epsilon_{n,m} \hat{X}_n^\dagger \hat{X}_m + \frac{U}{2} \sum_n \hat{X}_n^\dagger \hat{X}_n^\dagger \hat{X}_n \hat{X}_n + \sum_n \frac{\hat{p}_n^2}{2} + \frac{1}{2} \omega^2 \hat{R}_n^2 + \gamma \sum_n \hat{X}_n^\dagger \hat{X}_n \hat{R}_n + \sum_k \hat{a}_k^\dagger \hat{a}_k \omega_k + \sum_{n,k} \frac{\Omega_k}{\sqrt{N}} \left[\hat{a}_k^\dagger \hat{X}_n e^{-ik \cdot r_n} + \hat{a}_k \hat{X}_n^\dagger e^{ik \cdot r_n} \right].$$

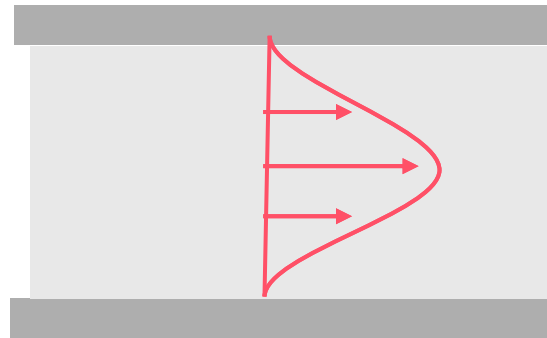


Light-Matter Hamiltonian

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Material
Excitation
(quantum)

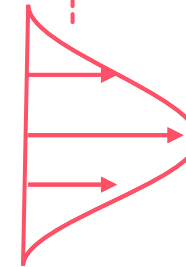
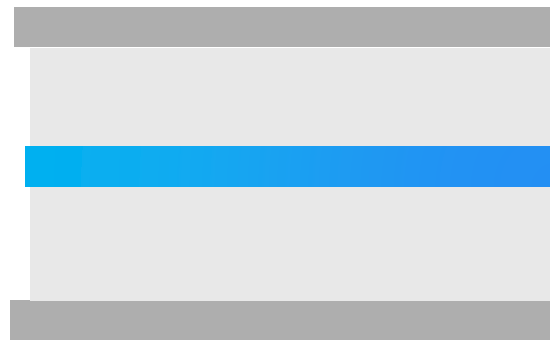


Light-Matter Hamiltonian

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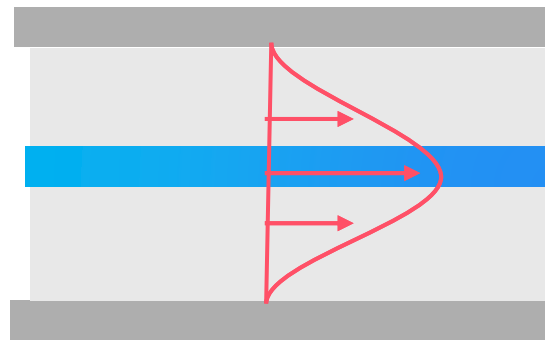
Cavity
Photons
(quantum)



Light-Matter Hamiltonian

$$\hat{H}_{\text{LM}} = \hat{H}_e + \hat{H}_p + \hat{H}_{e-p} + \hat{H}_c + \hat{H}_{e-c}.$$

$$= \sum_{n,m} \epsilon_{n,m} \hat{X}_n^\dagger \hat{X}_m + \frac{U}{2} \sum_n \hat{X}_n^\dagger \hat{X}_n^\dagger \hat{X}_n \hat{X}_n + \sum_n \frac{\hat{p}_n^2}{2} + \frac{1}{2} \omega^2 \hat{R}_n^2 + \gamma \sum_n \hat{X}_n^\dagger \hat{X}_n \hat{R}_n + \sum_k \hat{a}_k^\dagger \hat{a}_k \omega_k + \sum_{n,k} \frac{\Omega_k}{\sqrt{N}} \left[\hat{a}_k^\dagger \hat{X}_n e^{-ik \cdot r_n} + \hat{a}_k \hat{X}_n^\dagger e^{ik \cdot r_n} \right]$$



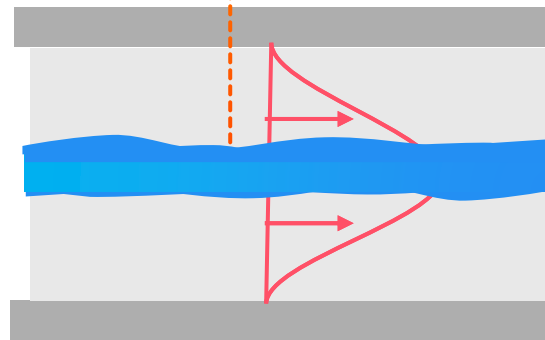
Light Matter Interactions

Light-Matter Hamiltonian

$$\hat{H}_{\text{LM}} = \hat{H}_e + \hat{H}_p + \hat{H}_{e-p} + \hat{H}_c + \hat{H}_{e-c}.$$

$$= \sum_{n,m} \epsilon_{n,m} \hat{X}_n^\dagger \hat{X}_m + \frac{U}{2} \sum_n \hat{X}_n^\dagger \hat{X}_n^\dagger \hat{X}_n \hat{X}_n + \sum_n \left[\frac{\hat{P}_n^2}{2} + \frac{1}{2} \omega^2 \hat{R}_n^2 + \gamma \sum_n \hat{X}_n^\dagger \hat{X}_n \hat{R}_n \right] + \sum_k \hat{a}_k^\dagger \hat{a}_k \omega_k + \sum_{n,k} \frac{\Omega_k}{\sqrt{N}} \left[\hat{a}_k^\dagger \hat{X}_n e^{-ik \cdot r_n} + \hat{a}_k \hat{X}_n^\dagger e^{ik \cdot r_n} \right].$$

Phonon interactions
(quasi-classical)

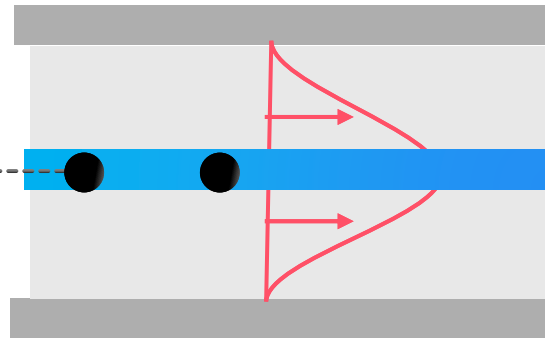


Light-Matter Hamiltonian

$$\hat{H}_{\text{LM}} = \hat{H}_e + \hat{H}_p + \hat{H}_{e-p} + \hat{H}_c + \hat{H}_{e-c}.$$

$$= \sum_{n,m} \epsilon_{n,m} \hat{X}_n^\dagger \hat{X}_m + \boxed{\frac{U}{2} \sum_n \hat{X}_n^\dagger \hat{X}_n^\dagger \hat{X}_n \hat{X}_n} + \sum_n \frac{\hat{P}_n^2}{2} + \frac{1}{2} \omega^2 \hat{R}_n^2 + \gamma \sum_n \hat{X}_n^\dagger \hat{X}_n \hat{R}_n + \sum_k \hat{a}_k^\dagger \hat{a}_k \omega_k + \sum_{n,k} \frac{\Omega_k}{\sqrt{N}} \left[\hat{a}_k^\dagger \hat{X}_n e^{-ik \cdot r_n} + \hat{a}_k \hat{X}_n^\dagger e^{ik \cdot r_n} \right].$$

Exciton-Exciton
interactions (**on-site**)



Choice of Ansatz

$$\hat{H}_{\text{LM}} = \hat{H}_e + \hat{H}_p + \hat{H}_{e-p} + \hat{H}_c + \hat{H}_{e-c}.$$

$m=n \rightarrow$ on-site energy (ϵ_0)
 $m \neq n \rightarrow$ hopping (τ)

$$= \sum_{n,m} \epsilon_{n,m} \hat{X}_n^\dagger \hat{X}_m + \frac{U}{2} \sum_n \hat{X}_n^\dagger \hat{X}_n^\dagger \hat{X}_n \hat{X}_n.$$

On-site many body interactions

$$+ \sum_n \frac{\hat{P}_n^2}{2} + \frac{1}{2} \omega^2 \hat{R}_n^2 + \gamma \sum_n \hat{X}_n^\dagger \hat{X}_n \hat{R}_n$$

$$+ \sum_k \hat{a}_k^\dagger \hat{a}_k \omega_k + \sum_{n,k} \frac{\Omega_k}{\sqrt{N}} \left[\hat{a}_k^\dagger \hat{X}_n e^{-ik \cdot r_n} + \hat{a}_k \hat{X}_n^\dagger e^{ik \cdot r_n} \right].$$

Mixed Quantum Classical Dynamics

$$\dot{P}_n(t) = - \left\langle \Psi(t) \left| \frac{d\hat{H}_{\text{LM}}}{dR_n} \right| \Psi(t) \right\rangle, \quad \dot{R}_n(t) = P_n(t)$$

$$i\dot{|\Psi(t)\rangle} = \left[\hat{H}_{\text{LM}} - \frac{1}{2} \sum_n \frac{P_n^2}{2} \right] |\Psi(t)\rangle$$

Intractable

Mean-field Ansatz

$$|\Psi(t)\rangle = \frac{(\hat{B}_0^\dagger(t))^{N_{\text{ex}}}}{\sqrt{N_{\text{ex}}!}} |0\rangle.$$

$$\hat{B}_0^\dagger(t) = \sum_n \phi_n(t) \cdot \hat{X}_n^\dagger + \sum_k \phi_k(t) \cdot \hat{a}_k^\dagger$$

EOM

Finally, using the Dirac-Frenkel Variational Principle,

**EOM of the exciton-photon
quantum subsystem**

$$i\dot{\psi}_n(t) = \epsilon_0\psi_n(t) - \tau(\psi_{n+1}(t) + \psi_{n-1}(t)) + \gamma\psi_n(t)R_n(t) + \sum_k \frac{\Omega_k}{\sqrt{N}} \phi_k(t) e^{-ik \cdot r_n} \\ + U(N_{\text{ex}} - 1)|\psi_n(t)|^2\psi_n(t).$$

$$i\dot{\phi}_k(t) = \phi_k(t)\omega_k + \sum_n \frac{\Omega_k}{\sqrt{N}} \psi_n(t) e^{-ik \cdot r_n}$$

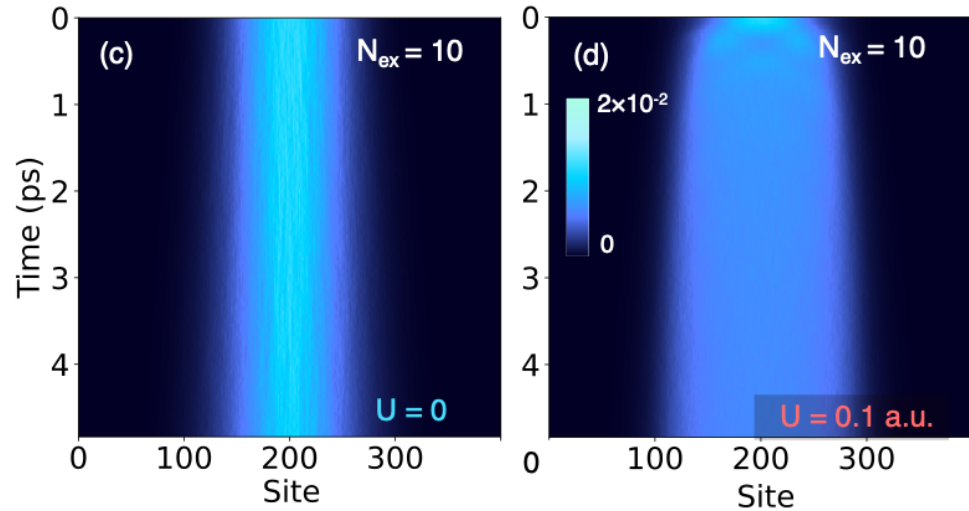
EOM of the phonons

$$\dot{R}_n(t) = P_n(t)$$

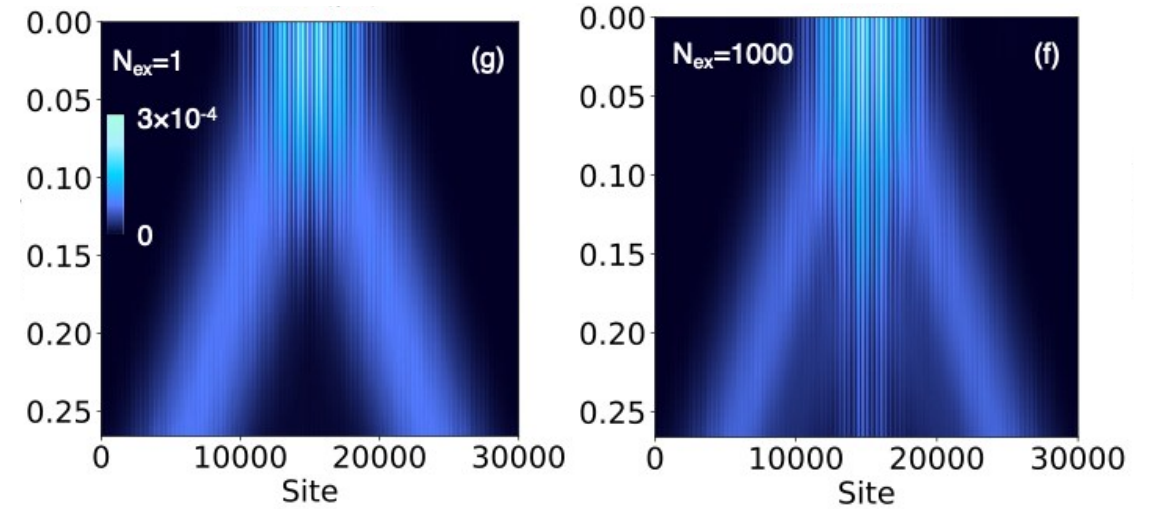
$$\dot{P}_n(t) = -\omega^2 R_n - \gamma|\psi_n(t)|^2 N_{\text{ex}}$$

Results

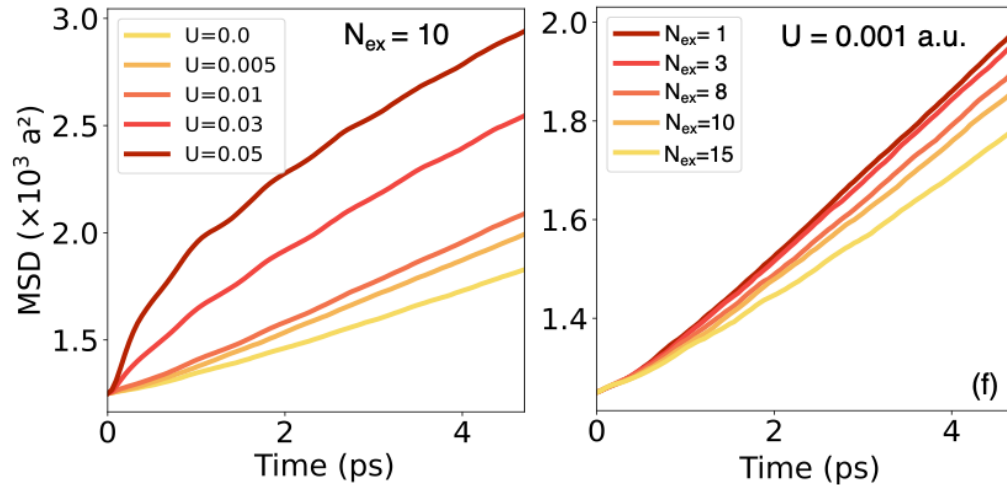
Quantum Dynamics of Bare Excitons



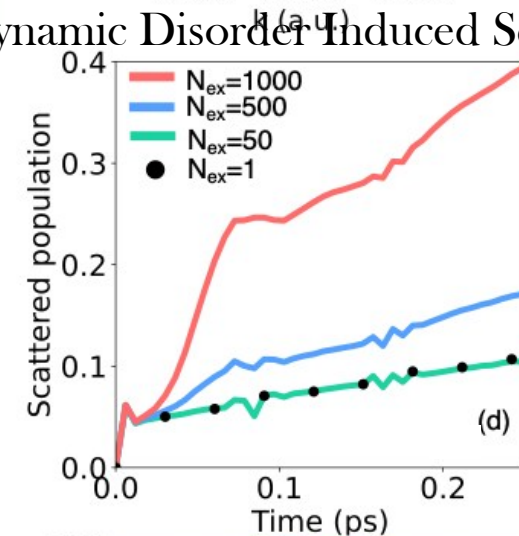
Quantum Dynamics of Exciton-polaritons (Non-interacting case)



Mean Square Displacement

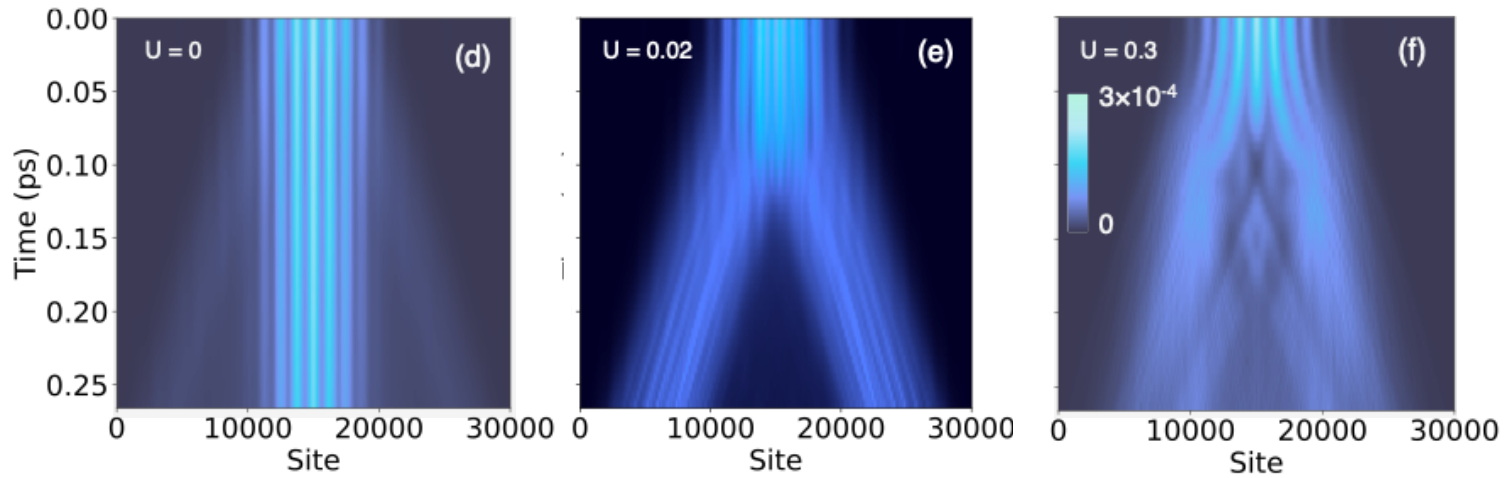


Dynamic Disorder Induced Scattering

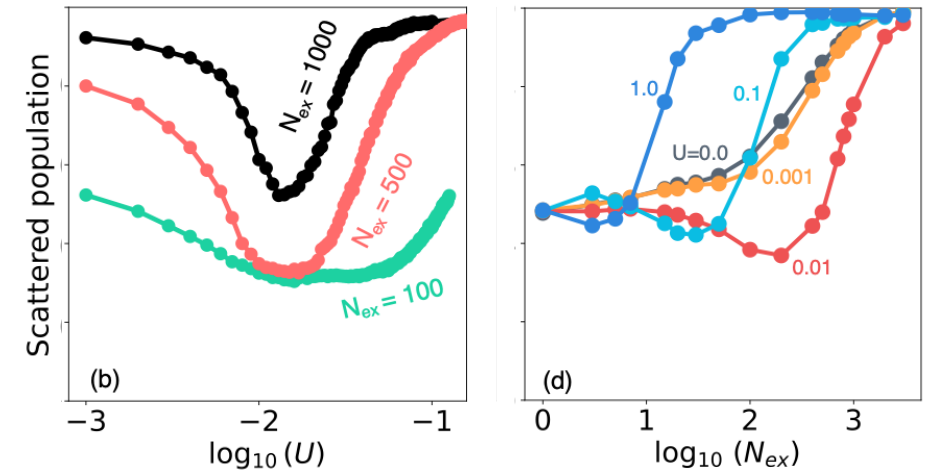


Results

Quantum Dynamics of Interacting Exciton-Polaritons



Interaction Strength Dependent Non-linearity



Many-body interactions can **combat** phonon induced decoherence.



Light & Matter Research Group
at Texas A&M

Thank you



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