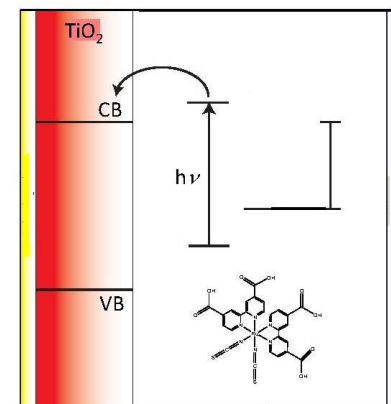
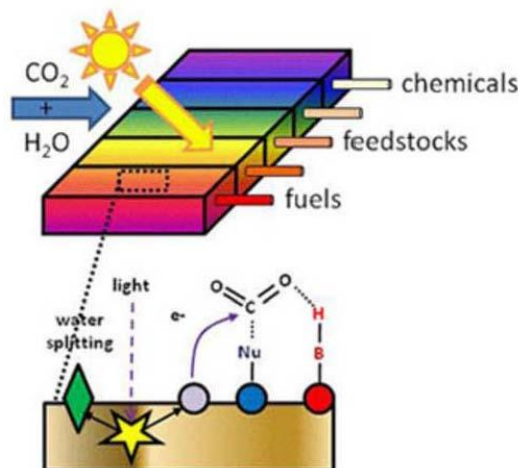
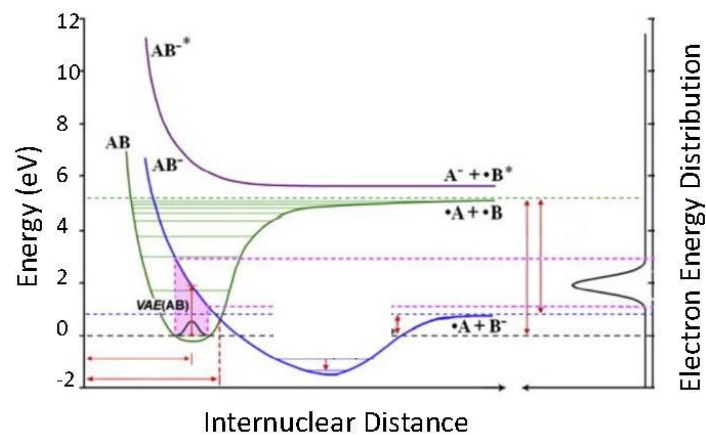
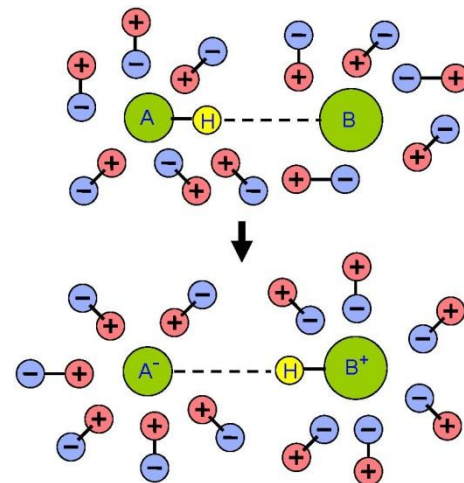
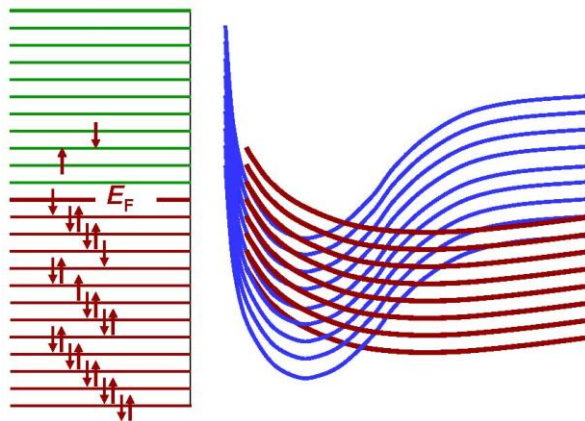
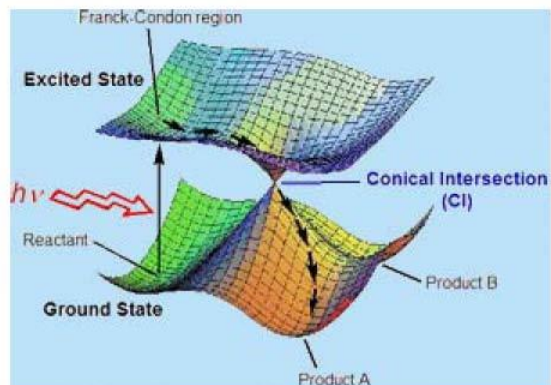


Adventures with Surface Hopping



VISTA: Virtual International Seminar on Theoretical Advancements

Adventures with Surface Hopping

A historical view of my role in some of the
origins and developments of surface hopping

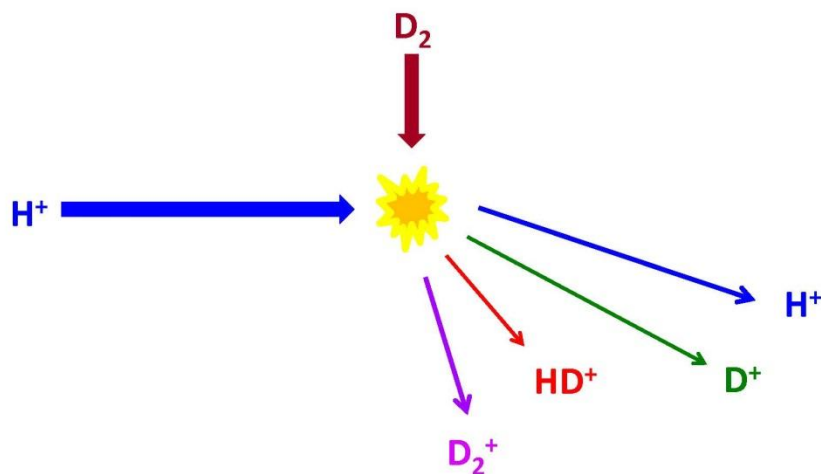
This will not be a scholarly review:

I will omit many important contributions by others

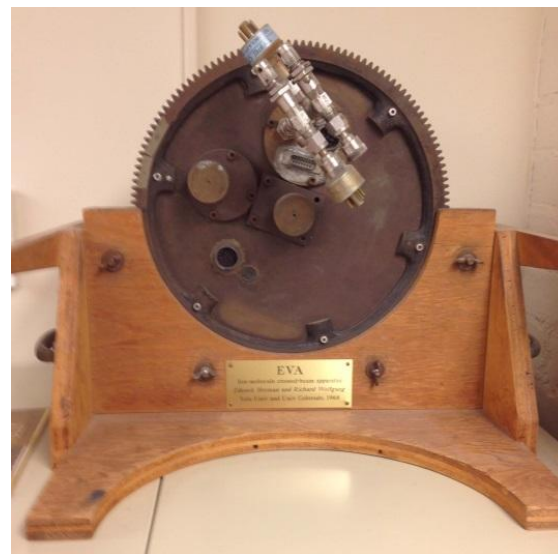
Molecular Beam Experiments of the Reaction of H^+ with D_2

Collaboration with Zdenek Herman, Richard Wolfgang, ...

1969 Univ Colorado Boulder



*Detect the angular and
velocity distributions of ions*



Theory: How hard can it be?

H_3^+ : 3 nuclei, 2 electrons

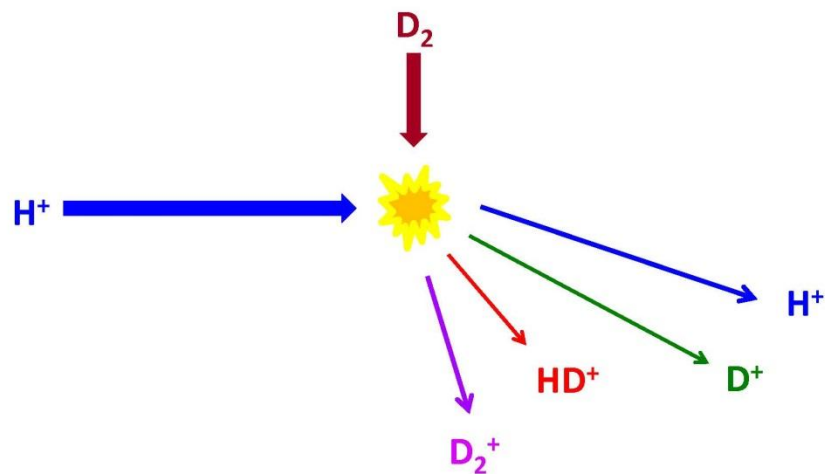
simplest chemical reaction ?

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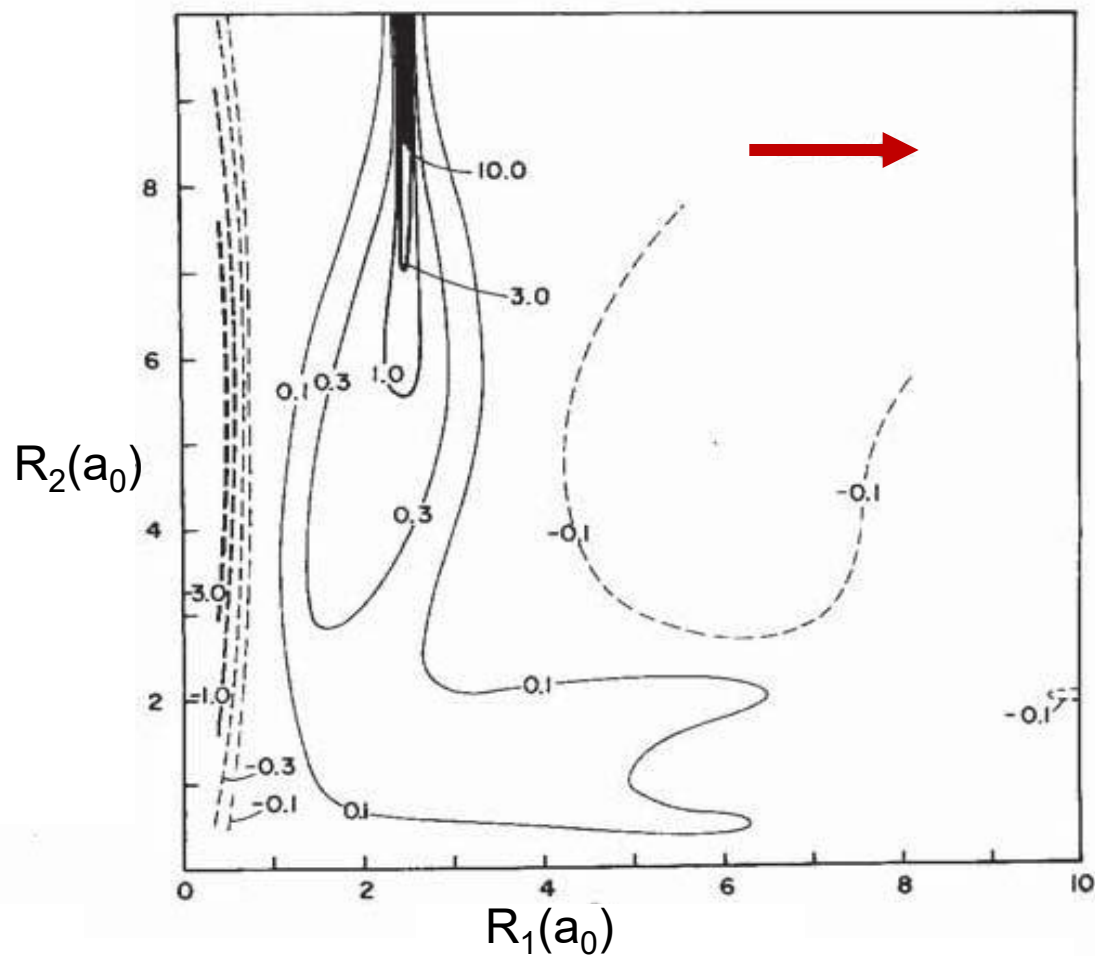
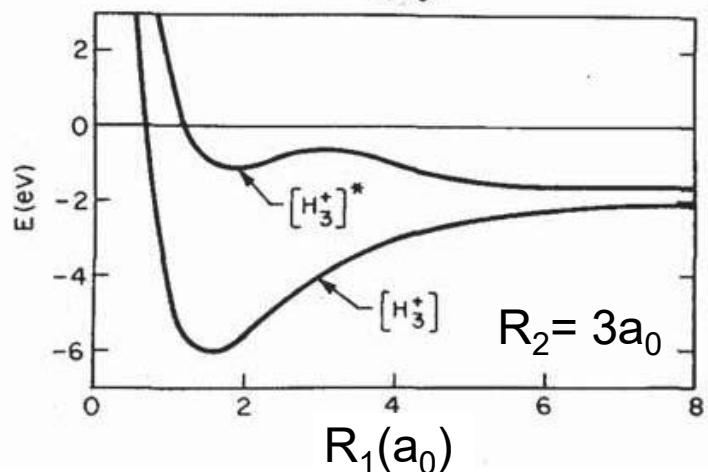
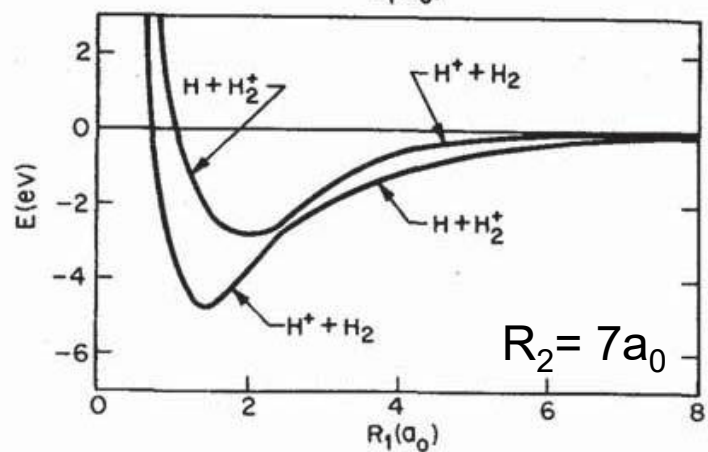
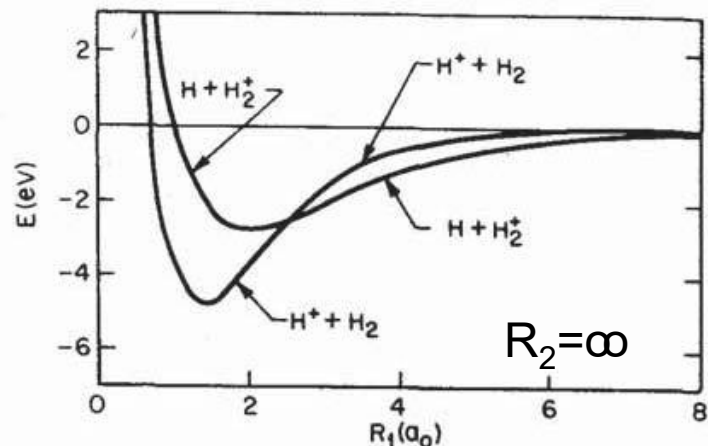
simplest chemical reaction ?

However: 2 potential energy surfaces !



Compute Potential Energy Surfaces and Nonadiabatic Coupling Vectors

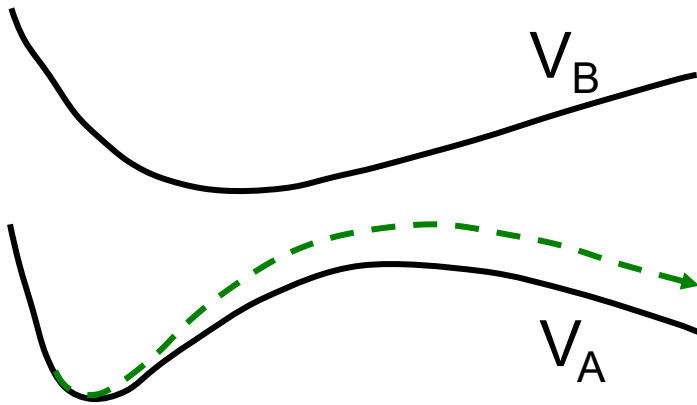
Semiempirical *Diatomics in Molecules*



Nonadiabatic coupling vector

January 14, 2026

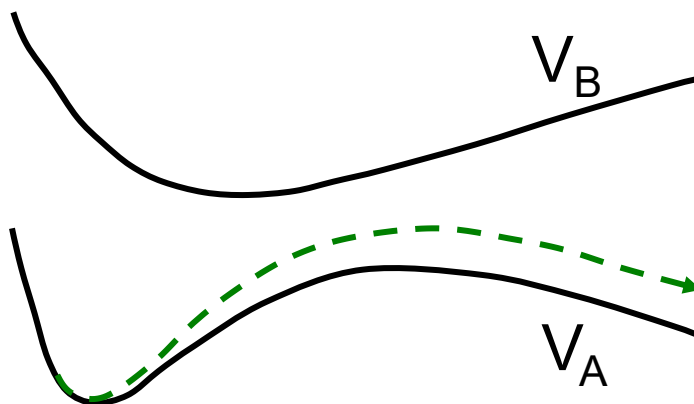
MIXED QUANTUM-CLASSICAL STRATEGIES FOR NONADIABATIC DYNAMICS



Ehrenfest
(self-consistent field)

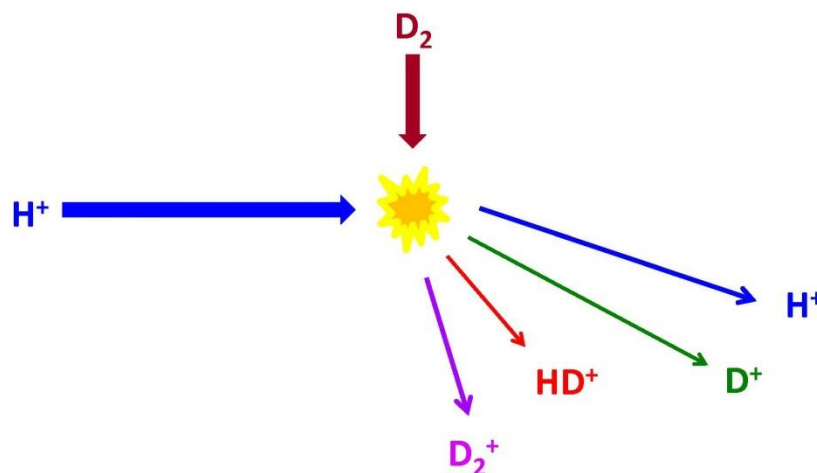
e.g., A. D. McLachlan, *Mol. Phys.* 1964

MIXED QUANTUM-CLASSICAL STRATEGIES FOR NONADIABATIC DYNAMICS

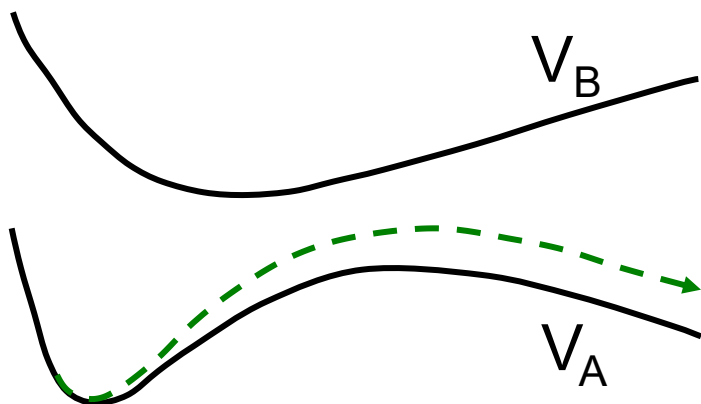


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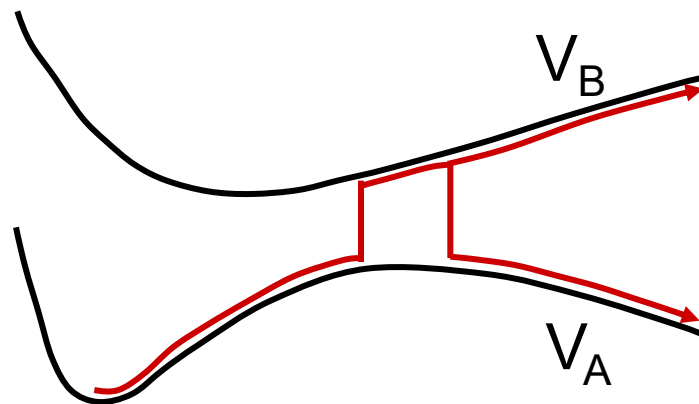
e.g., A. D. McLachlan, *Mol. Phys.* 1964



MIXED QUANTUM-CLASSICAL STRATEGIES FOR NONADIABATIC DYNAMICS



Ehrenfest
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Trajectory Surface-Hopping



motion on individual
potential energy surfaces

Surface Hopping (1971 Version)

1] $M \ddot{R}(t) = -\nabla_R \mathcal{E}_k$, i.e., classical motion on a single adiabatic p.e.s.

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2] Electronic amplitudes along path: $dc_j/dt = -\frac{i}{\hbar} V_{jj} c_j - \dot{R} \cdot \sum_i \langle \varphi_j | \nabla_R \varphi_i \rangle c_i$

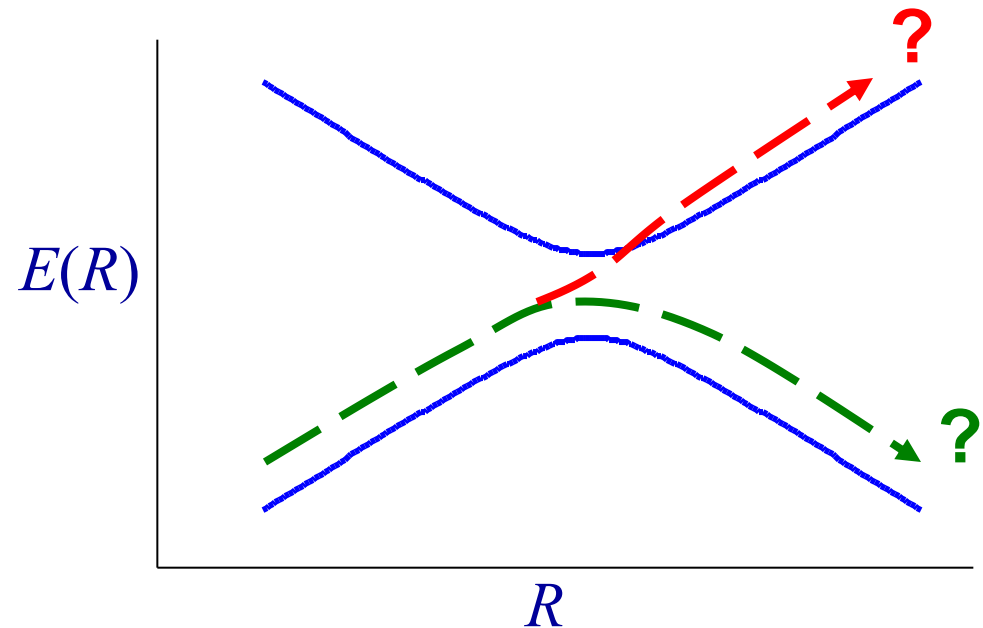
Approximate $|c_k|^2$ by Landau-Zener formula

Landau-Zener Approximation: 1932

2 x 2 Hamiltonian matrix

Assumptions:

1. H_{11} and H_{22} linear
2. H_{12} constant
3. Velocity constant



$$P_{\text{nonad}} \approx \exp \left[\frac{-2\pi H_{12}^2}{\hbar \dot{R} |\partial(H_{11} - H_{22}) / \partial R|} \right]$$

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3] Select electronic state by random number according to $|c_k|^2$

Hops occur only at crossing points

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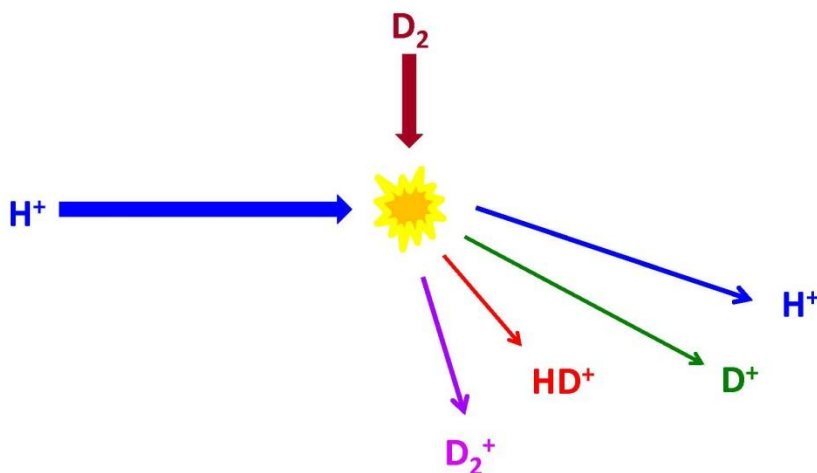
4] Scale momentum in *direction of coupling vector* to conserve energy

P. Pechukas, Phys. Rev., 1969

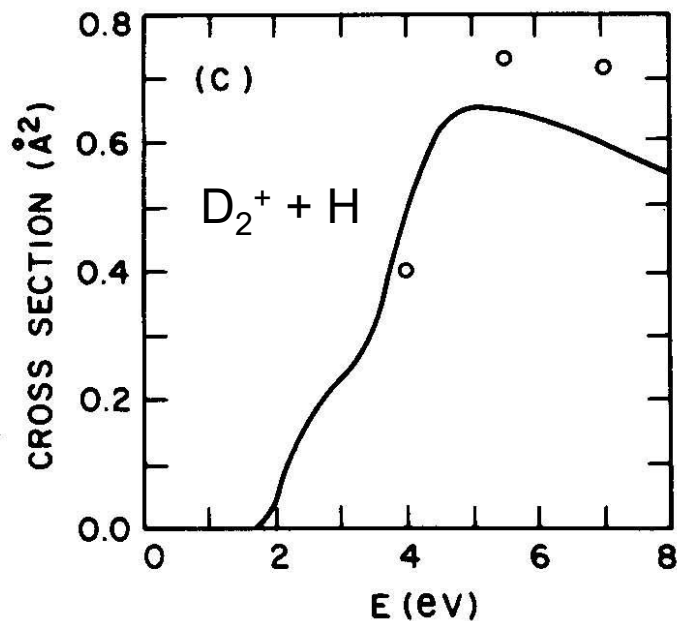
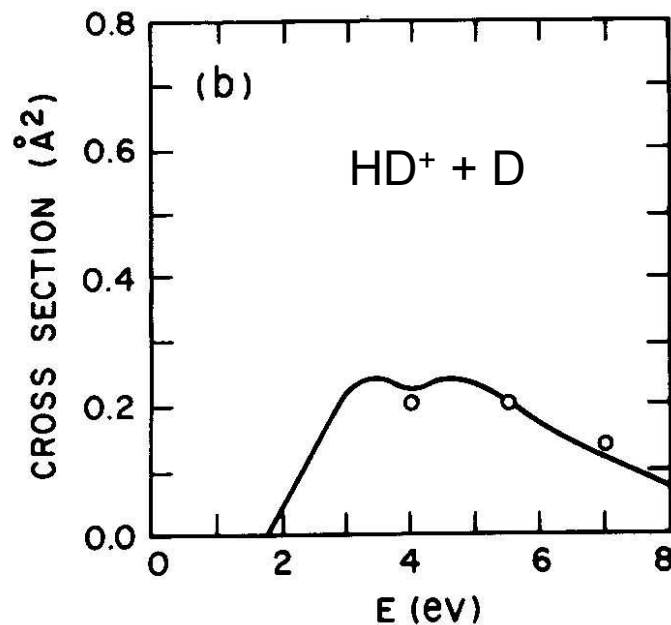
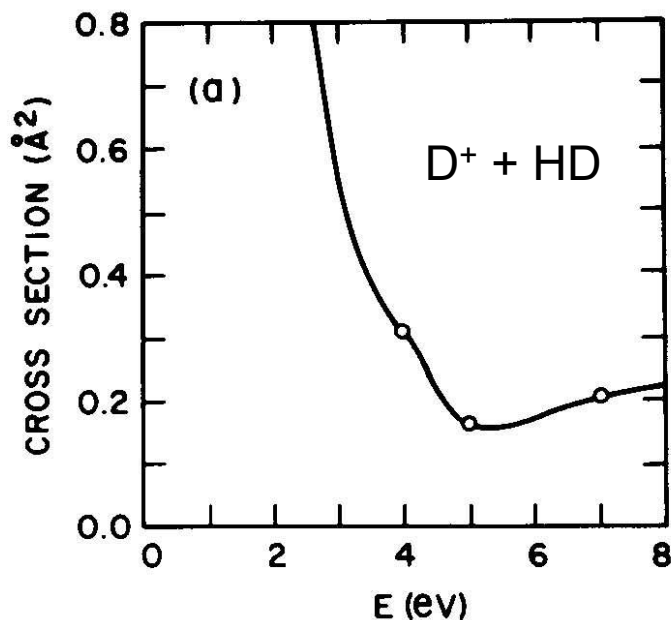
Trajectory Surface Hopping Approach to Nonadiabatic Molecular Collisions: The Reaction of H^+ with D_2

1971 with graduate student Richard Preston

Used Brookhaven National Labs CDC 6600 supercomputer



“It takes a laser physicist to be a farmer” Richard Preston, 2010



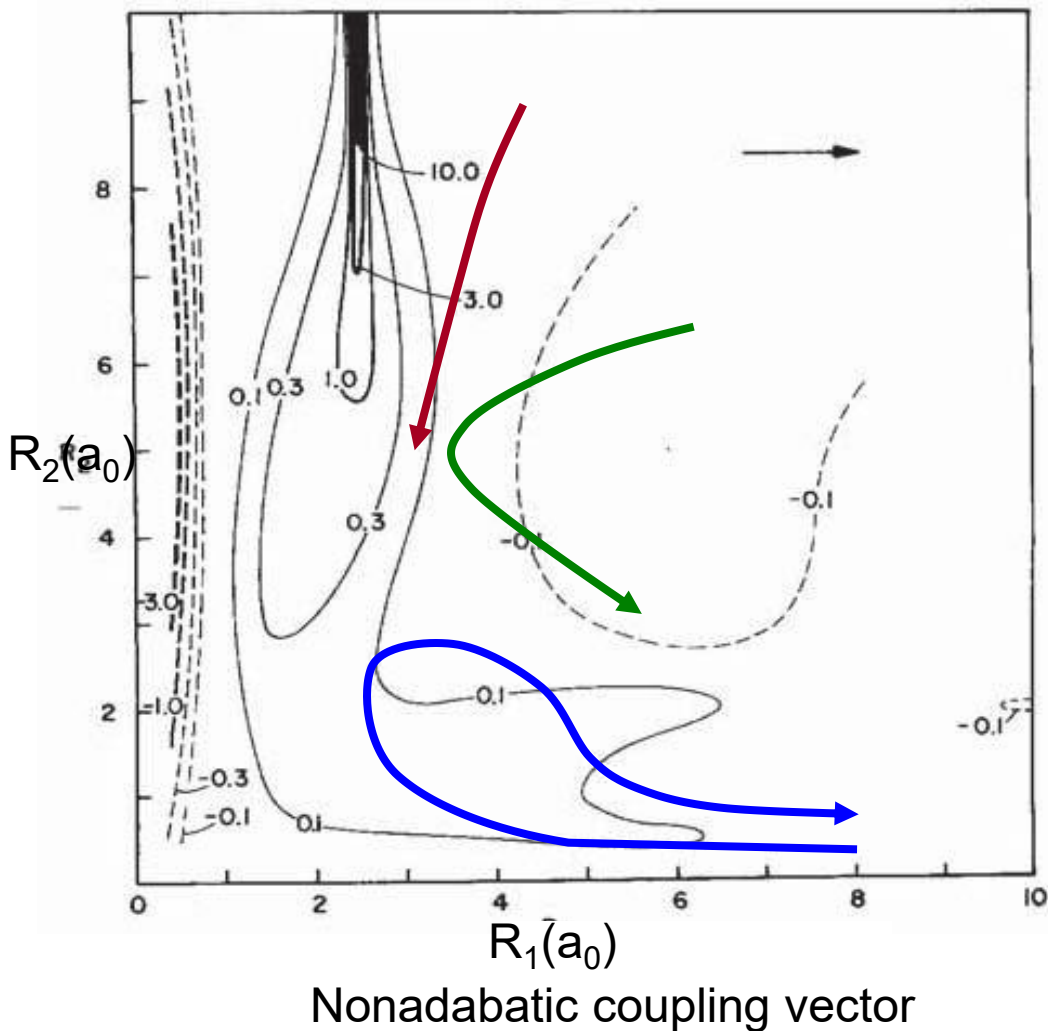
Absolute cross sections for
the reaction of H^+ with D_2
as a function of incident energy

Points: surface hopping,
Tully - Preston, 1971

Solid curves: experiment,
Ochs -Teloy, 1974

Reference 1007

Fast Forward to 1990, Bell Labs:



Turning points

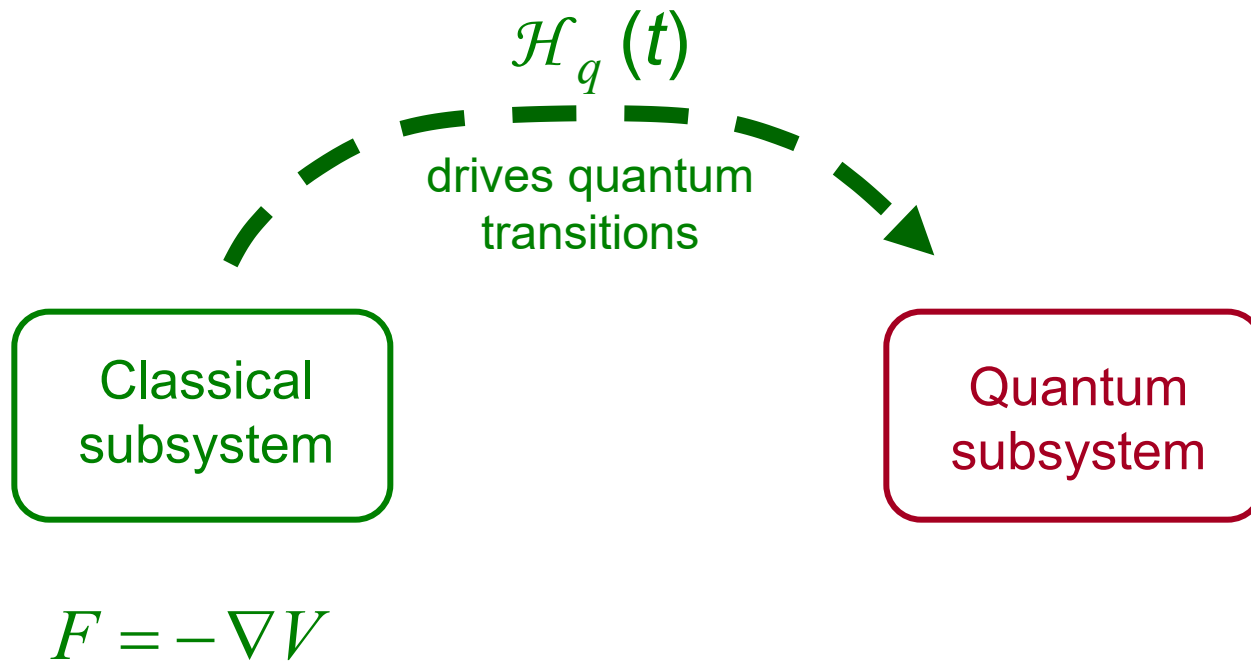
No distinct crossing points

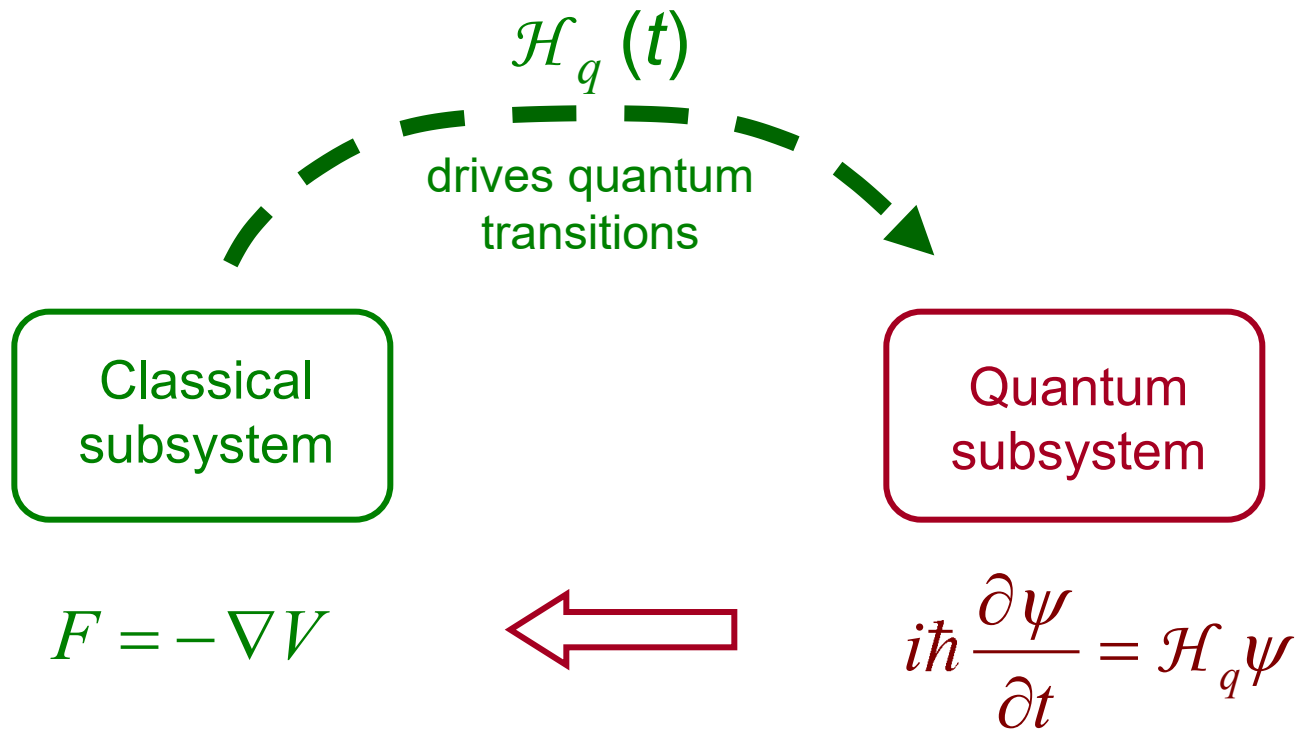
Nearly parallel surfaces

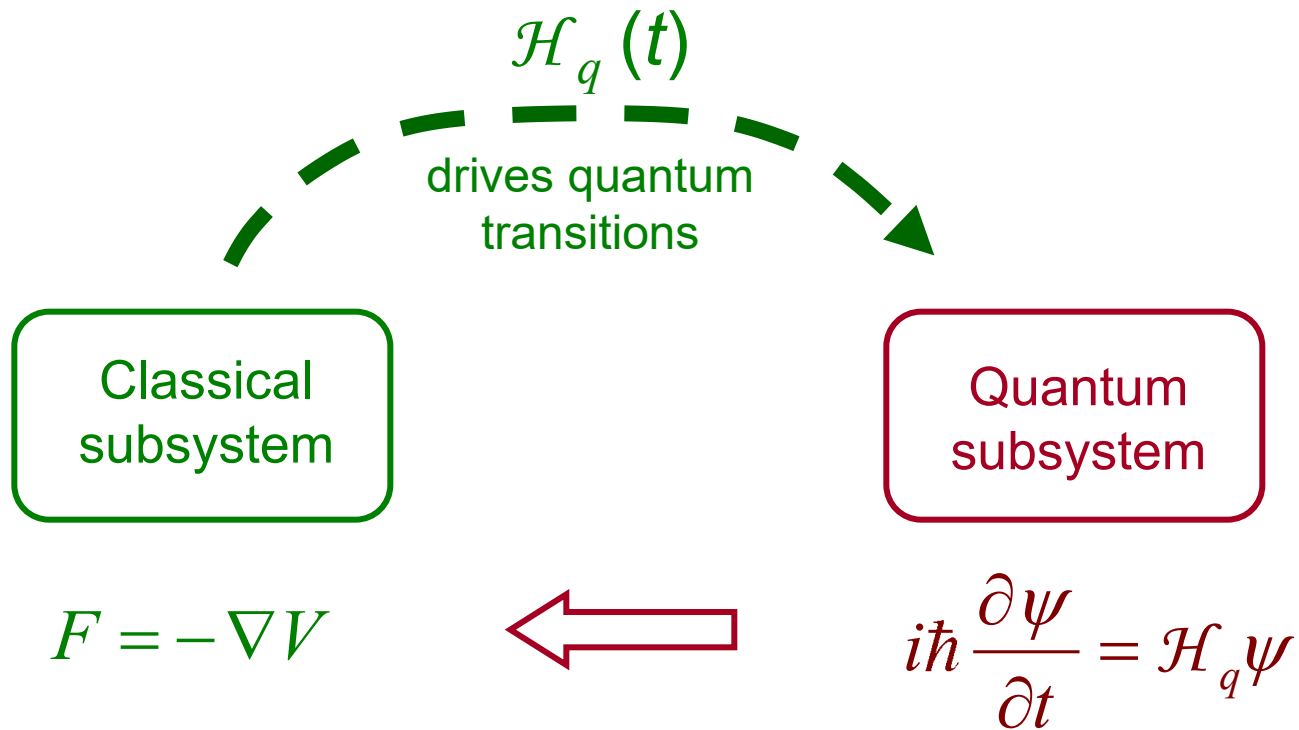
Many electronic states

Broad coupling region

➔ *Need to go beyond
1971 assumption of hopping
only at crossing points*

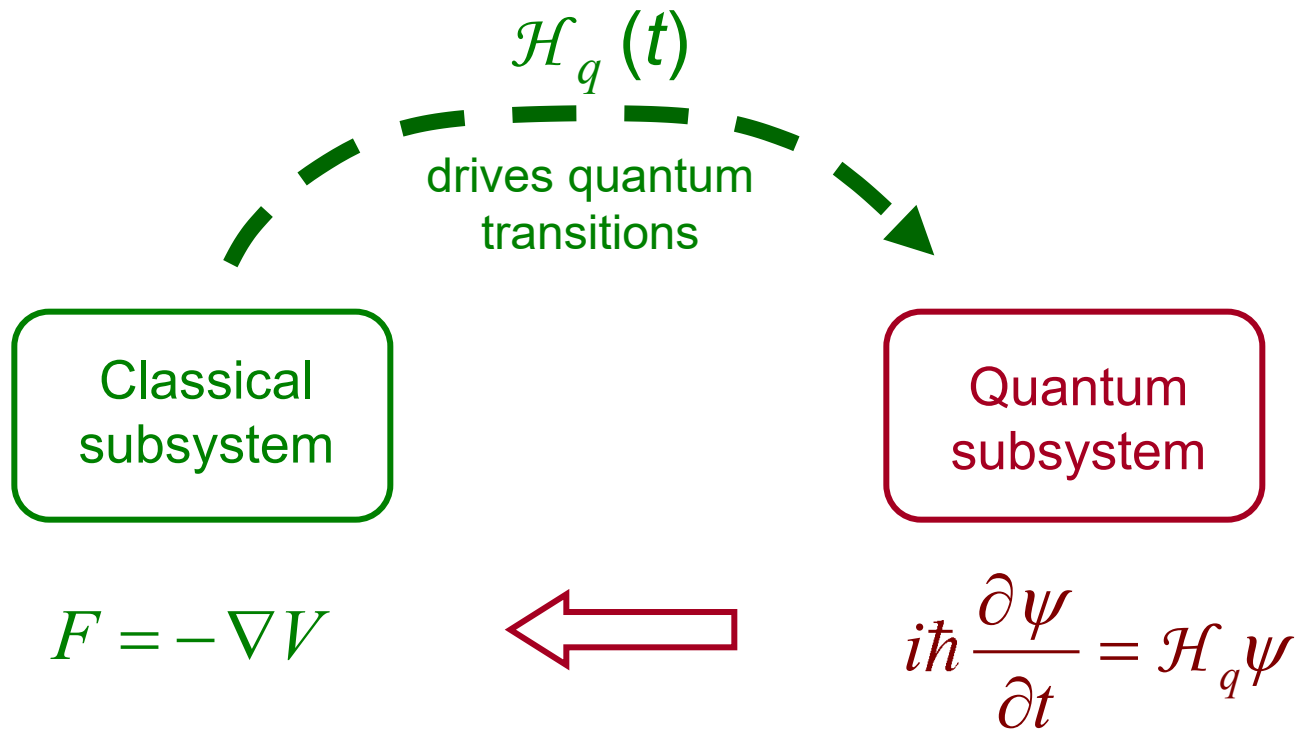






Key: quantum force: “*back-reaction*”

How do quantum transitions effect classical path?



Key: quantum force: “*back-reaction*”

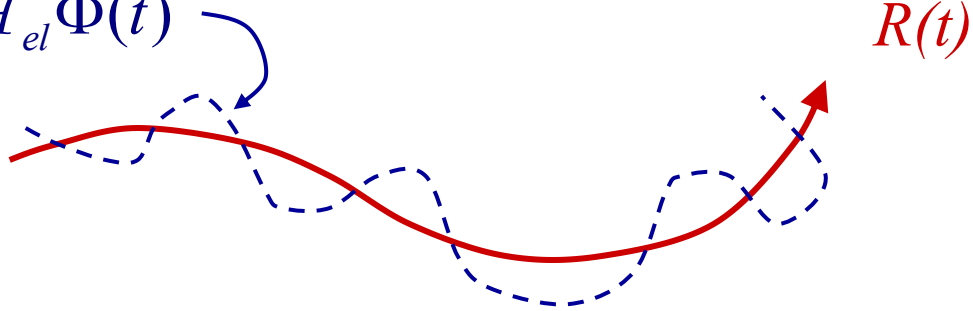
How do quantum transitions effect classical path?

How can we treat back-reaction more accurately, and yet retain the feature that trajectories end up in distinct states rather than mixtures?

Surface Hopping (1990 Version)

1] $M \ddot{R}(t) = -\nabla_R \mathcal{E}_k$, i.e., motion on a single adiabatic p.e.s.

2] Quantum amplitudes along path: $dc_j/dt = -\frac{i}{\hbar} V_{jj} c_j - \dot{R} \cdot \sum_i \langle \varphi_j | \nabla_R \varphi_i \rangle c_i$

$$i\hbar \frac{\partial \Phi(t)}{\partial t} = \mathcal{H}_{el} \Phi(t)$$


Surface Hopping (1990 Version)

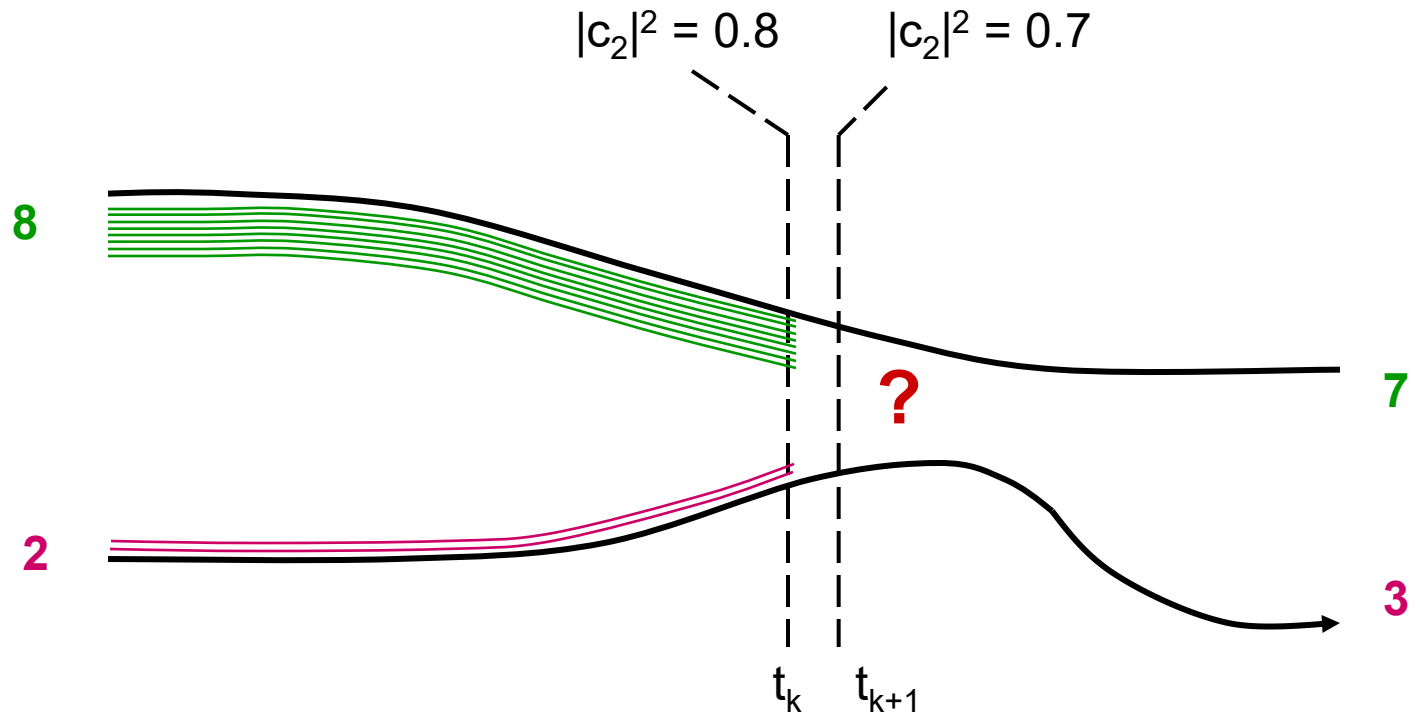
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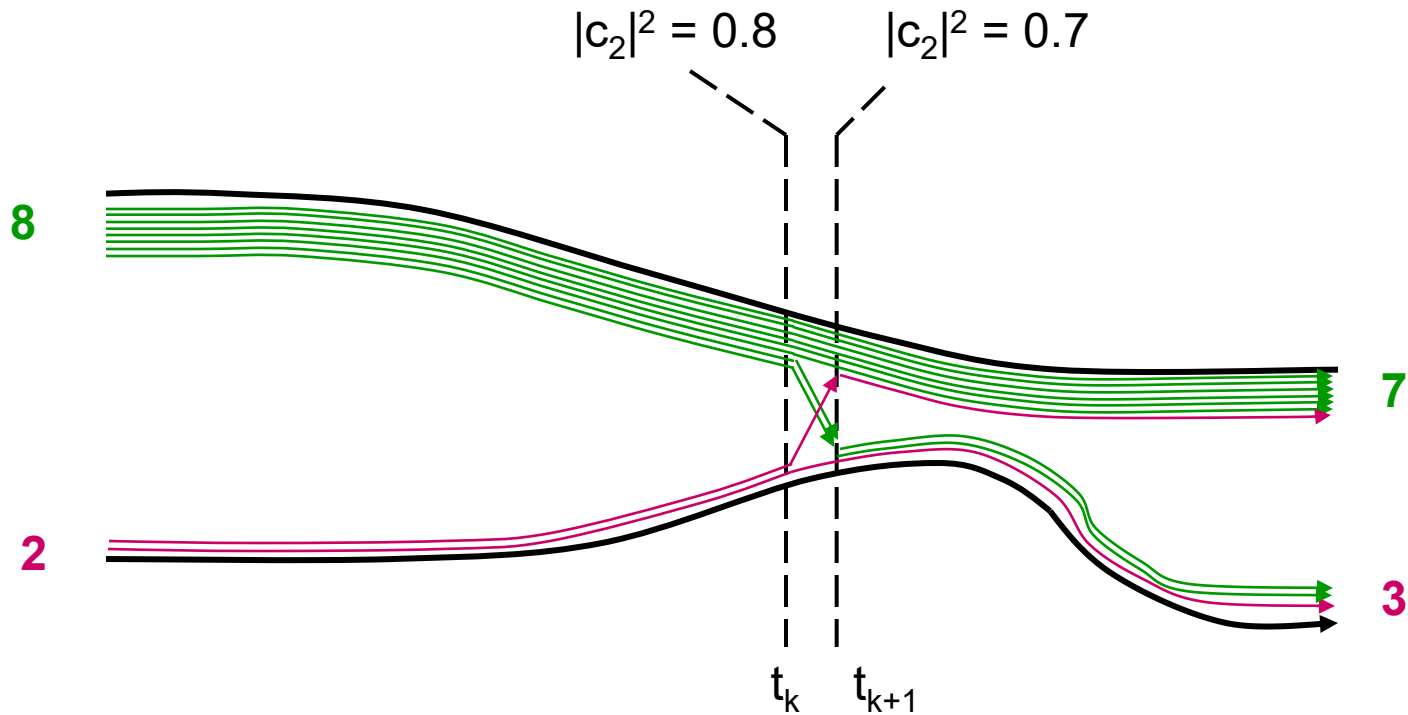
3] Stochastic “hops” between states so that avg probability = $|c_k(t)|^2$

“fewest switches”

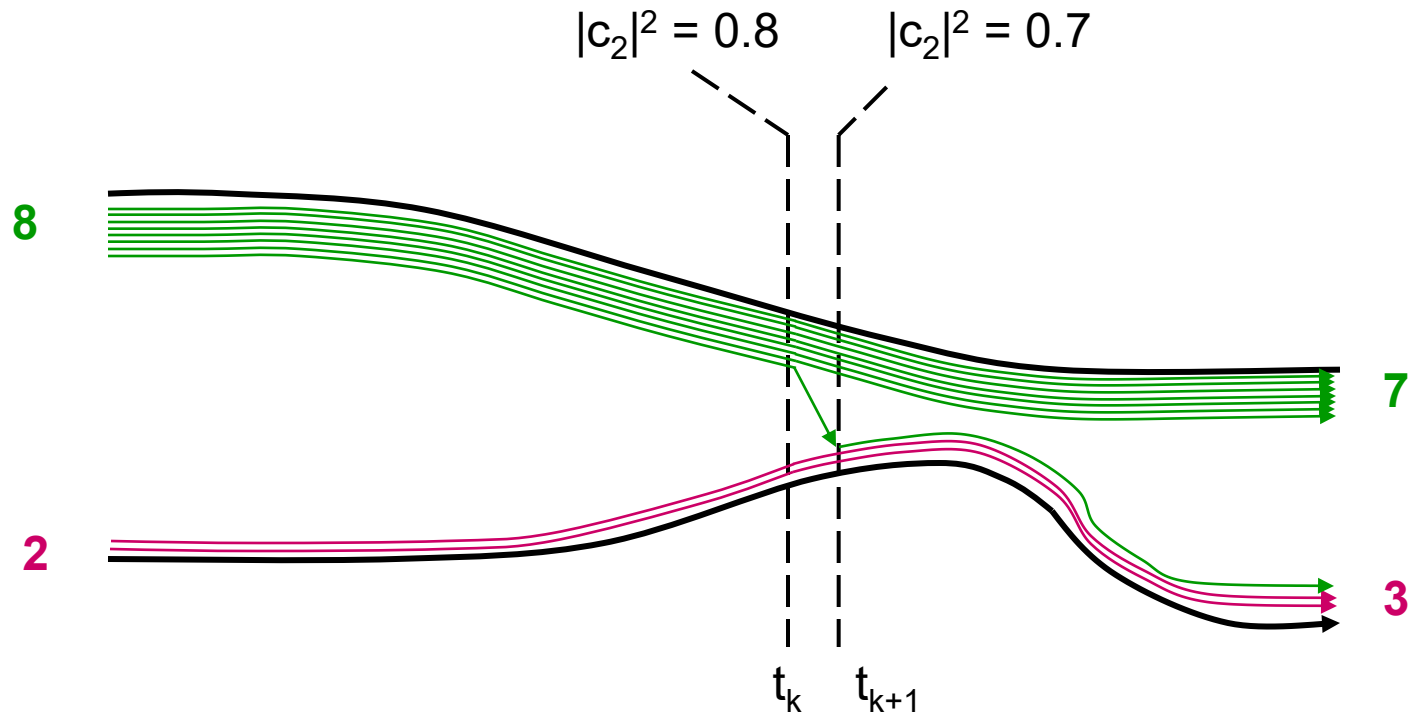
Stochastic *Fewest Switches* algorithm (2-state):



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$$P_{2 \rightarrow 1} = \begin{cases} \frac{|c_2(k)|^2 - |c_2(k+1)|^2}{|c_2(k)|^2}, & |c_2(k)|^2 > |c_2(k+1)|^2 \\ 0, & |c_2(k)|^2 \leq |c_2(k+1)|^2 \end{cases}$$

Surface Hopping (1990 Version)

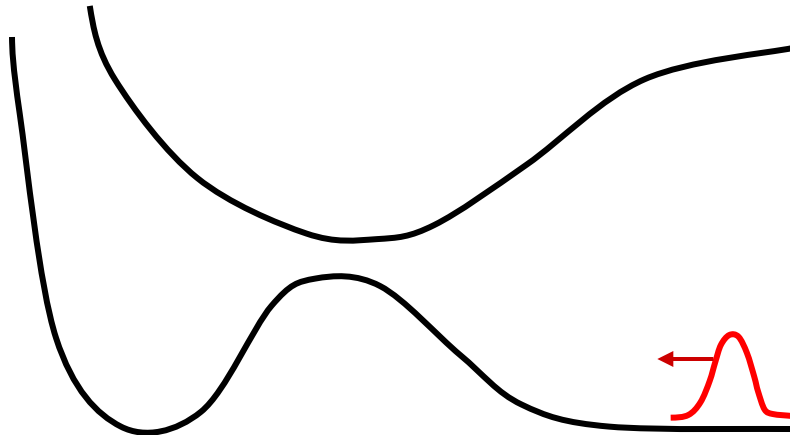
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“Pechukas Force”

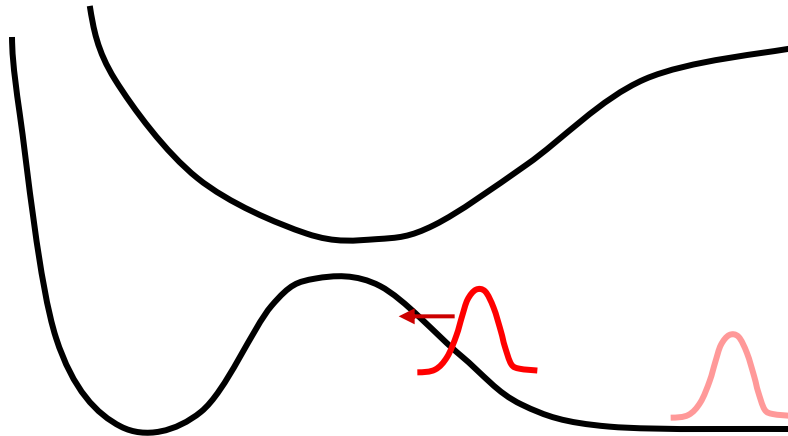
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- 5] Include decoherence

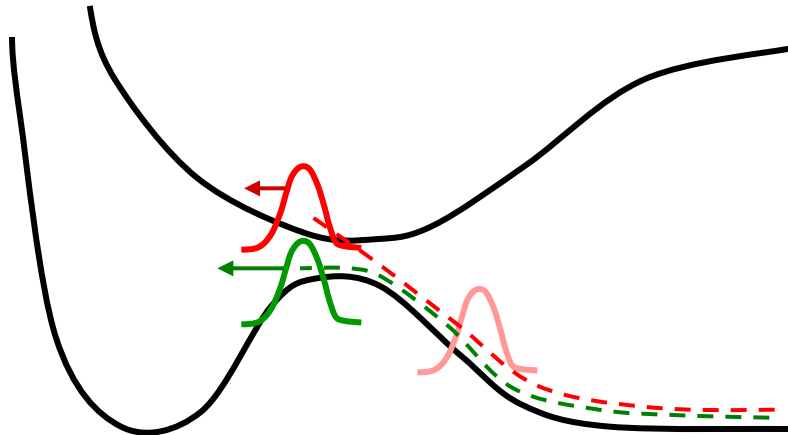
Decoherence



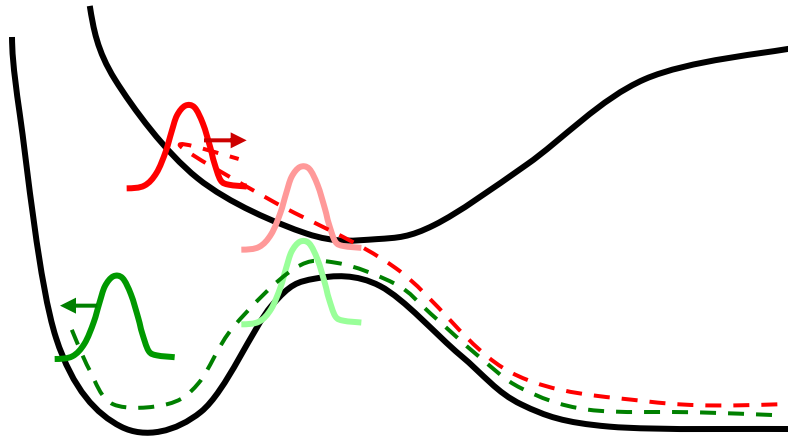
Decoherence



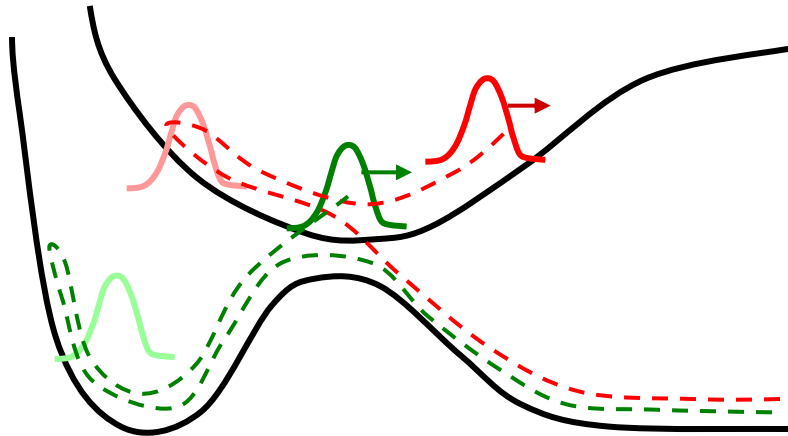
Decoherence



Decoherence



Decoherence



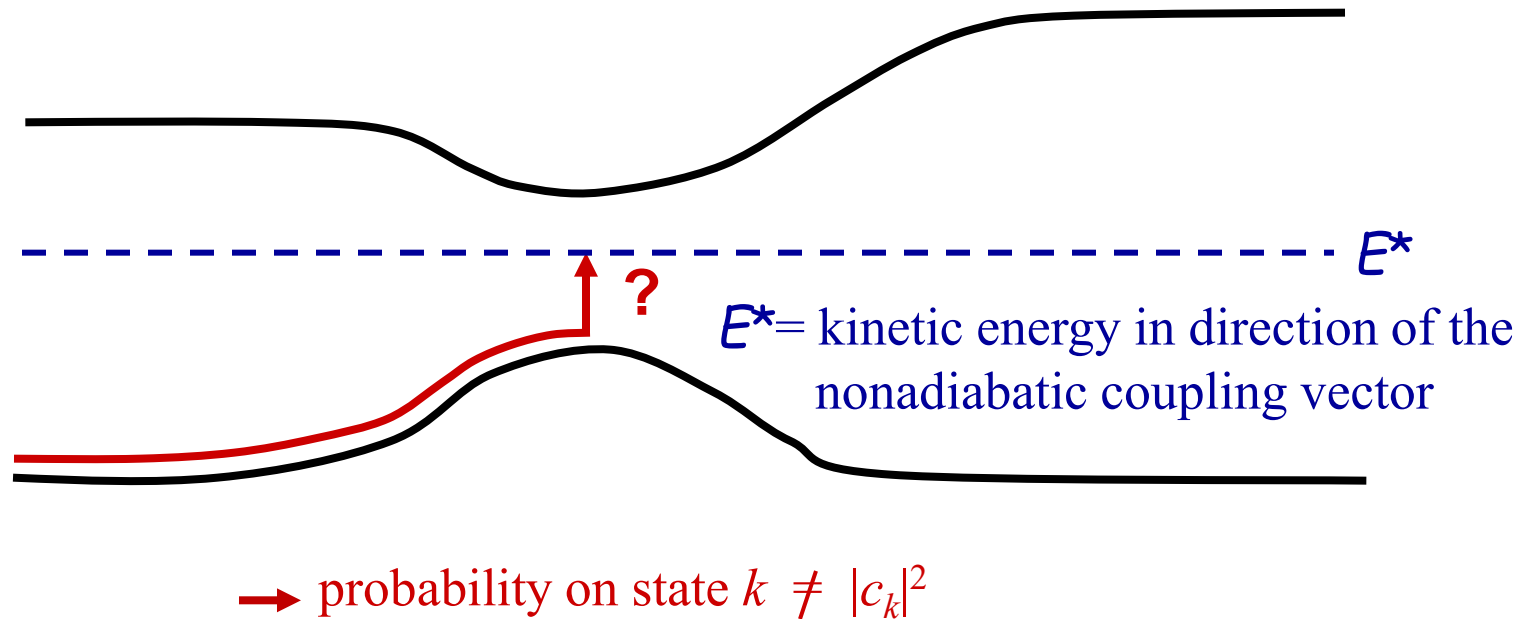
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Algorithm calls for a hop but there is insufficient energy in the degree of freedom in the direction of the nonadiabatic coupling vector

Frustrated Hops

Algorithm calls for a hop but there is insufficient energy in the degree of freedom in the direction of the nonadiabatic coupling vector



Critical element of Surface Hopping: Branching probabilities are determined not by $|c_k|^2$, but by the fraction of trajectories on each state.

Frustrated Hops and Detailed Balance

Detailed Balance:

$$\mathcal{N}_1 P_{1 \rightarrow 2} = \mathcal{N}_2 P_{2 \rightarrow 1} \rightarrow \text{Equilibrium}$$

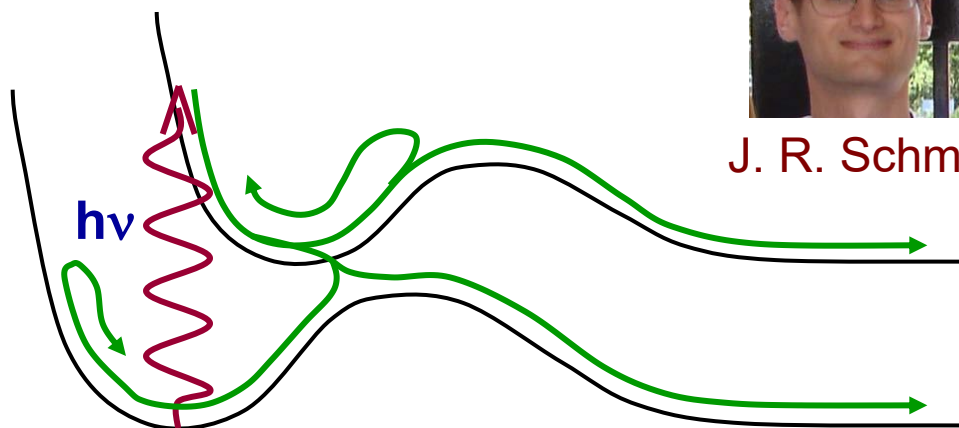


Priya Parandekar



J. R. Schmidt

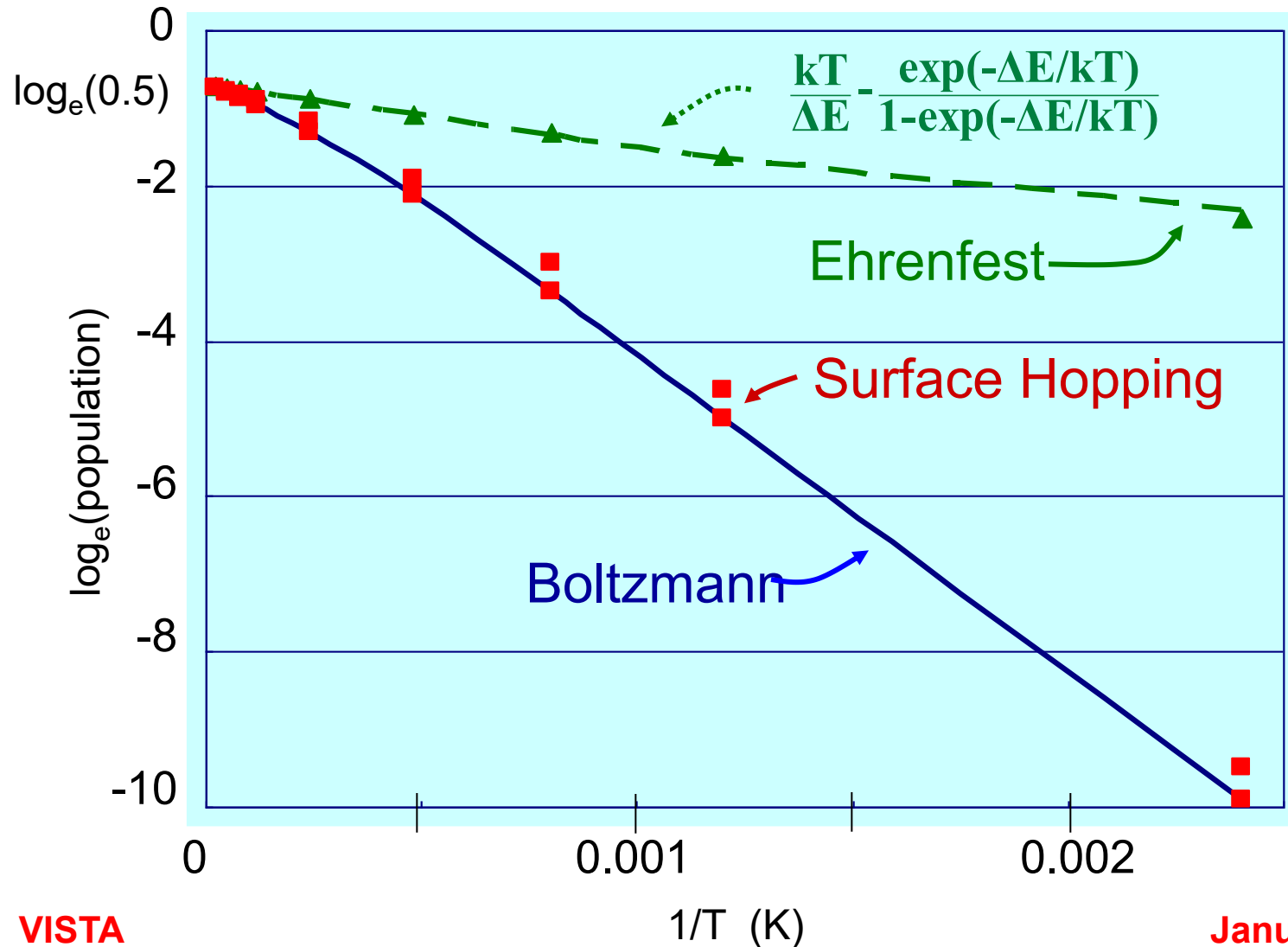
- Multiple Transitions
- Relaxation Processes
- Infrequent events
- Condensed Phases
- *Long Timescales*



e.g., nonradiative transition
vs. reaction on excited state

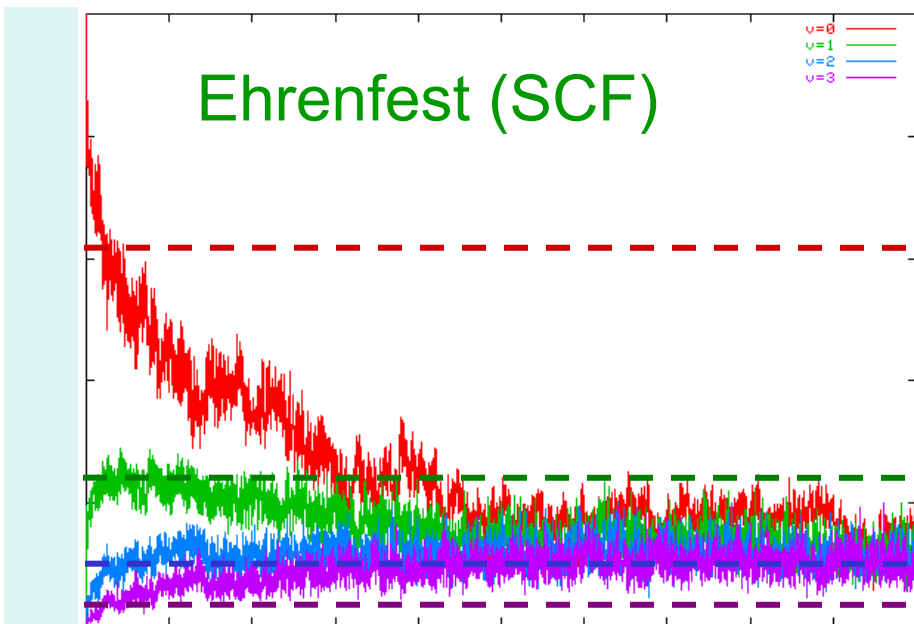
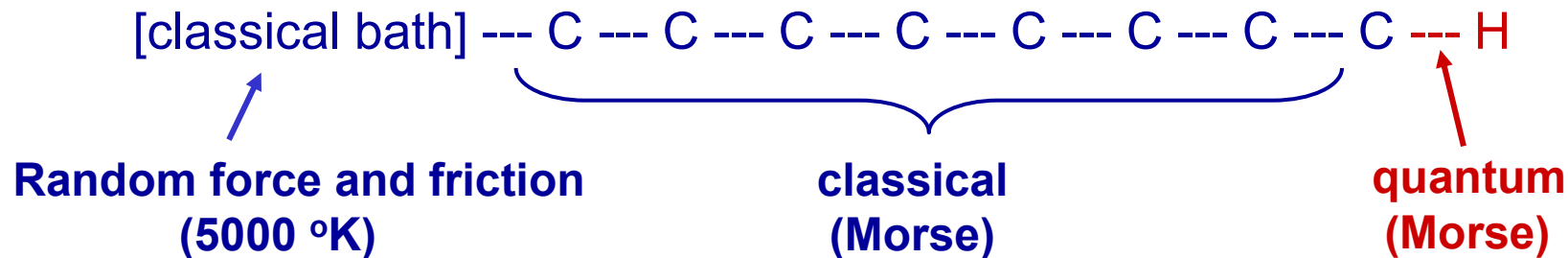
2-Electronic State Model $\Delta E = 34.6$ kJ/mole

Excited State Population vs. Inverse Temperature

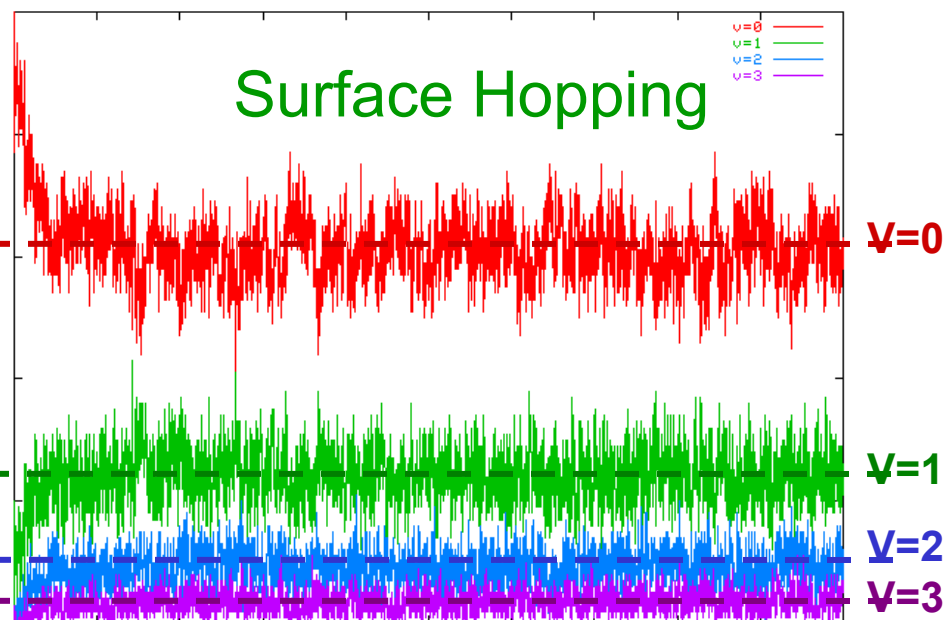


Many Quantum States

2005 Priya Parandekar



time



time

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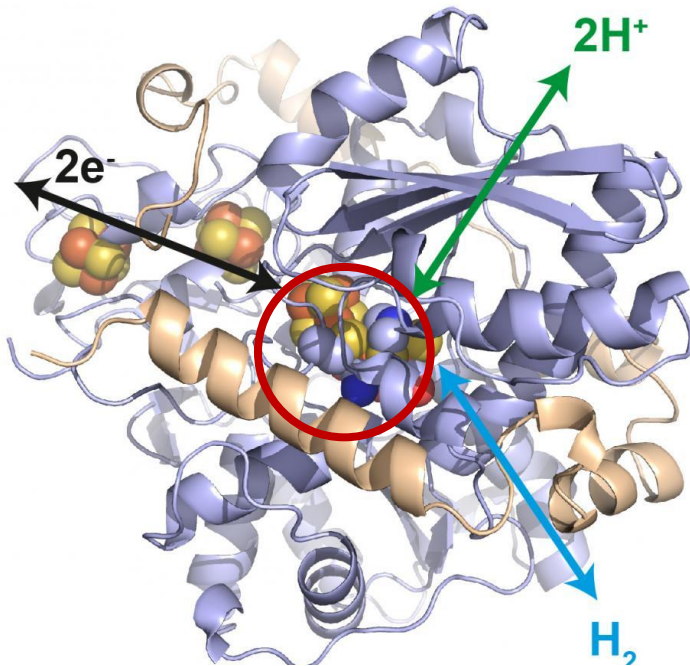
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Treat Nuclear Motion by Quantum Mechanics:

- tunneling
- zero-point motion
- quantized energy levels

Prohibitive for large number of nuclei

Compromise strategy: a few quantum nuclei, the rest classical

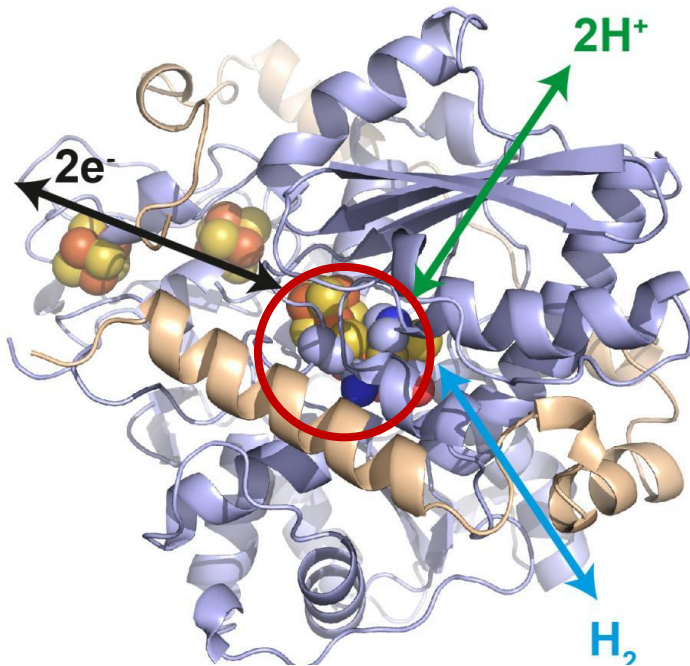


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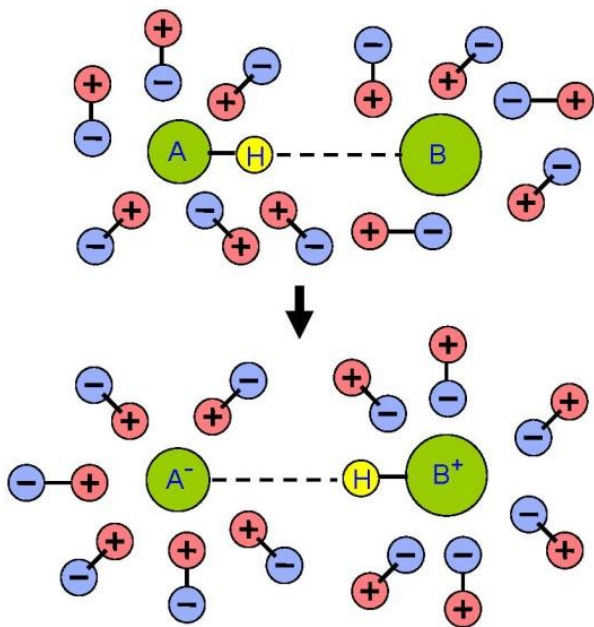
Compromise strategy: a few quantum nuclei, the rest classical



***However:
justification of the
Born-Oppenheimer
approximation ???***

Mixed Quantum-Classical Nuclear Motion

Sharon Hammes-Schiffer and JCT, *J. Chem. Phys.* **101**, 4657 (1994)



Quantum Effects:

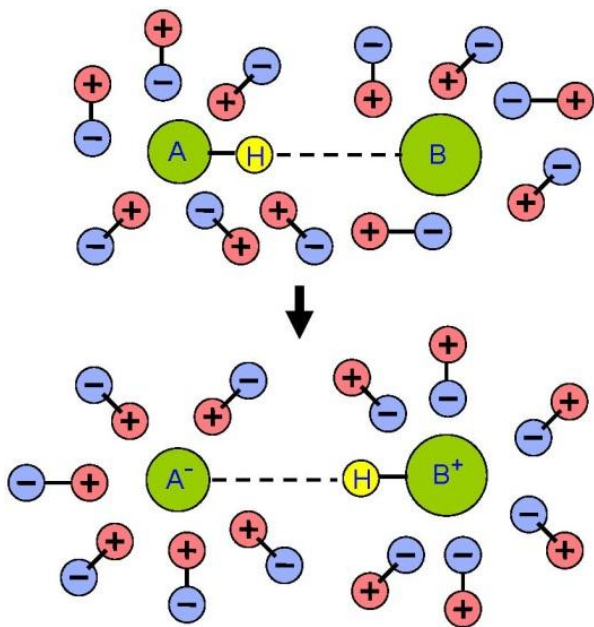
Zero-Point Energy

Quantized Energy Levels

Tunneling

Mixed Quantum-Classical Nuclear Motion

Sharon Hammes-Schiffer and JCT, *J. Chem. Phys.* **101**, 4657 (1994)



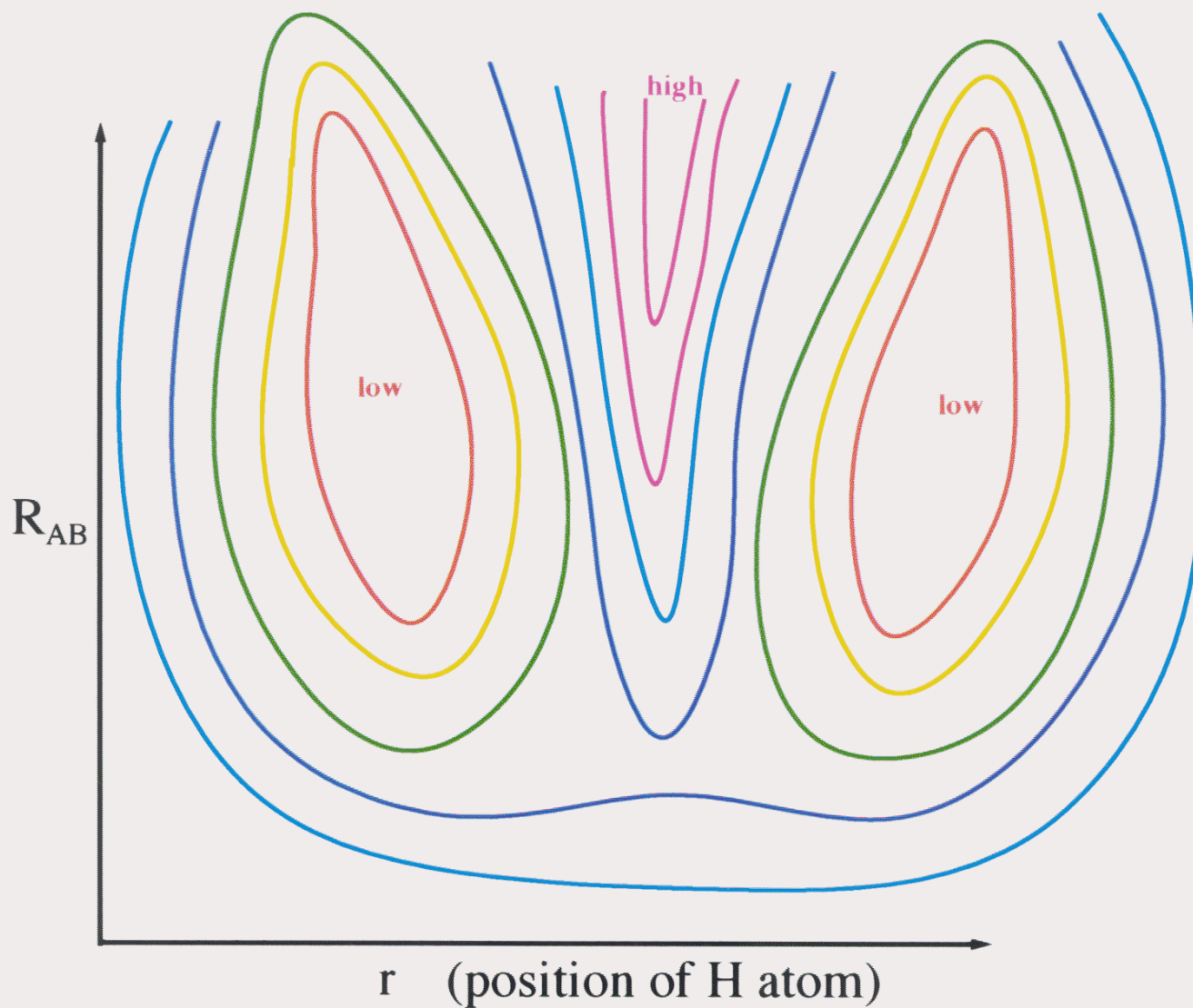
Quantum Effects:

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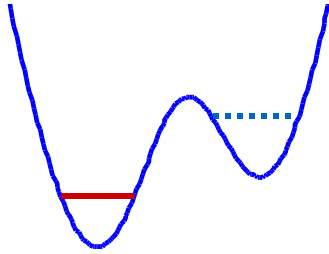
Quantized Energy Levels

Tunneling

CONTOUR PLOT OF FREE ENERGY

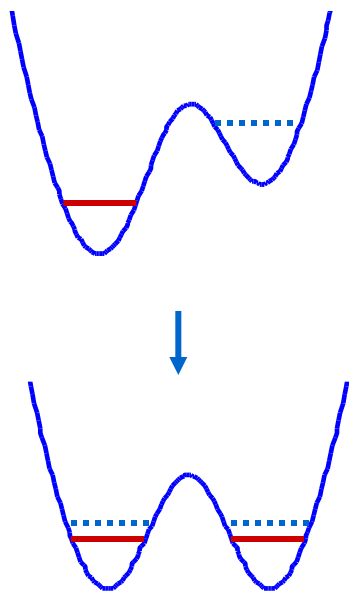


Adiabatic vs. Nonadiabatic (Sharon Hammes-Schiffer, JCT)

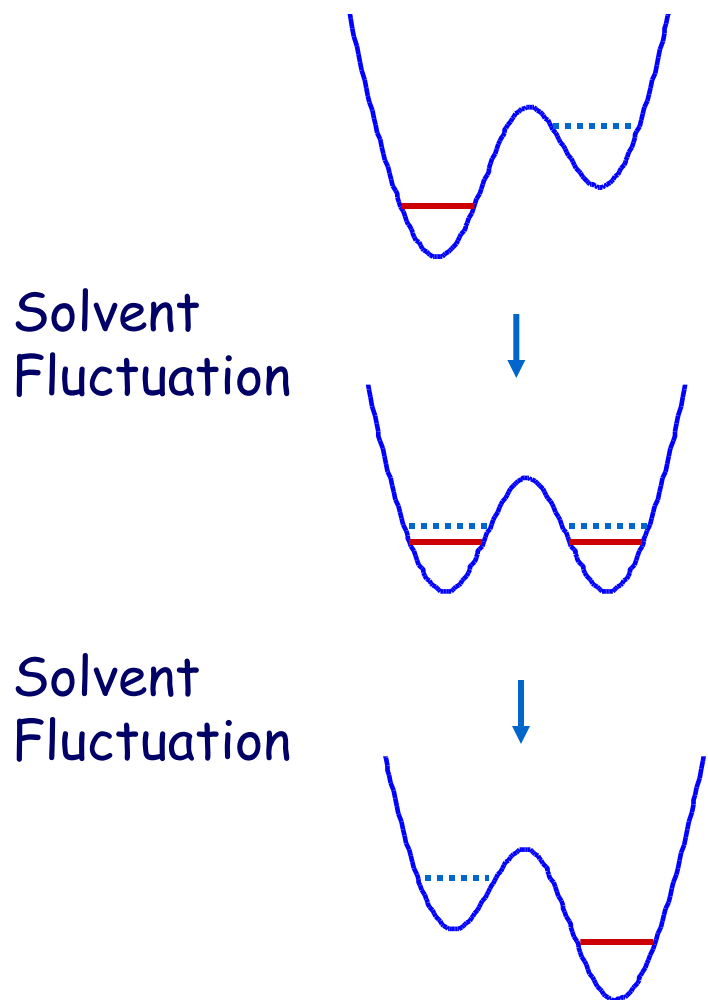


Adiabatic vs. Nonadiabatic (Sharon Hammes-Schiffer, JCT)

Solvent
Fluctuation

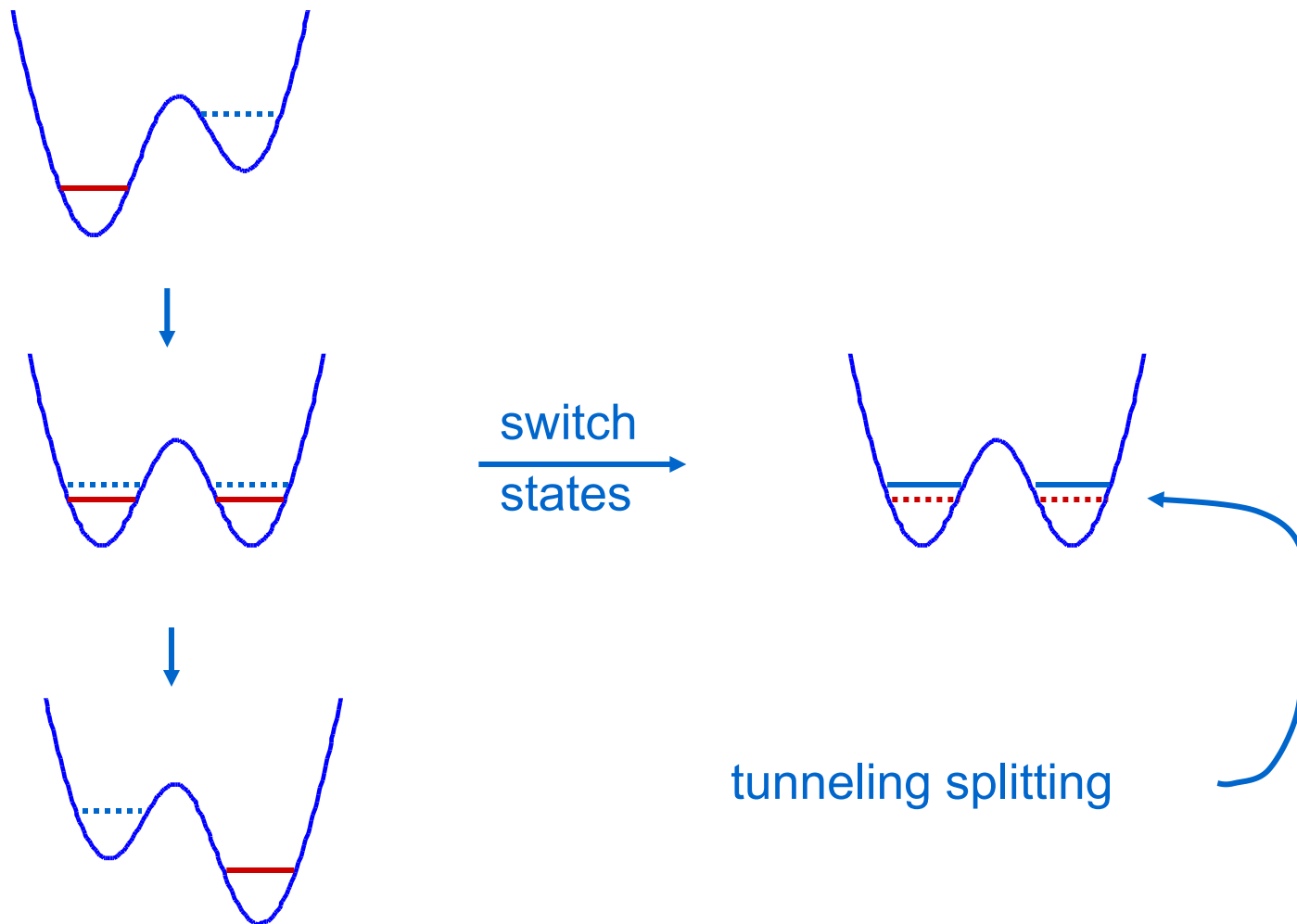


Adiabatic vs. Nonadiabatic (Sharon Hammes-Schiffer, JCT)



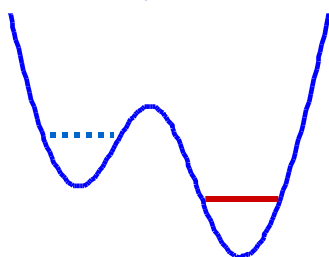
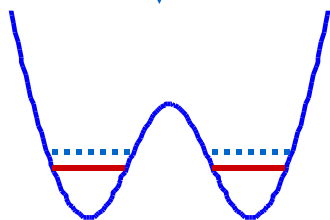
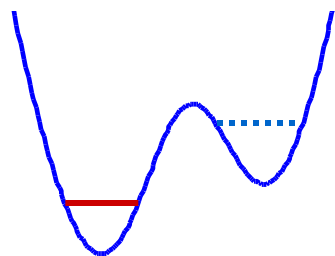
Adiabatic \rightarrow Reaction

Adiabatic vs. Nonadiabatic (Sharon Hammes-Schiffer, JCT)



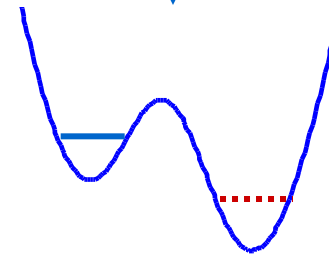
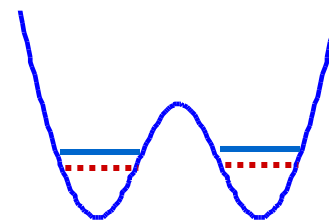
Adiabatic → Reaction

Adiabatic vs. Nonadiabatic (Sharon Hammes-Schiffer, JCT)



Adiabatic → Reaction

switch
states

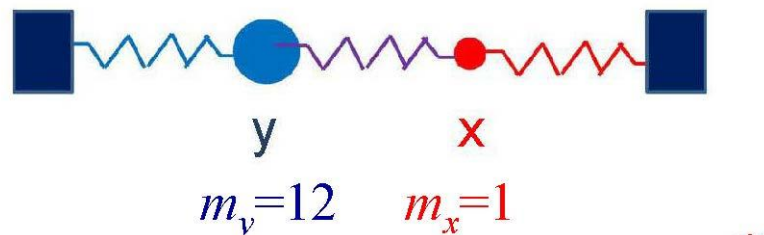


Nonadiabatic → No Reaction

Mixed Quantum-Classical Nuclear Motion

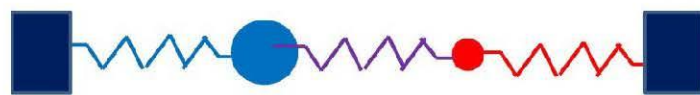
Vibrational Spectra ? JCT, *J. Chem. Phys.* 2023

Simple 2 atom linear model



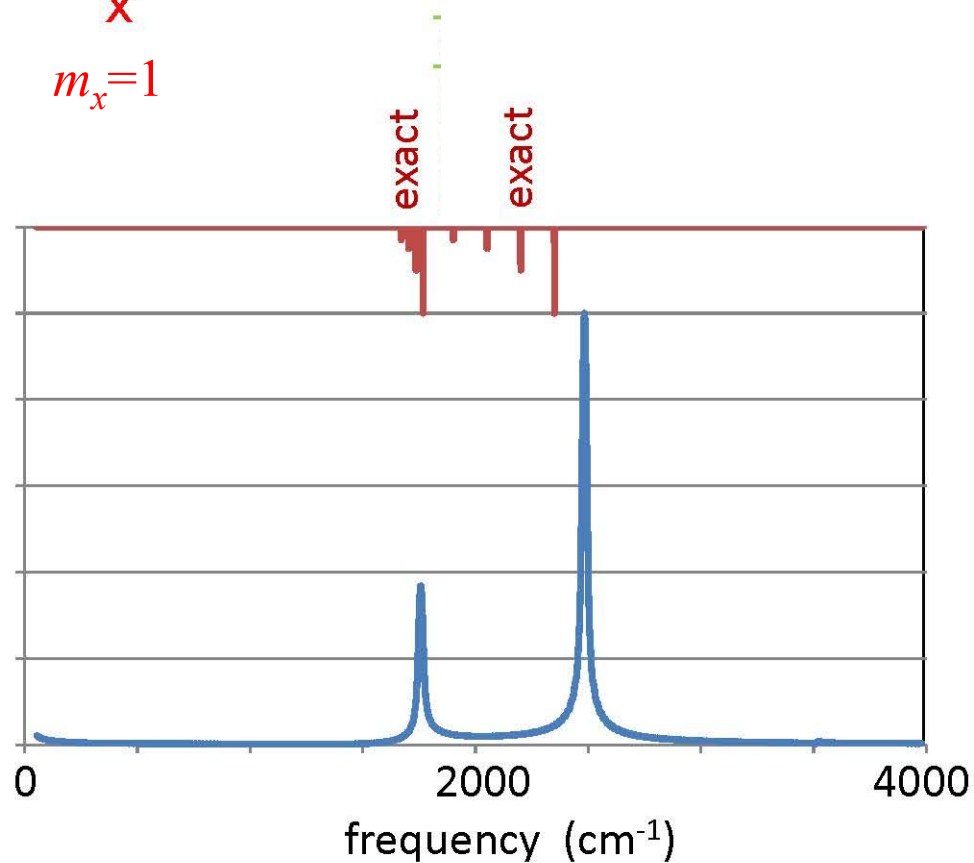
Compute spectra by Fourier Transform
of Dipole Moment Correlation Function

Morse Potentials



$$m_y=12 \quad m_x=1$$

Both atoms classical



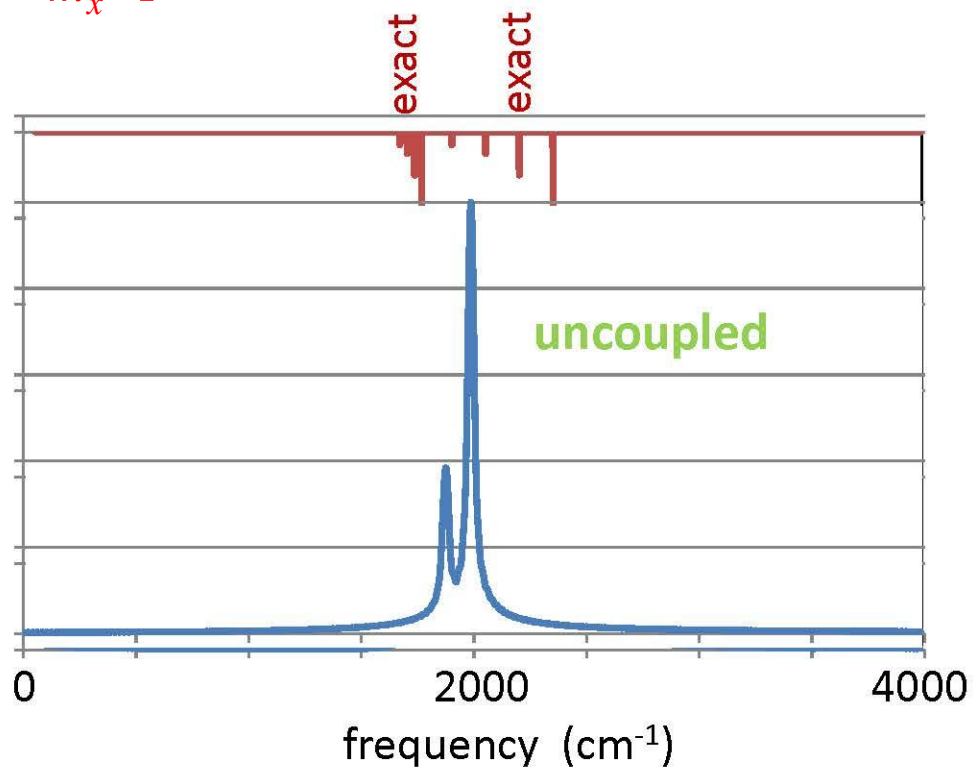
Fourier transform of dipole moment autocorrelation function

Morse Potentials



$m_y=12$ $m_x=1$

Light atom quantum: Adiabatic



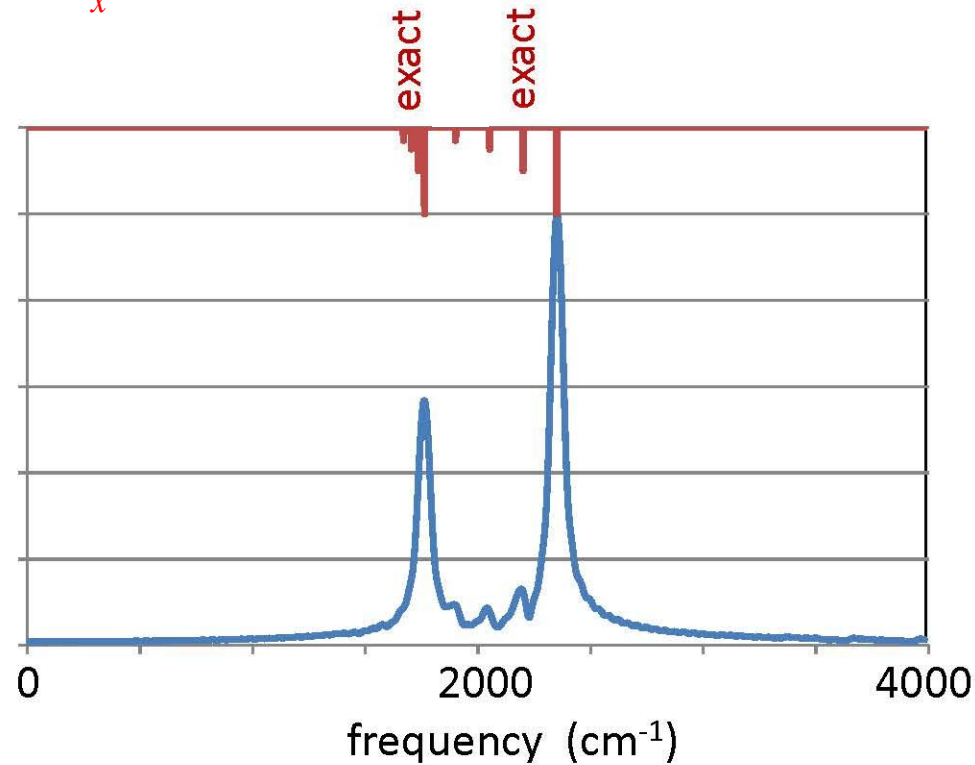
Fourier transform of dipole moment autocorrelation function

Morse Potentials



$m_y=12$ $m_x=1$

Light atom quantum: Ehrenfest

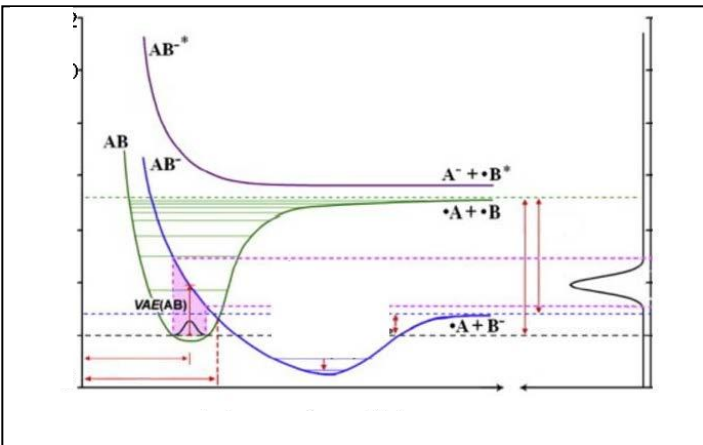


Fourier transform of dipole moment autocorrelation function

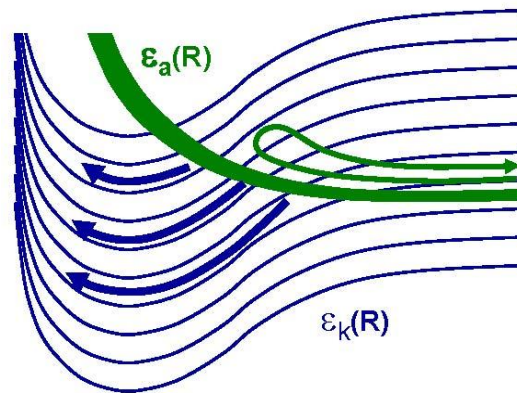
Conclusions about treating some nuclei quantum, others classical

1. Requires nonadiabaticity to achieve coupled motions (e.g., normal modes)
2. True also for electron-nuclear coupled motion, but less dramatic because of mass disparities
3. Ehrenfest incorporates coupled motion
4. *Surface Hopping is problematical*

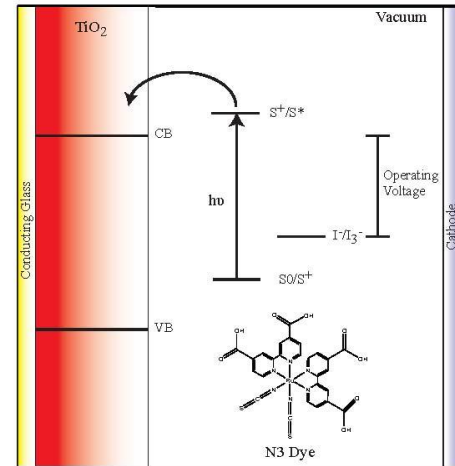
electronic continuum ?



ionization

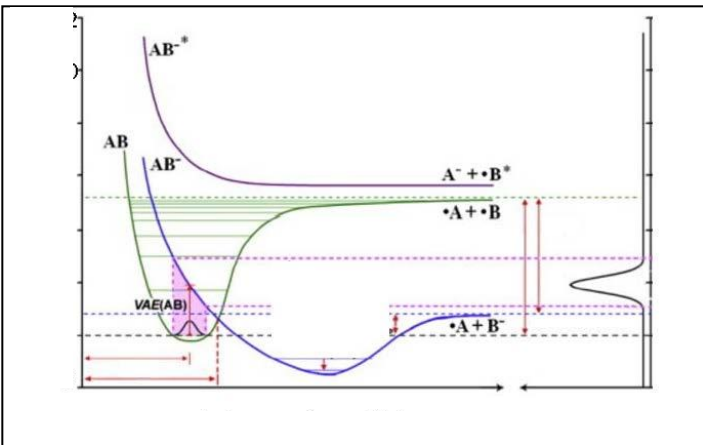


metals

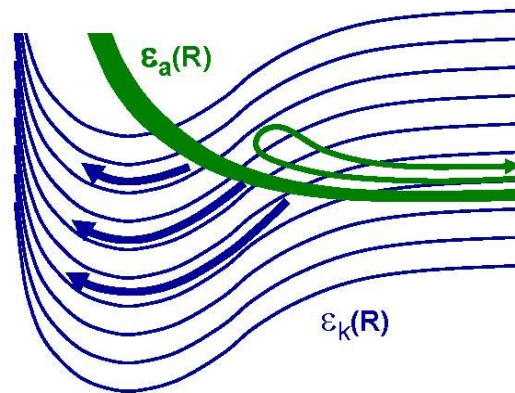


semiconductors

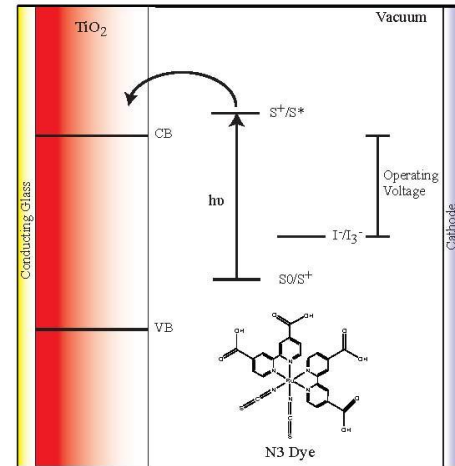
electronic continuum ?



ionization



metals

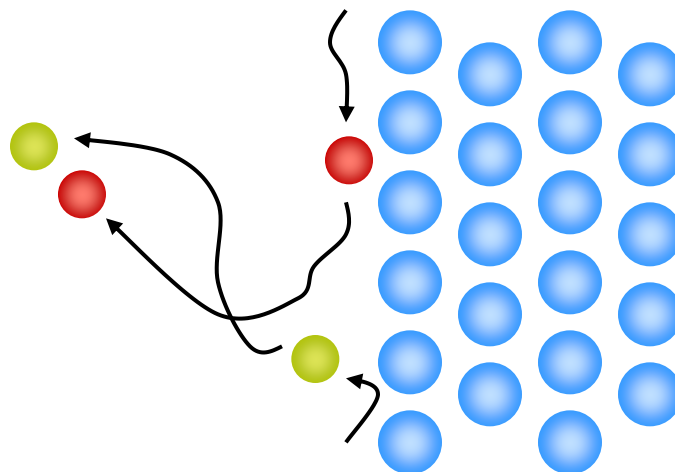


semiconductors

Nonadiabatic dissipation at metal surfaces

Weak Coupling Approx to Ehrenfest:

Molecular Dynamics with Electronic Friction, Martin Head-Gordon, JCT, *J. Chem. Phys.* **103**, 10137 (1995)



$$M \ddot{x} = -\nabla V(x) - \int_0^t dt' \Lambda(t, t') \dot{x}(t') + R(t),$$

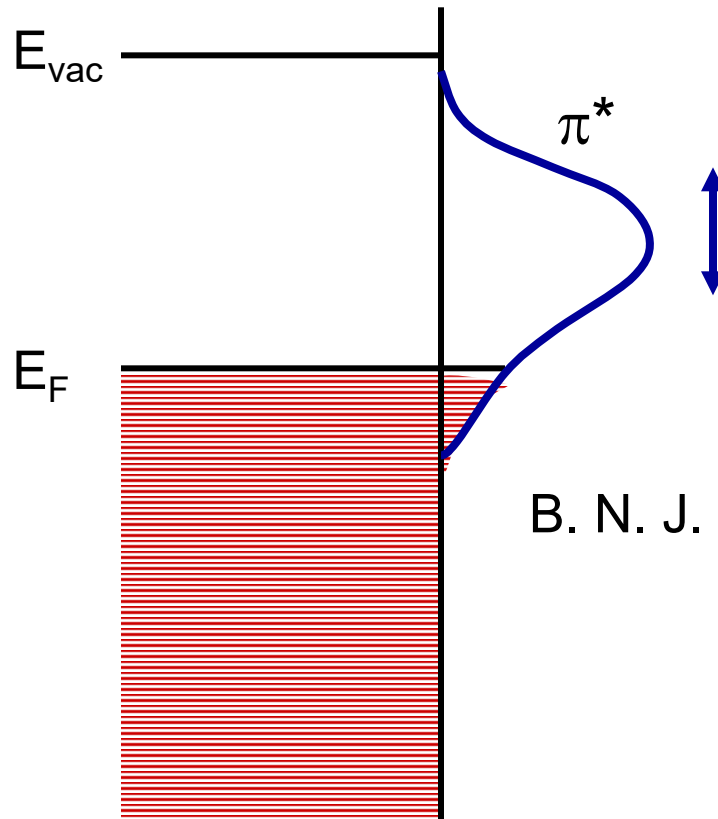
friction kernel fluctuating force

VIBRATIONAL LIFETIME: CO on Cu(100) (ps.)

MODE	EXPERIMENT	CALC (0K) (<i>Phonons</i>)	CALC(0K) (<i>Electrons</i>)
C – O stretch	2.5 +/- 0.5 (A. Harris et al.)	>100000	3.3 +/- 0.5

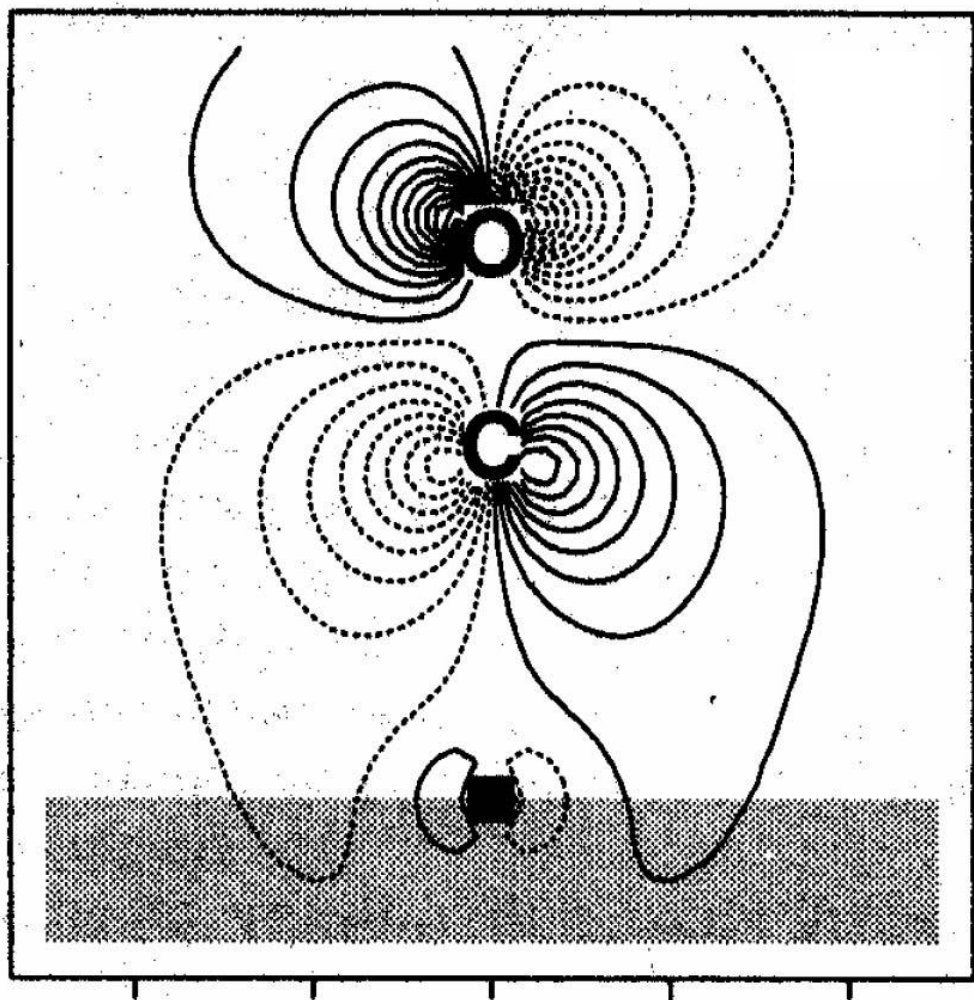
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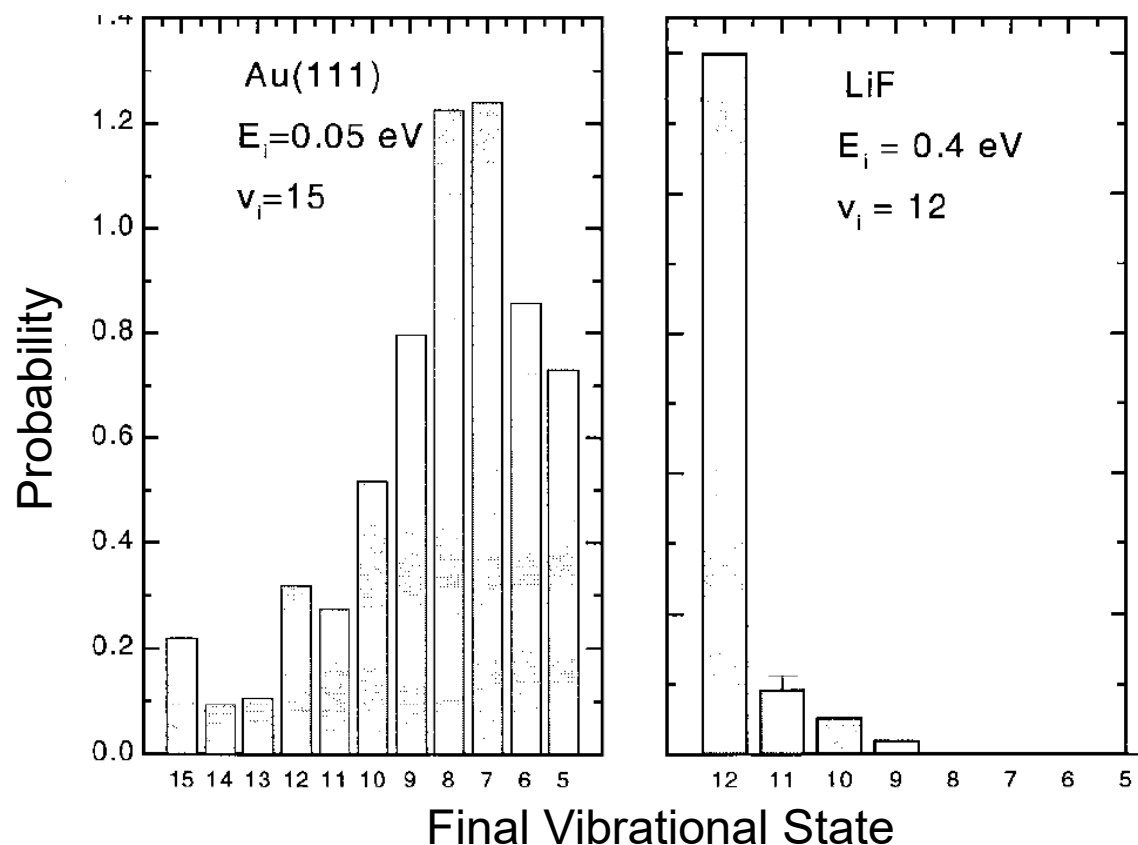
B. N. J. Persson and M. Persson, (1980)

$d\Psi_{\text{HOMO}}/dR$ (C -- O Stretch)



Scattering of vibrationally excited NO from Au(111)

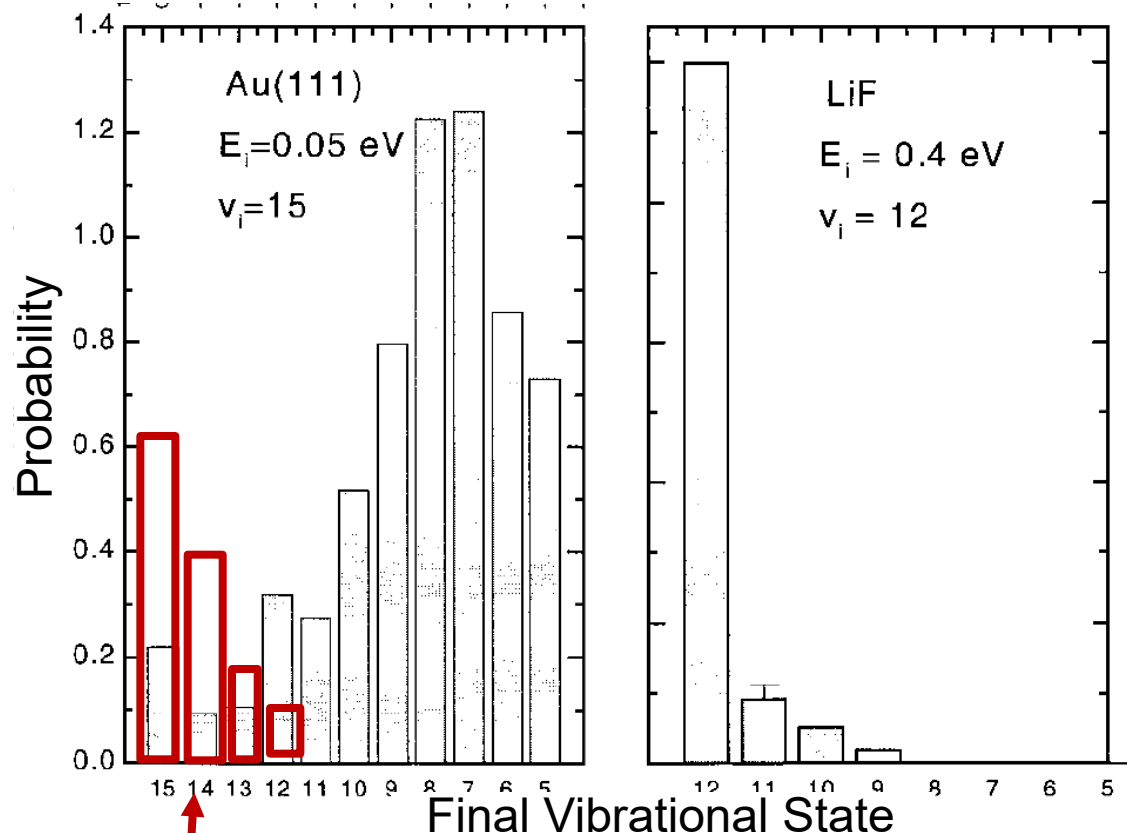
electron transfer: strong coupling



Huang, Rettner, Auerbach, Wodtke, *Science* **2000**, 290, 111.

Scattering of vibrationally excited NO from Au(111)

electron transfer: strong coupling



Friction:

Weak Coupling:

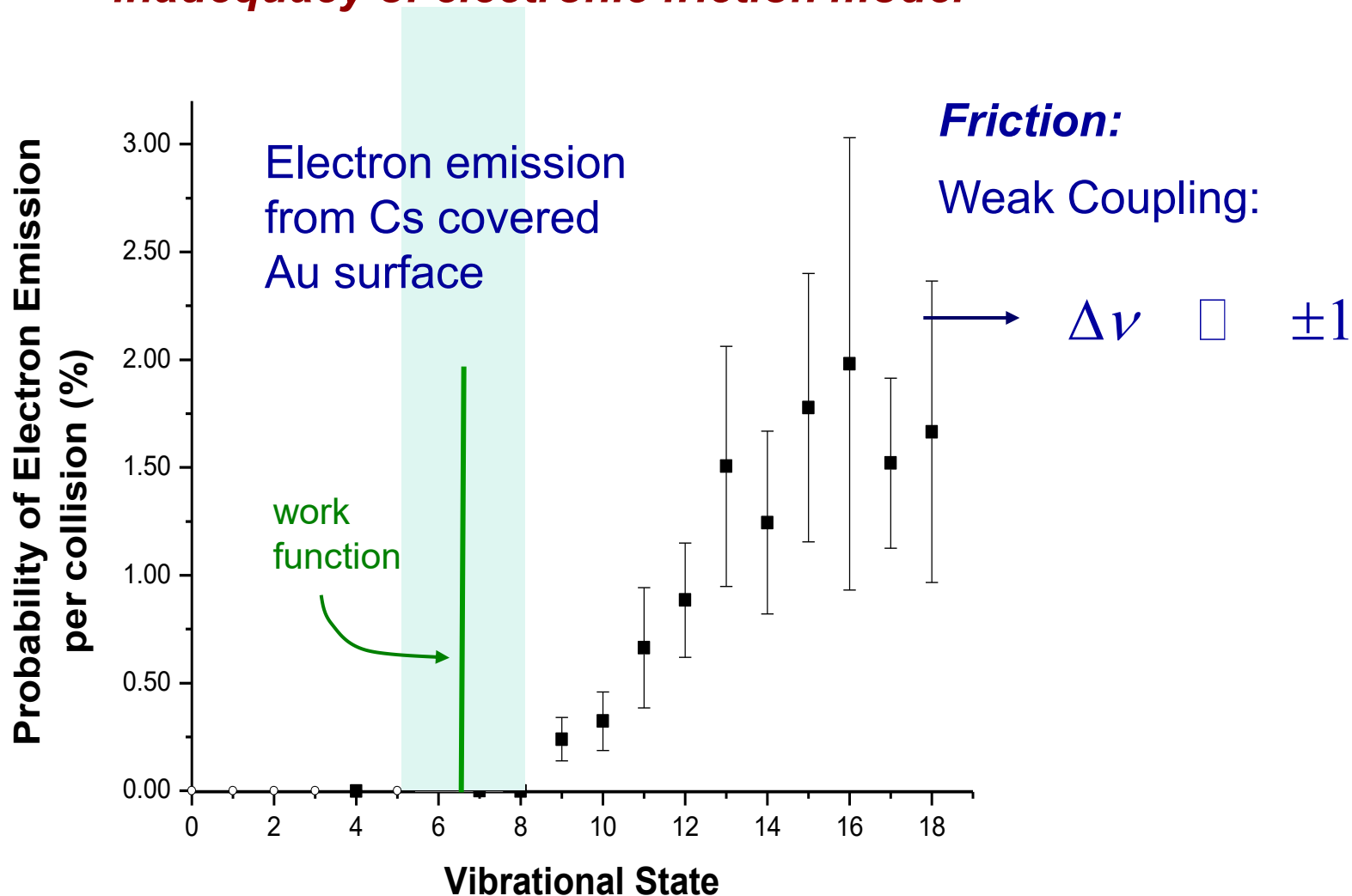
$$\longrightarrow \Delta v \approx \pm 1$$

*Electronic Friction
cannot account for
the observed huge
energy transfer*

Friction: from my memory !

Electron Emission:

→ *inadequacy of electronic friction model*

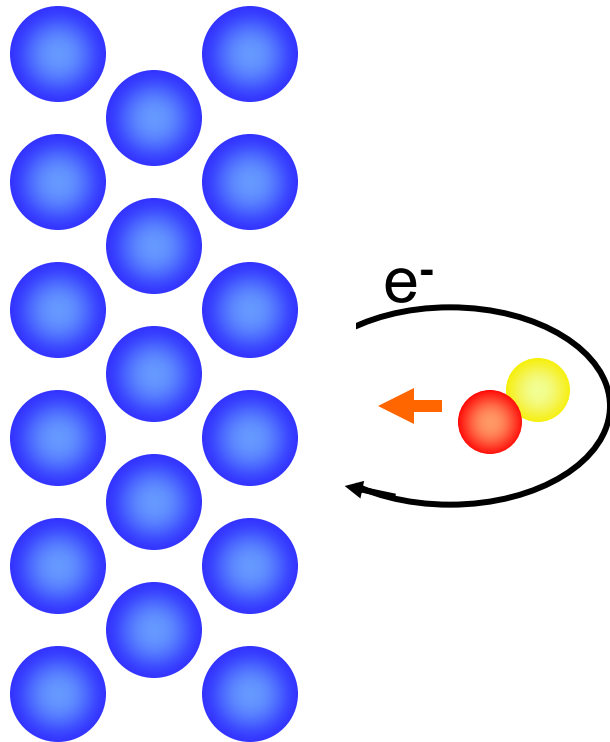


White, J. D.; Chen, J.; Auerbach, D. J.; Wodtke, A. M. *Nature* **2005**, 433, 503.

Electronic excitations at metal surfaces?

Inelastic Electron Scattering Picture:

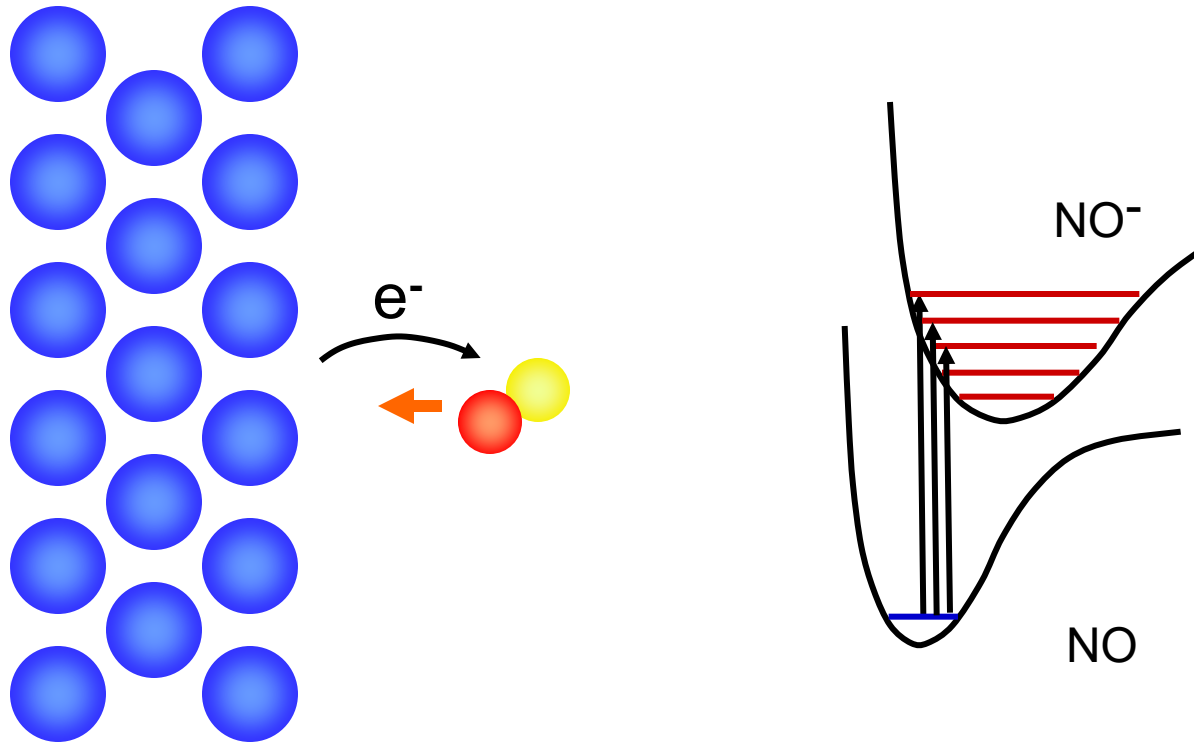
a. Non-Resonant \rightarrow *friction model*



Electronic excitations at metal surfaces?

Inelastic Electron Scattering Picture:

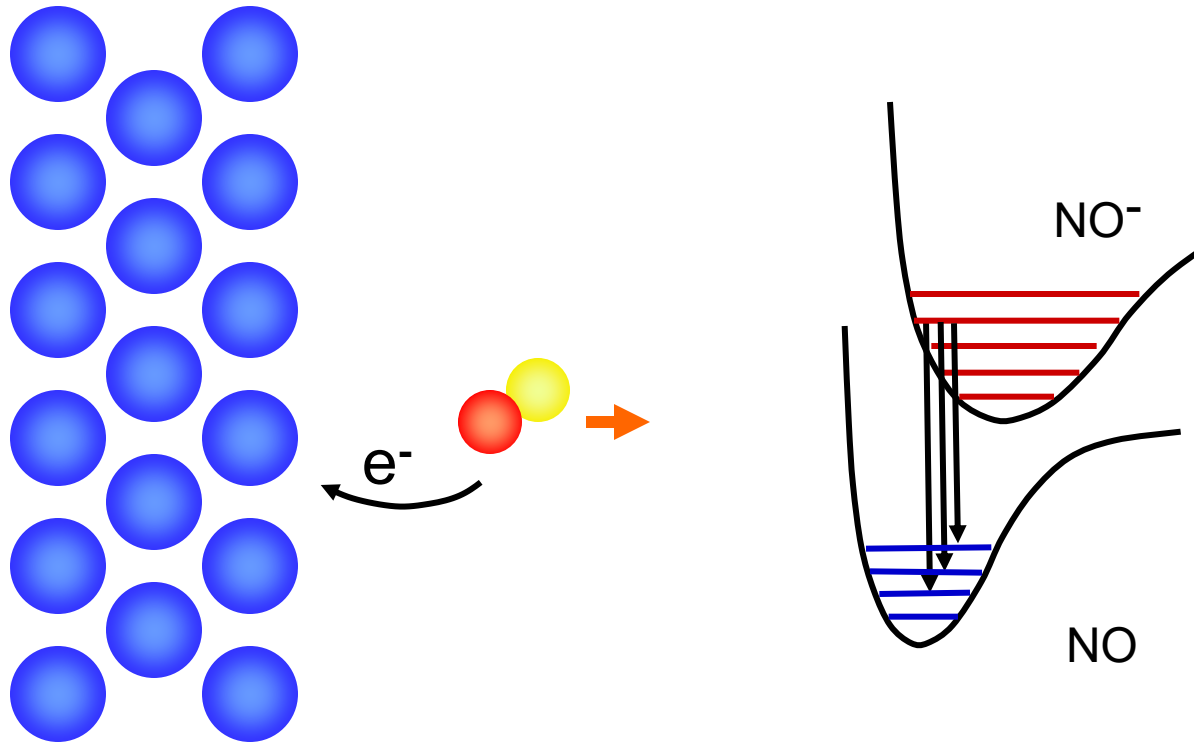
b. Resonant \rightarrow *transient negative ion*



Electronic excitations at metal surfaces?

Inelastic Electron Scattering Picture:

b. Resonant \rightarrow *transient negative ion*

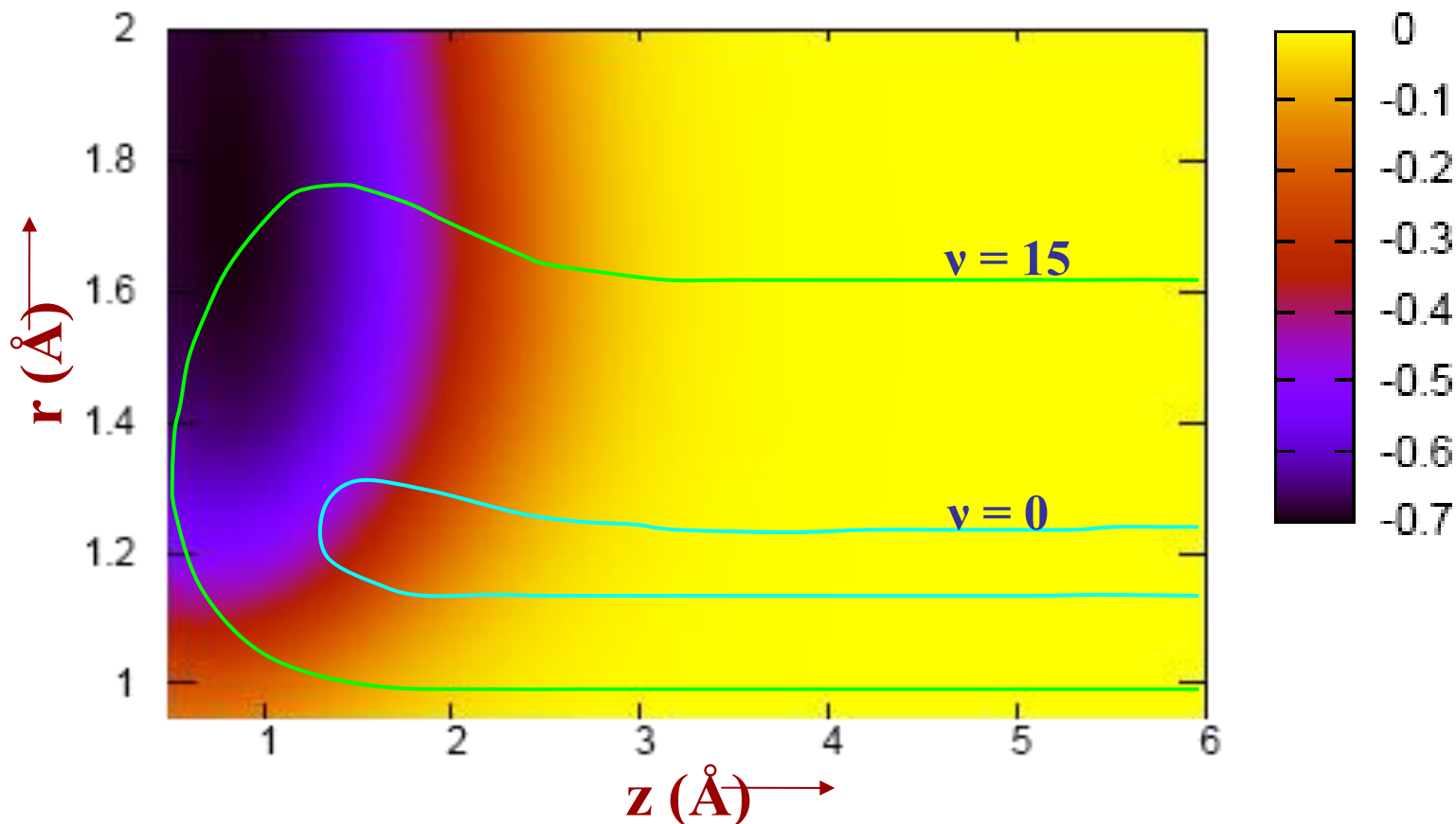


Franck-Condon picture (J. W. Gadzuk)

NO approaching a three-fold site on Au(111)

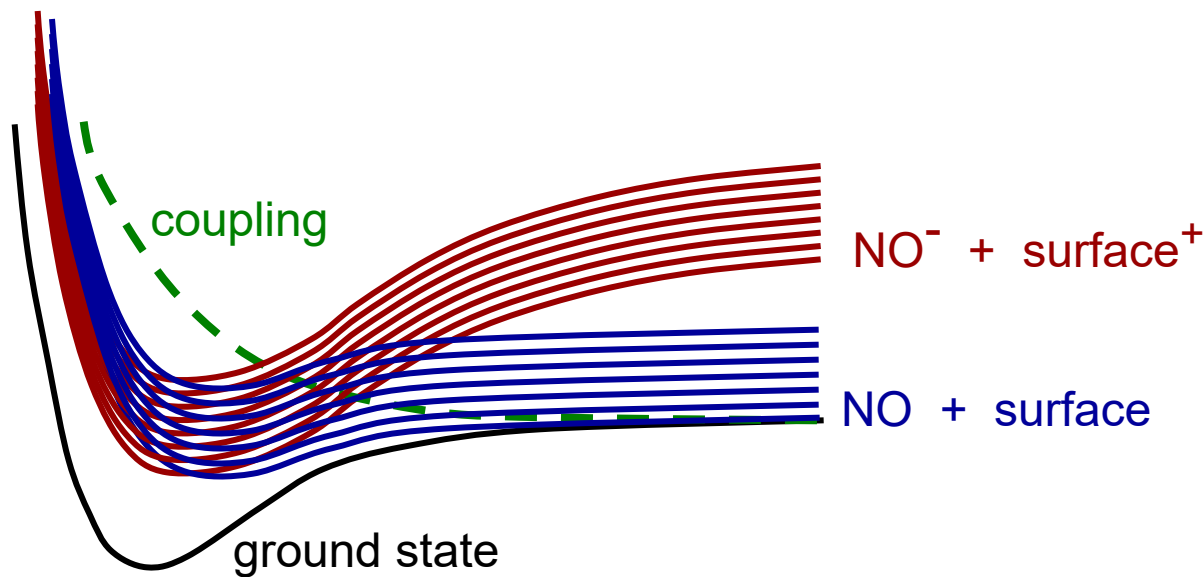
Ground state charge on NO

DFT (VASP, PW91)



Scattering of Nitric Oxide from a Metal Surface (Gold)

1. Calculations of potential energy surfaces and nonadiabatic couplings? *Newns-Anderson Hamiltonian*
2. Propagation of surface hopping trajectories? *Independent Electron Surface Hopping (IESH)*



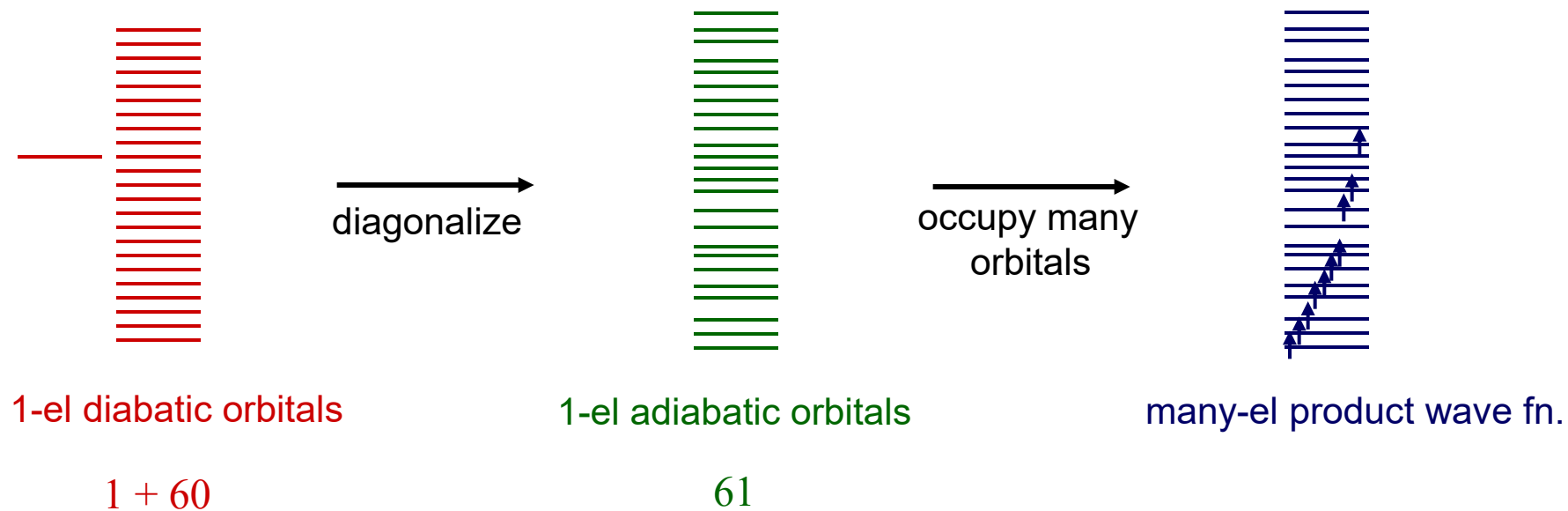
Sharani Roy



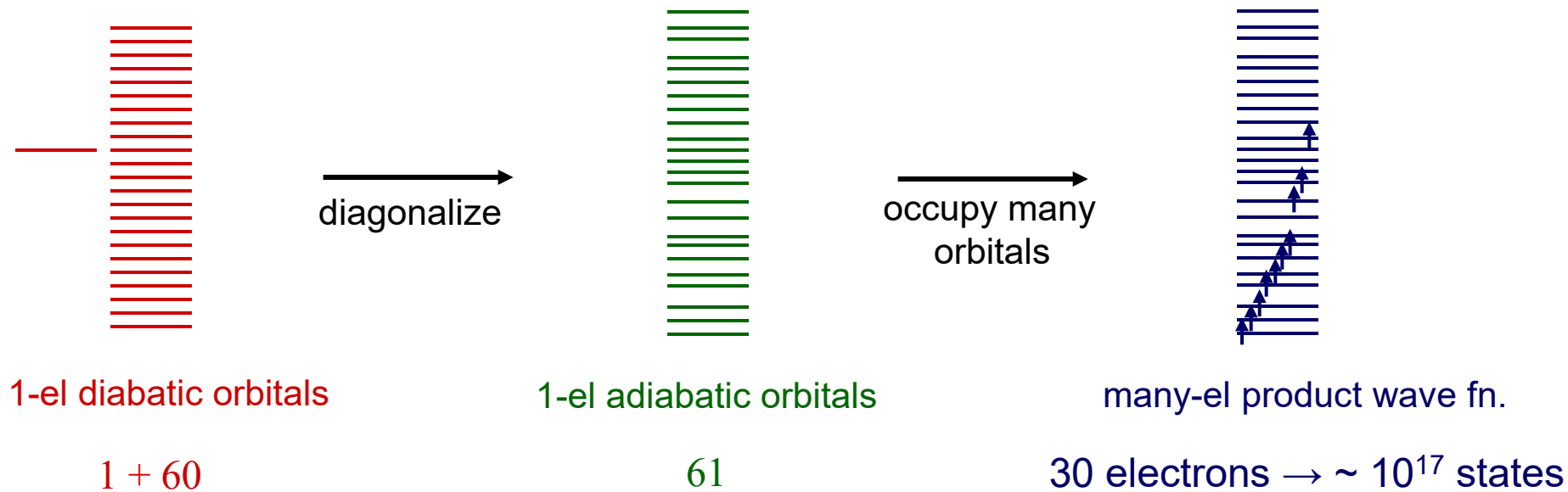
Neil Shenvi

January 14, 2026

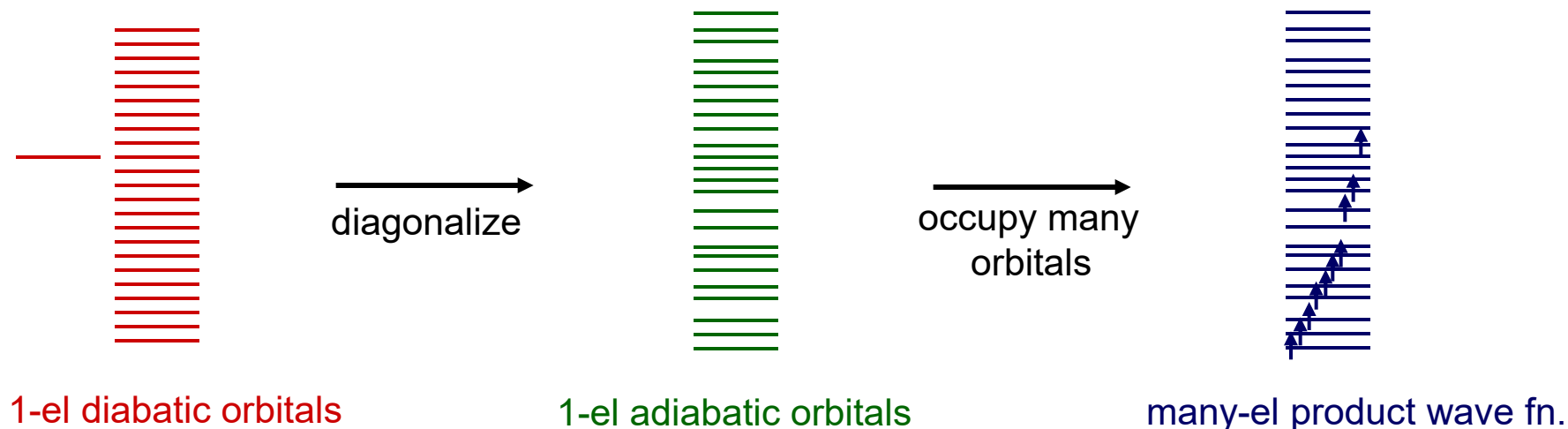
1. Map onto Anderson-Newns (tight binding) picture:



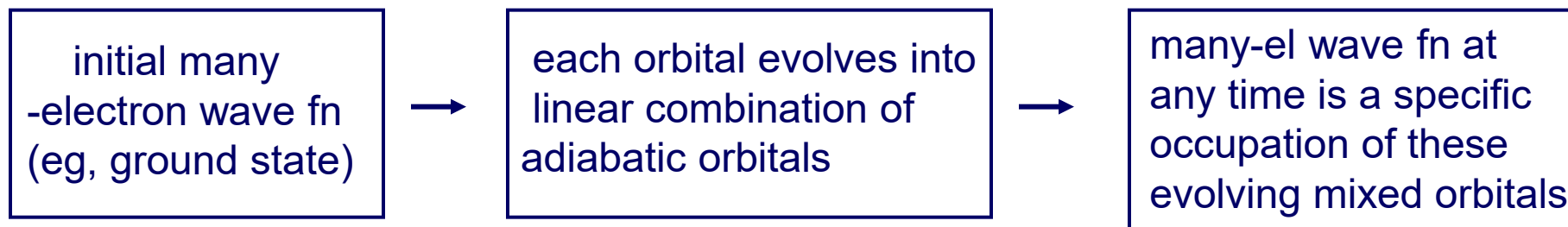
1. Map onto Anderson-Newns (tight binding) picture:



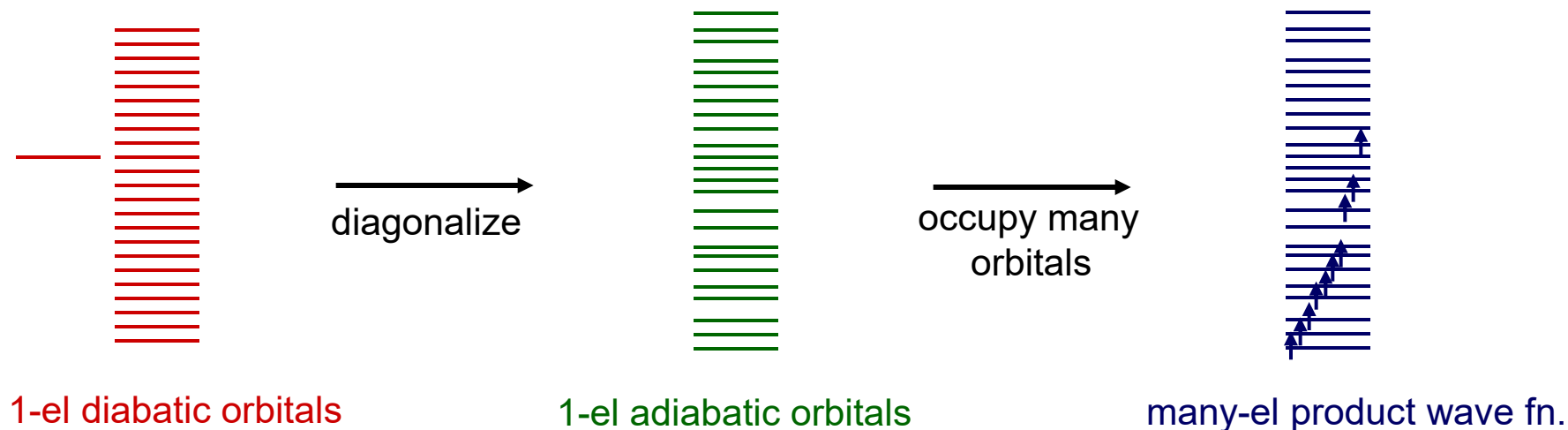
1. Map onto Anderson-Newns picture:



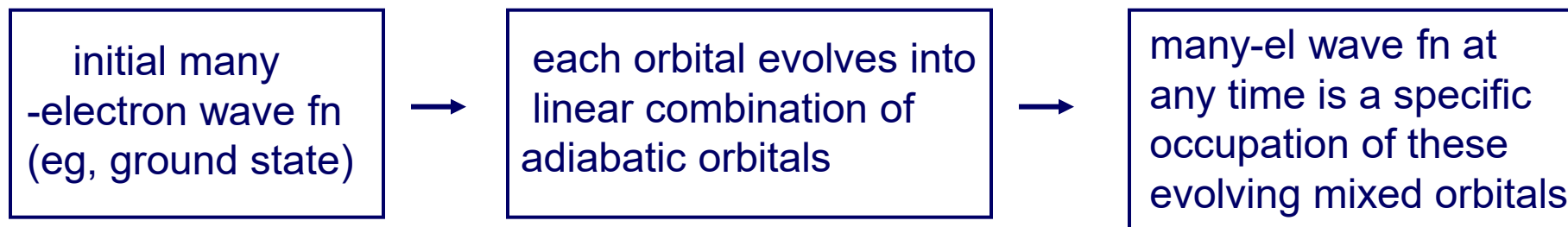
2. Evolve many-electron wave function along trajectory:



1. Map onto Anderson-Newns picture:



2. Evolve many-electron wave function along trajectory:

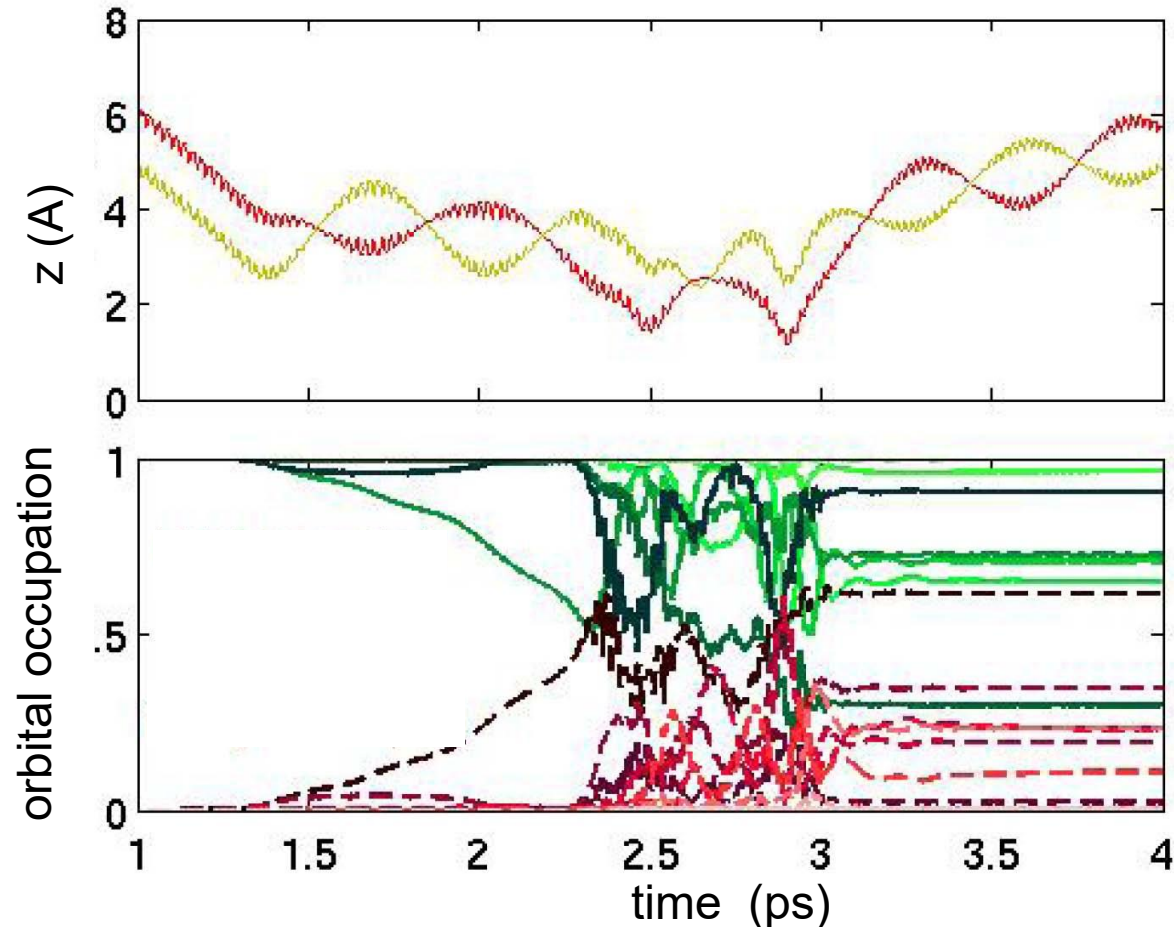


3. Fewest Switches Surface Hop among adiabatic orbitals

each electron independently

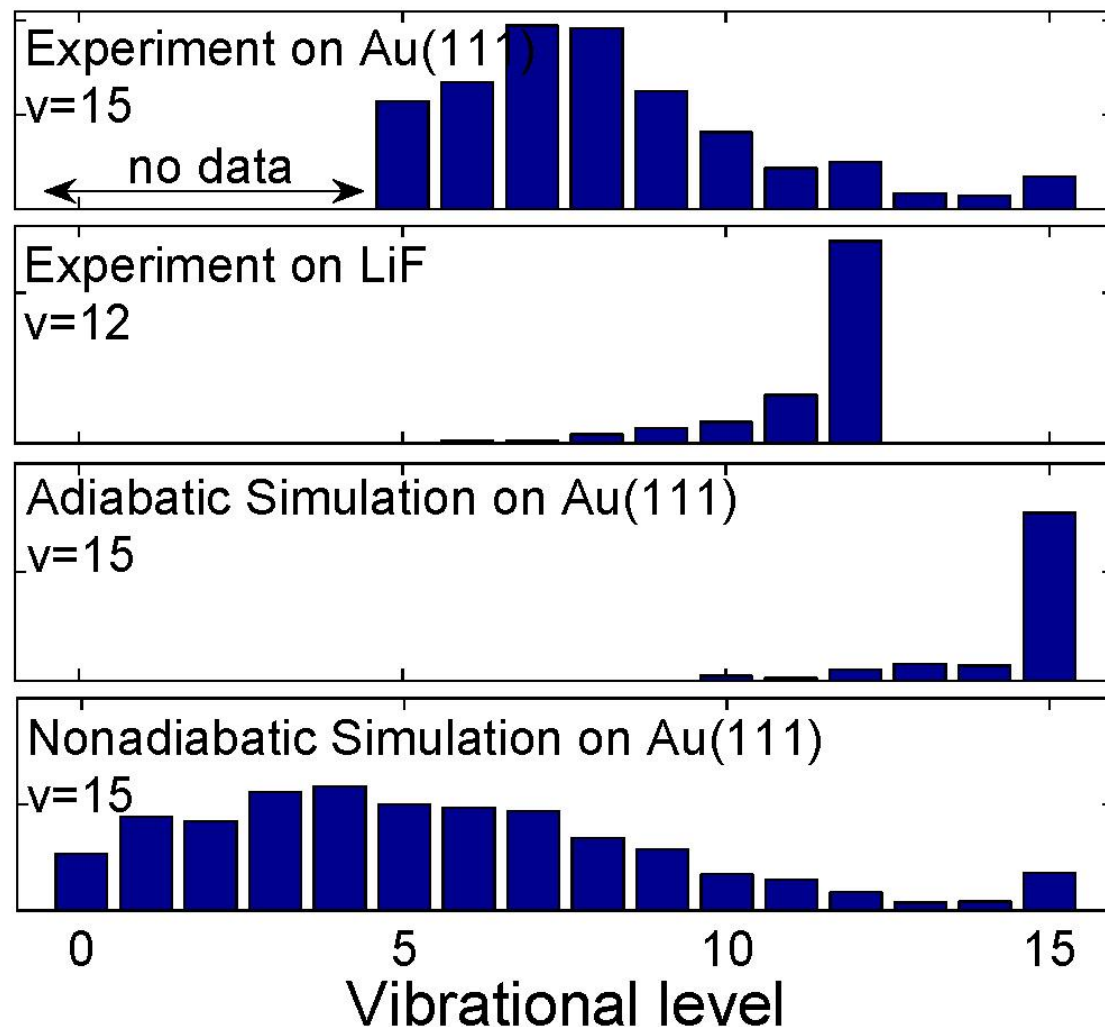
NO($\nu=15$) scattered from Au(111) (*surface hopping trajectory*)

132 moving Au atoms, accurate phonon spectrum, periodic boundaries, etc.

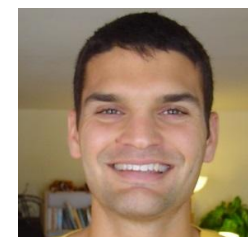


~ 100 surface hops typically encountered

Population



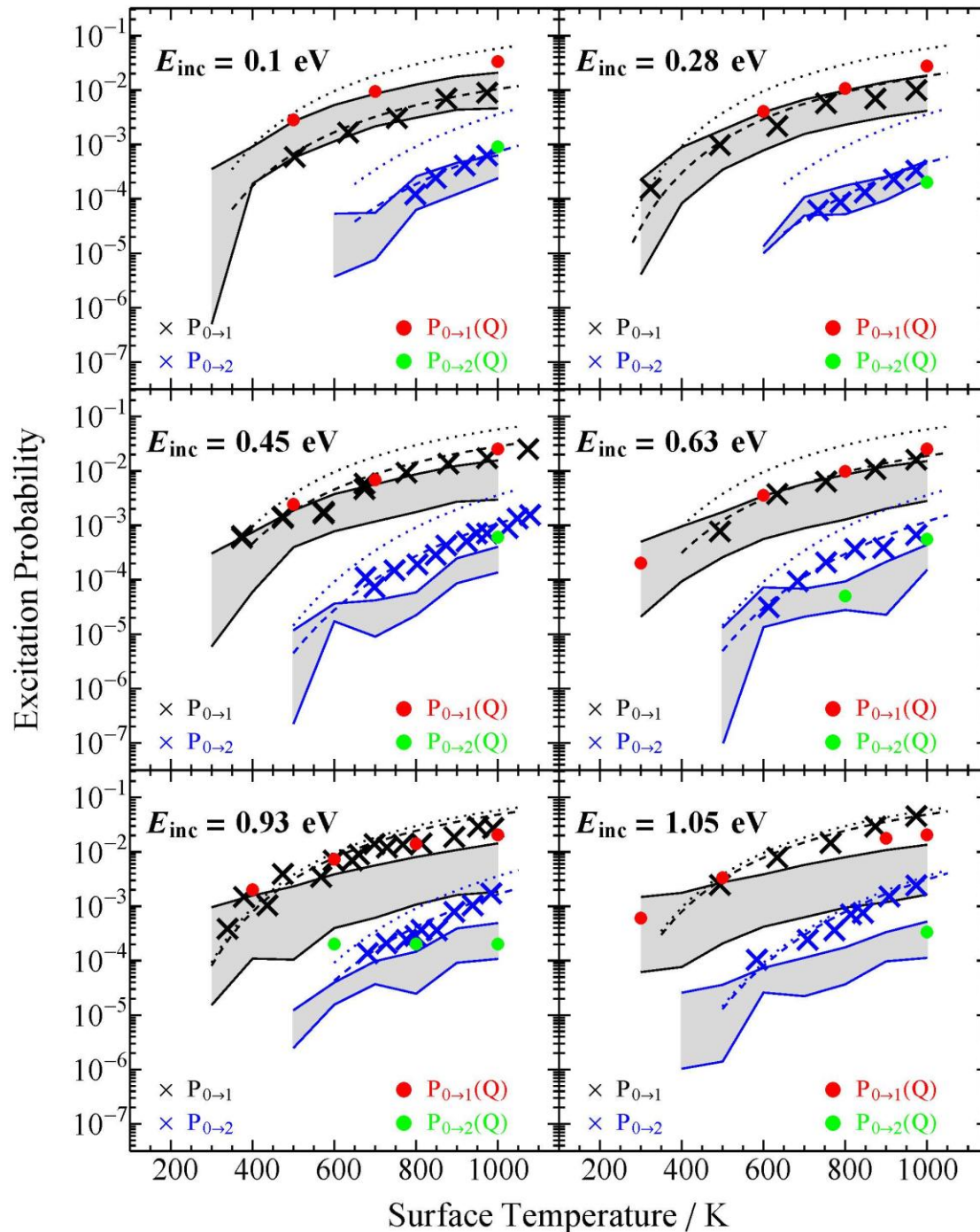
Sharani Roy



Neil Shenvi

Expts: Huang, Rettner, Auerbach, Wodtke,
Science **2000**, 290, 111.

NO(v=0) scattered from Au(111)



XXXX $P_{0 \rightarrow 1}$

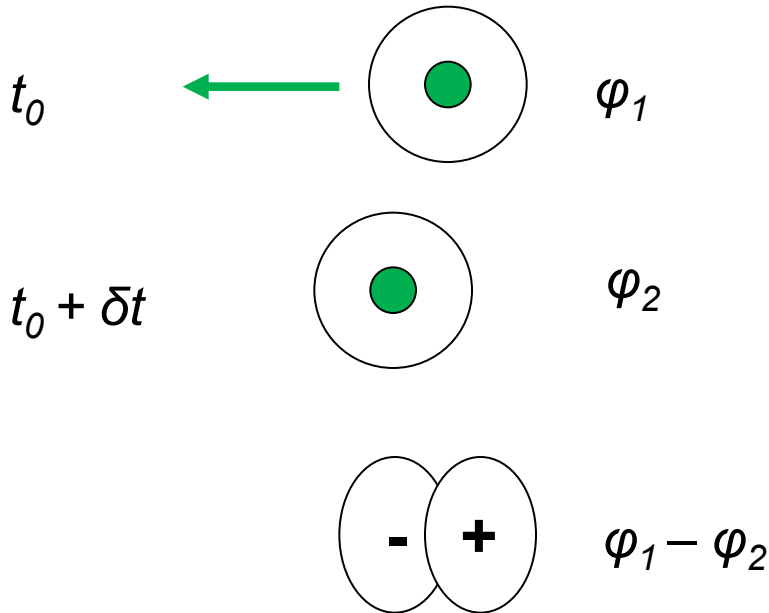
XXXX $P_{0 \rightarrow 2}$



Sascha
Kandratsenka

January 14, 2026

Translational and Angular Momentum Conservation ?



moving nucleus but
stationary electron

$$\varphi(z) \longrightarrow \varphi(z) \exp(i k_e z)$$

Translational and Angular Momentum Conservation ?

Direct Dynamics in Cartesian coordinates:

E_{trans} and J are not conserved

This is true for all nonadiabatic (Cartesian) direct dynamics methods that are based on separation of “fast” and “slow” particles, where energies and couplings of the fast particles are computed at fixed positions of the slow particles, and motion of the slow coordinates is governed by nonadiabatic coupling.

Examples:

Ehrenfest (mean field)

Surface Hopping

Exact Factorization

Nonadiabatic wave packet evolution

Translational and Angular Momentum Conservation ?

Some Strategies:

- Electron translational factors (Subotnik et al. 2023)

$$\varphi(z) \longrightarrow \varphi(z) \exp(i k_e z)$$

- Direct Dynamics with approximations to nonadiabatic couplings that depend only on internal coordinates
- Projection Operator to remove overall translation and rotation from nonadiabatic couplings (Truhlar et al. 2020)
- Fitted potentials and nonadiabatic couplings that depend only on internal coordinates. e.g., machine learning

VISTA: Virtual International Seminar on Theoretical Advancements

Thanks to **Alexey Akimov** for
inaugurating and leading **VISTA**:
100 Seminars!!



